



Equation of state for strongly interacting matter within a HTL quasiparticle model

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- QPM with $\text{Im} \Pi_i \neq 0$, plasmons and plasminos from HTL
- extrapolation of lattice QCD to large baryon densities, e.g. CBM@FAIR
- EOS for hydrodynamic phase of heavy-ion collisions (SPS, RHIC, LHC)

From QCD to thermodynamics

QCD: $\mathcal{L}$

→ propagators

→ self-energies



Blaizot, Rebhan, Iancu

thermodyn. potential Ω

→ state variables: $p, s, n_q, \dots$

→ EOS $e = e(p)$

→ $T^{\mu\nu}$, hydrodynamics

lattice QCD

→ available @ $\mu \simeq 0$

→ $\mu \neq 0$: sign problem

CJT formalism

- effective action

$$\Gamma[D, S] = I - \frac{1}{2} \left\{ \text{Tr} [\ln D^{-1}] + \text{Tr} [D_0^{-1} D - 1] \right\} \\ + \left\{ \text{Tr} [\ln S^{-1}] + \text{Tr} [S_0^{-1} S - 1] \right\} + \Gamma_2[D, S]$$

- translation-invariant systems, no broken symmetries

$$\frac{\Omega}{V} = \text{tr} \int \frac{d^4 k}{(2\pi)^4} n_B(\omega) \text{Im}(\ln D^{-1} - \Pi D) \\ + 2 \text{tr} \int \frac{d^4 k}{(2\pi)^4} n_F(\omega) \text{Im}(\ln S^{-1} - \Sigma S) - \frac{T}{V} \Gamma_2$$

2-loop QCD thermodynamics

- truncate Γ_2 at 2-loop order

$$\Gamma_2 = \frac{1}{12} \text{ (gluon self-energy)} + \frac{1}{8} \text{ (ghost self-energy)} - \frac{1}{2} \text{ (quark self-energy)}$$

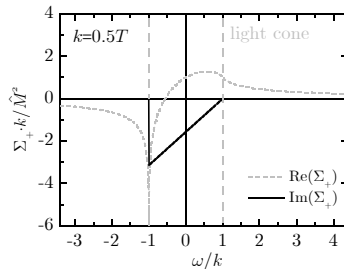
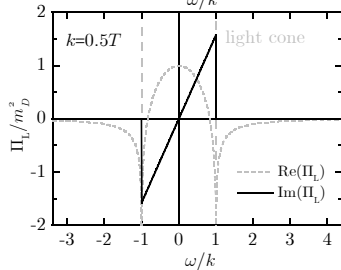
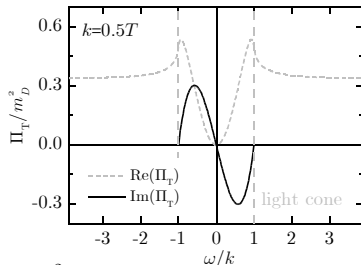
→ self-energies of 1-loop order

$$\Pi = \frac{1}{2} \text{ (gluon self-energy)} + \frac{1}{2} \text{ (ghost self-energy)} - \text{ (quark self-energy)}$$

- gauge invariance: HTL self-energies instead

HTL self-energies

- $\text{Im} \Pi \neq 0$ below the lightcone (solid lines)



→ Landau damping

Entropy

- stationarity of Ω

$$s := -\frac{1}{V} \left. \frac{\partial \Omega}{\partial T} \right|_{\mu} = -\frac{1}{V} \left(\left. \frac{\partial \Omega}{\partial T} \right|_{\text{expl.}} + \underbrace{\frac{\delta \Omega}{\delta D}}_0 \frac{\partial D}{\partial T} \right)_{\mu}$$

$$= s_{g,T} + s_{g,L} + s_{q,Pt} + s_{q,P1} + s', \quad s' = 0$$

Vanderheyden, Baym: JSP'98

- e.g. gluons:

$$s_{g,T} \sim \int_{d^4k} \frac{\partial n_B}{\partial T} \left\{ \underbrace{\pi \varepsilon(\omega) \Theta(-\text{Re} D_T^{-1})}_{\text{qp contribution}} + \underbrace{\text{Re} D_T \text{Im} \Pi_T - \arctan\left(\frac{\text{Im} \Pi_T}{\text{Re} D_T^{-1}}\right)}_{\text{damping terms}} \right\}$$

