

Ekman boundary layers in a fluid filled precessing cylinder

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Astro MHD

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HzDR

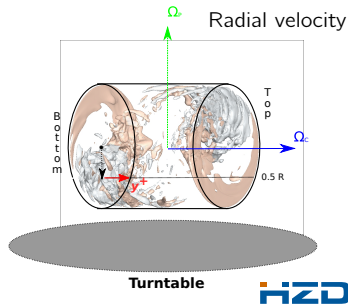
 **HELMHOLTZ**
| ZENTRUM DRESDEN
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- 1 Scope of the study
- 2 Ekman layer theory
- 3 Application to precessing cylinder endwall and results from simulation
- 4 Extrapolation and connection with former results

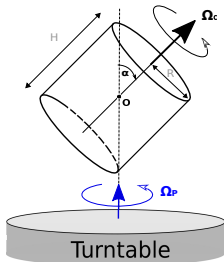
Research questions and problems

- What is the behaviour of the boundary layer on the cylinder endcaps
- How it evolves in the phase space of Ek (cylinder rotations) and Po (precession force strength) numbers? → Instabilities, turbulence...
- Does the endwall boundary layers interact/adapt with the bulk flow?

Problem: finding reference scales for the study of boundary layer (thickness, reference velocity). The precession driven flows are intrinsically 3D



Math formulation and tools

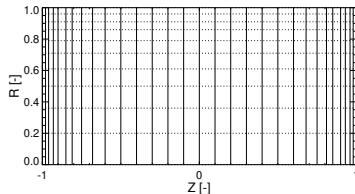
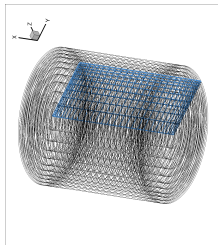


$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} = -\nabla p - 2Po \hat{\mathbf{k}} \times \mathbf{u} + Ek \nabla^2 \mathbf{u}$$

$$\nabla \cdot \mathbf{u} = 0$$

$$Po = \frac{\Omega_p}{\Omega_c} \quad Ek = \frac{\nu}{\Omega_c R^2} \quad \Gamma = \frac{H}{R}$$

⇓ Numerical



Solving through DNS code SEMTEX <http://users.monash.edu.au/~bburn/semtex.html>

General theory of Ekman layer

- Ekman boundary layer: type of boundary layer developed in rotating systems, characterized by the balance between Coriolis and viscous forces
- Study methods: math (stability analysis), DNS, lab experiment, atmospheric measurements. Linear laminar theory: scaling with $\sqrt{Ek}$

Caldwell, Van Atta, J. Fluid Mech. (1970), 44

Main parameter to classify the Ekman layer:

$$Re_{\delta} = \frac{U_{ref} \delta_{Ek}}{\nu}$$

- U_{ref} is the reference velocity (geostrophic) at the BL edge

- δ_{Ek} the thickness
 $\approx \sqrt{Ek} = \sqrt{\nu/(\Omega L^2)}$

Type of instability	Quantity	Theory		Experiment		
		Faller & Kaylor	Lilly	Faller, Faller & Kaylor	Tatro <i>et al.</i>	Present work
Type II (class A)	Critical Reynolds number	55	55	< 70	56.3+	56.7
	Wavelength, in Ekman depths	24	21	22 to 33	27.8 ± 2.0	110.8 Ro
	$\epsilon \uparrow$	-15°	-20°	+5° to -20°	0 to -8°	61% of that predicted by Lilly)
	Velocity divided by geostrophic	0.50	0.57	—	0.16	
Type I (class B)	Critical Reynolds number	118	110	125 ± 5	124.5+	—
	Wavelength, in Ekman depths	11	11.9	10.9	11.8	—
	$\epsilon \uparrow$	10°-12°	8°	+14.5 ± 2.0°	+14.8° ± 0.8°	—
	Velocity divided by geostrophic	0.33(11°)	0.094	0.023(14.5°)	0.034	—

$\epsilon \uparrow$ is the angle between wavefront and geostrophic wind.

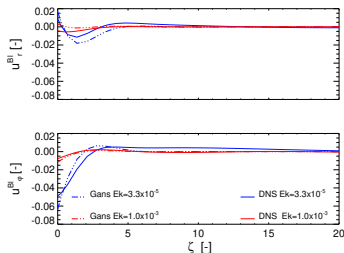
TABLE 1. Summary of Ekman layer instabilities

Fully turbulent Ekman layer for $Re_{\delta} > 150$

Idea for precession problem (weak regime)

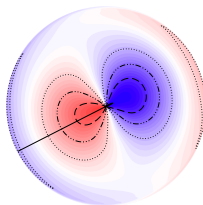
First step: decomposing and filtering the DNS.

$$\begin{aligned}
 u_{DNS} = & \underbrace{\sum_{k=1}^K A_{00k} u_{00k}(r)}_{\text{axisymmetric}} + \underbrace{\sum_{m=1}^M \sum_{k=1}^K \frac{1}{2} [A_{m0k} u_{m0k}(r, \varphi)] + c.c.}_{\text{non-axisymmetric}} \\
 & \underbrace{\sum_{n=1}^N \sum_{k=1}^K \frac{1}{2} [A_{0nk} u_{0nk}(z, r