- HTL QPM model $\iff$ effective QPM
~~coll. modes, damping, asympt. disp. rel.~~

Blaizot, Iancu, Rebhan: PRD'01
 Rebhan, Romatschke: PRD'03

Bluhm, Kämpfer, RS, Seipt: EPJC'07

Pressure

$$s \sim \left(\frac{\partial \Omega}{\partial T} \right)_{\text{expl.}} + \underbrace{\frac{\delta \Omega}{\delta D}}_0 \frac{\partial D}{\partial T} \Big|_{\mu}$$

$$s_i \sim \int_{d^4k} \frac{\partial n_{B/F}}{\partial T} \left\{ \text{qp} + \text{damping} \right\}$$

- self-consistent formulation of the pressure

$$p = -\frac{\Omega}{V} := \sum_i p_i - B$$

$$p_i \sim \int_{d^4k} n_{B/F} \left\{ \text{qp} + \text{damping} \right\}$$

$$\frac{\partial B}{\partial T} := \sum_i \frac{\partial p_i}{\partial \Pi_i} \frac{\partial \Pi_i}{\partial T} \quad \left(\frac{\partial B}{\partial \mu} = \sum_i \ddot{\frac{\partial \Pi_i}{\partial \mu}} \right)$$

- entropy density

$$s = \frac{\partial p}{\partial T} = \sum_i s_i + \frac{\partial p_i}{\partial \Pi_i} \frac{\partial \Pi_i}{\partial T} - \frac{\partial B}{\partial T} = \sum_i s_i$$

- net quark density

$$n \sim \frac{\partial \Omega}{\partial \mu} \Big|_{\text{expl.}} + \underbrace{\frac{\delta \Omega}{\delta D}}_0 \frac{\partial D}{\partial \mu}$$

$$n_q = \frac{\partial p}{\partial \mu} \sim \int_{d^4k} \left(\frac{\partial n_F}{\partial \mu} + \frac{\partial n_F^A}{\partial \mu} \right) \left\{ \text{qp} + \text{damping} \right\}$$

Effective coupling

- fundamental parameter

$$g^2(x^2) = \frac{16\pi^2}{\beta_0 \ln(x^2)} \left(1 - \frac{4\beta_1}{\beta_0^2} \frac{\ln[\ln(x^2)]}{\ln(x^2)} \right)$$

- running coupling g^2 $x = \frac{\bar{\mu}}{\Lambda_{\text{QCD}}} \quad \bar{\mu} \sim T$

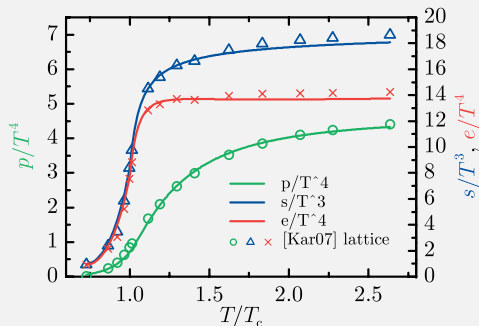


$$T > T_c, \mu = 0$$

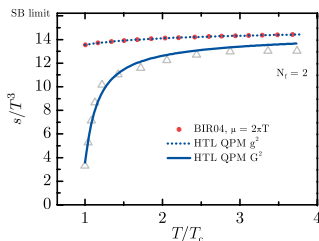
- effective coupling G^2 $x = \frac{\lambda}{T_c} (T - T_s)$

Adjustment @ $\mu=0$

- $\mu=0$: adjust to ℓ QCD
 - $\rightarrow T_s, \lambda$ fixed
 - $\rightarrow G^2(T, \mu=0)$

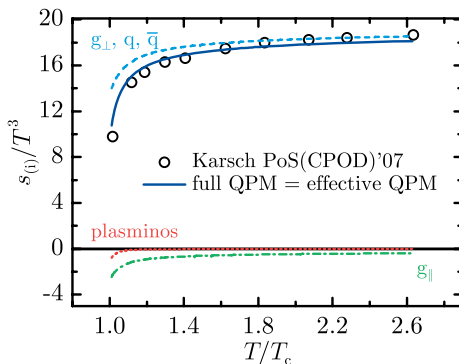


- $T_s=0$:

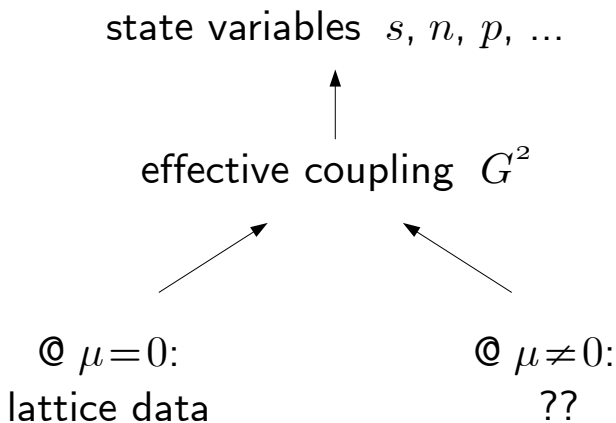


Influence of coll. modes + LD @ $\mu=0$

- individual entropy contributions



- Landau damping large close to T_c , decreases for higher temperatures



Into the T - μ -plane

- $\mu > 0$: stationary potential, self-consistent model
 $\rightarrow$ impose Maxwell's relation

$$\frac{\partial s}{\partial \mu} = \frac{\partial n}{\partial T} \quad \longrightarrow \quad a_T \frac{\partial G^2}{\partial T} + a_\mu \frac{\partial G^2}{\partial \mu} = b$$

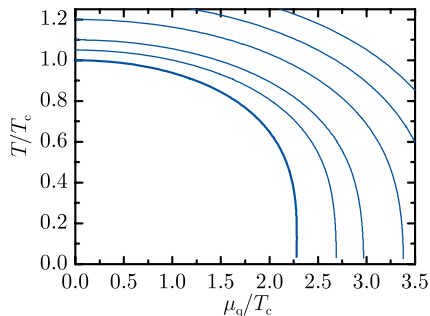
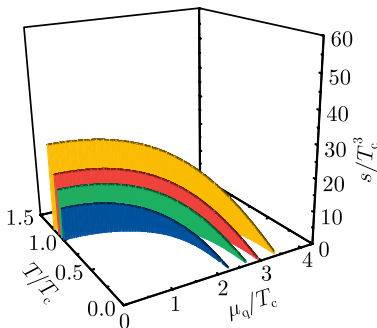
Peshier, Kämpfer, Soff: PRC'00, PRD'02

- solve quasilinear PDE for $G^2(T, \mu \neq 0)$ using method of characteristics
- test with lattice data for $\mu \simeq 0$: successful (eQPM)

Larger chemical potential



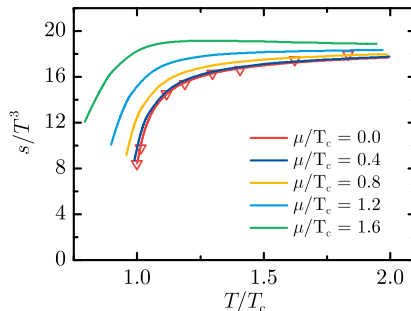
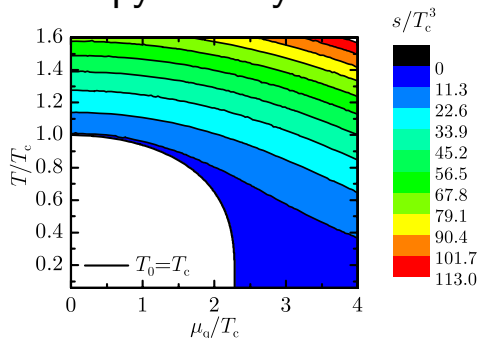
- full HTL QPM:
stable characteristics
- entropy density



RS, Bluhm, Kämpfer: EPJ ST'08

Results (1)

- entropy density

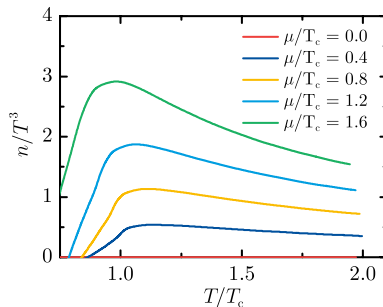
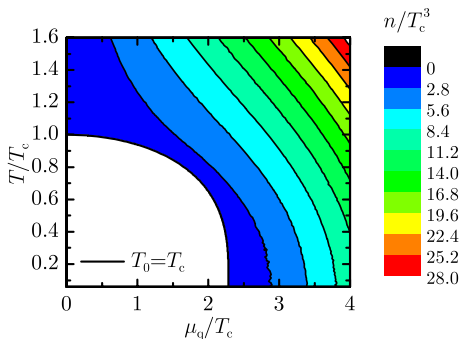


RS, Bluhm, Kämpfer: arXiv:0803.1571

- increases with temperature and chemical potential

Results (2)

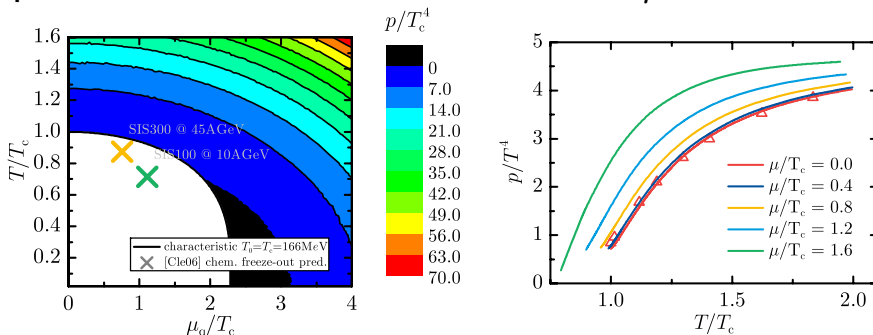
- net quark density



- small area of negative net quark density below transition line

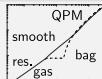
Results (3)

- pressure also increases with T and μ



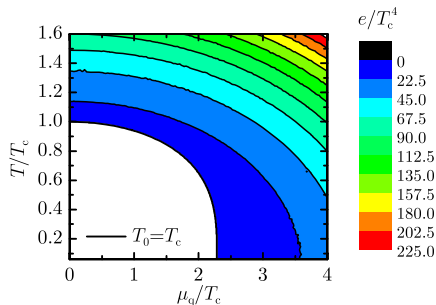
- small area of negative pressure also above transition line
 $\rightarrow$ no problems for EoS @ RHIC, LHC, SPS, FAIR

EOS for RHIC and LHC



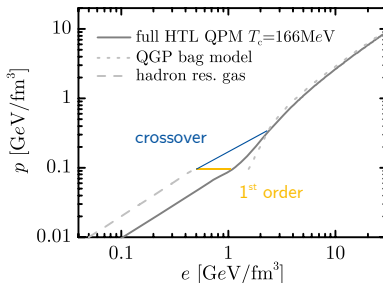
- energy density

$$e = -p + sT + \mu n$$



- EOS for LHC, RHIC

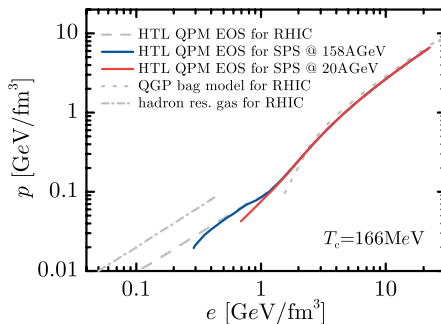
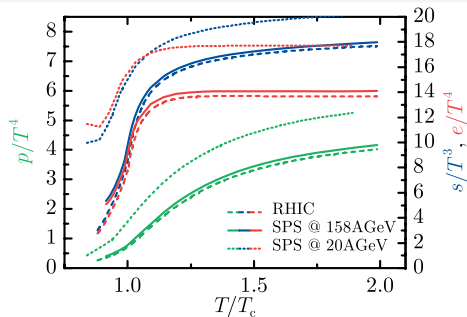
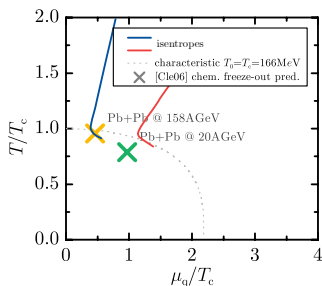
$$n_b/s \approx 0$$



EOS for SPS

PRELIMINARY

• SPS



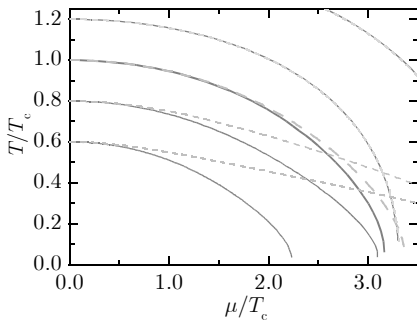
Summary & Outlook

- 2-loop Γ_2 + eff. coupling $G^2 \rightarrow$ HTL QPM
- Landau damping + collective modes
 $\rightarrow$ large μ accessible
- limitation: negative pressure for small T and μ
 $\rightarrow$ however deconfinement region @ LHC, RHIC, FAIR, SPS accessible
- outlook: EOS for FAIR, critical endpoint

Kämpfer, Bluhm, RS, Seipt: NPA'06

More effects of collective excitations

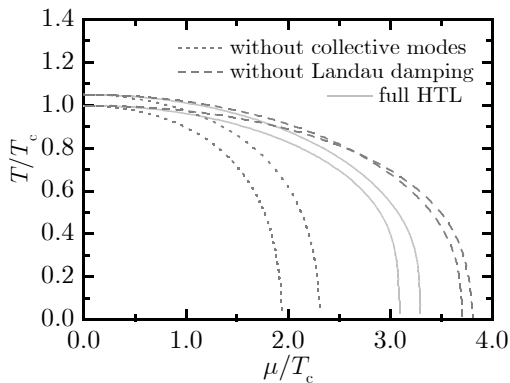
- collective modes
→ neg. entropy contrib.



situation improves

More effects of Landau damping

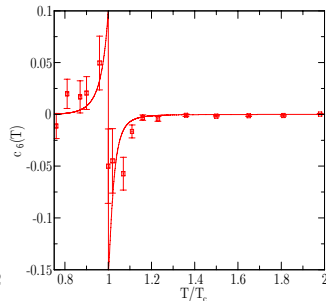
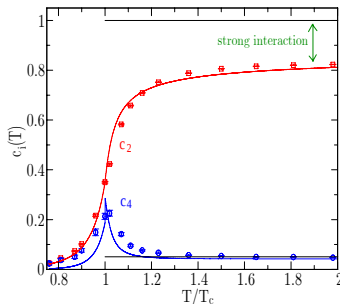
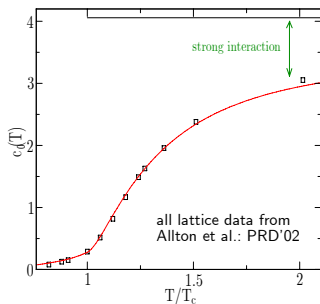
- only minor contribution at $\mu = 0$
- essential for $\mu > 0$



Small chemical potential

- $p(T, \mu \gtrsim 0)$ lattice data

$$p = T^4 \sum_n c_n(T) \left(\frac{\mu}{T}\right)^n \quad c_n(T) = \frac{1}{n!} \frac{\partial^n (p/T^4)}{\partial (\mu/T)^n}$$

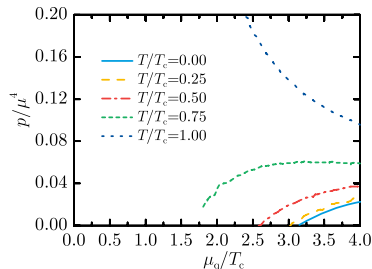
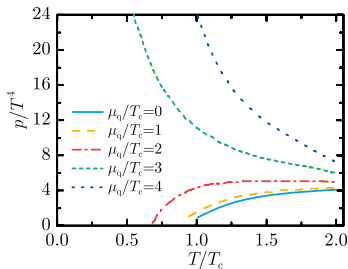
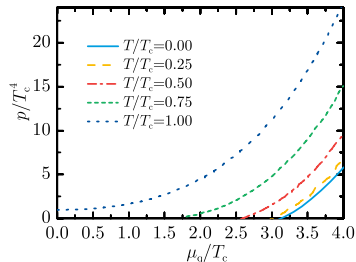
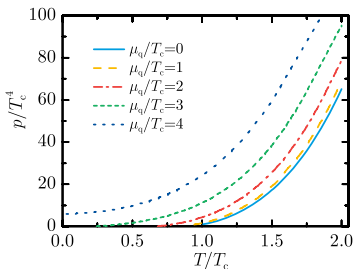


Bluhm, Kämpfer, Soff: PLB'05

- small $\mu \rightarrow$ effective QPM suffices

Results for the pressure (2)

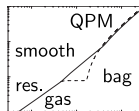
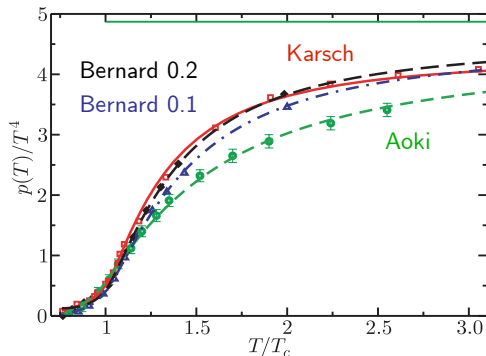
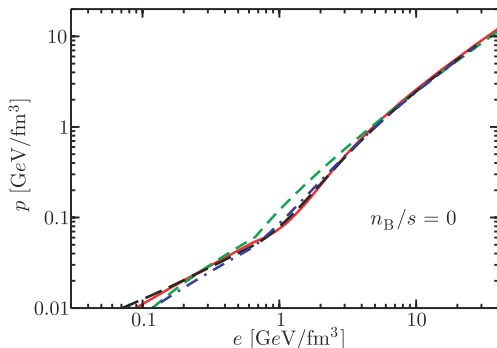
- pressure cuts



EOS for $N_f=2+1$

- RHIC, LHC:

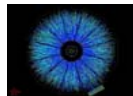
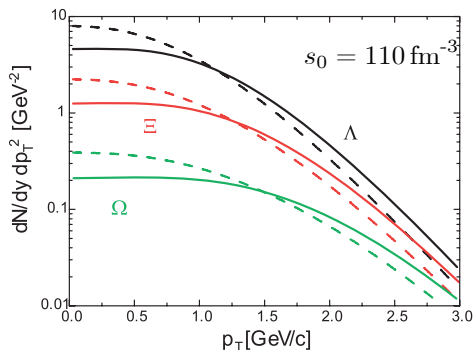
$$\mu = 0$$



Kämpfer, Bluhm, RS, Seipt, Heinz: NPA'05
 Bluhm, Kämpfer, RS, Seipt, Heinz: PRC'07

First test

- calculate elliptic flow using relativistic hydro code
- compare with experimental data (RHIC)



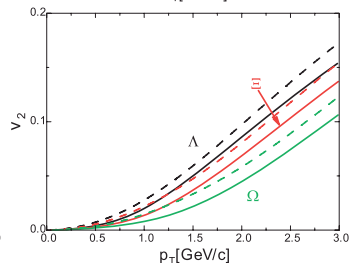
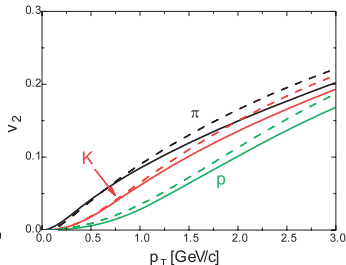
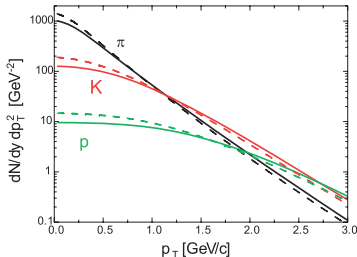
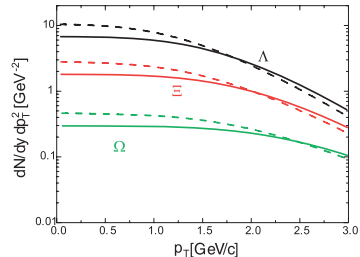
Predictions for LHC

- LHC Pb+Pb collisions - conservative guess:

$$s_0 = 330 \text{ fm}^{-3}, \quad \tau_0 = 0.6 \text{ fm}/c$$

$$b = 5.2 \text{ fm} \quad T_0 = 515 \text{ MeV}$$

- higher initial temperature
 $\rightarrow$ flatter p_T spectra
 $\rightarrow$ smaller v_2



More LHC predictions

- initial parameters translate to

$$e_0 = 127 \text{ GeV}, \quad p_0 = 42 \frac{\text{GeV}}{\text{fm}^3}, \quad T_0 = 515 \text{ MeV}$$

- LHC: higher initial temperature $\rightarrow$ longer fireball lifetime $\rightarrow$ stronger radial flow $\rightarrow p_T$ spectra flat

More full HTL quasiparticle model

- now: $\boxed{\text{Im } \Pi \neq 0}$ + collective excitations

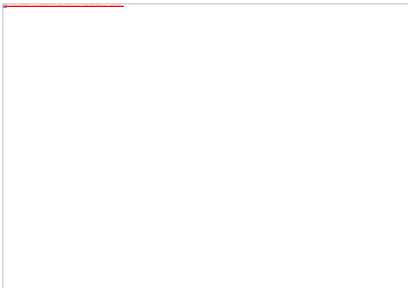
$$s = s_{qp} + s_{damp} = s_{qp} + (\tilde{s} - s_{qp})$$

$$\tilde{s} = \int d\omega \underbrace{\int dk \sigma(\omega, k)}_{\hat{= s_{qp}}(\omega)} \cdot F(\text{Im} \Pi(\omega, k)) \quad \xi := \frac{\text{Im} \Pi}{\text{Re} D^{-1}}$$

$$F := -\frac{1}{\pi} \left(\frac{\xi^2(\omega)}{(1+\xi^2(\omega))^2} + \frac{\xi^2(-\omega)}{(1+\xi^2(-\omega))^2} \right) \frac{\partial \xi}{\partial \omega}$$

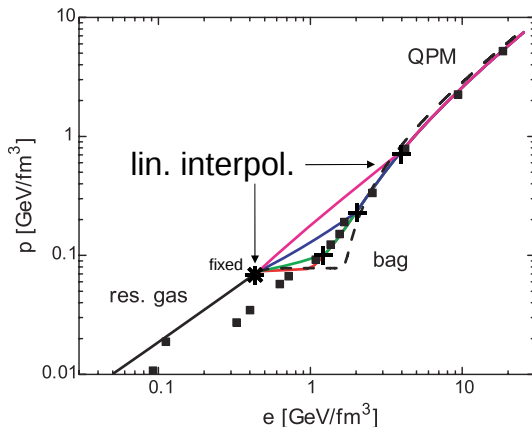
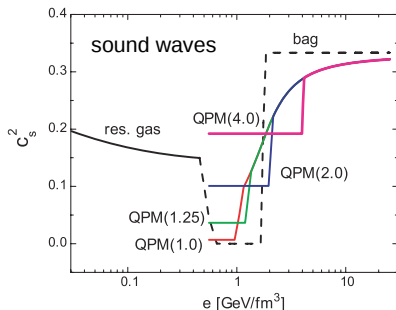
Backup

- model describes all available quantities:



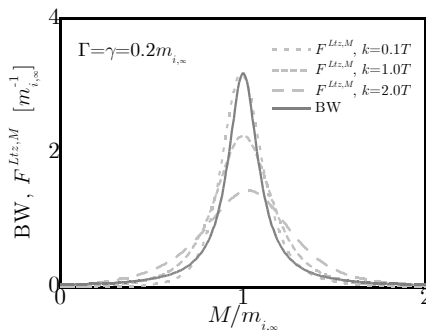
A family of EOS's $\mu_B \ll T$

- interpolate between hadron gas and QPM description



Backup: Inclusion of widths

- Peshier: $\text{Im } \Pi = 2\gamma\omega$

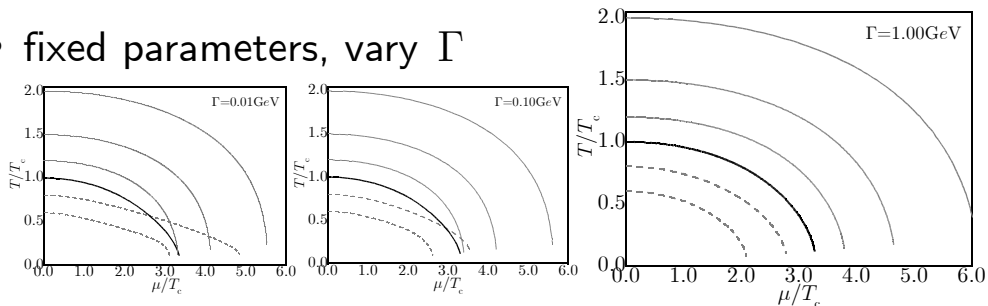


- ansatz $F(\omega, k) \rightarrow \text{BW}(m)$

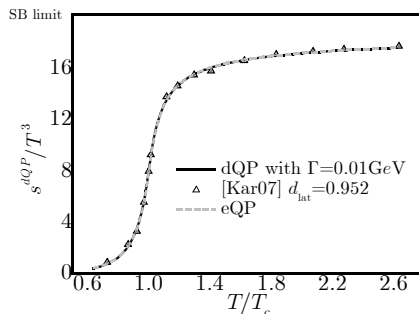
$$s(T) = \int dM s_{qp}(T, M) \text{BW}(m, M, \Gamma)$$

Backup: Distributed quasiparticle model

- fixed parameters, vary Γ



- adjustment to lattice
 $\Gamma = 0.01 \text{ GeV}$



Backup: Distributed quasiparticle model II

- bias adjustment $\Gamma \stackrel{!}{=} 1 \text{ GeV}$

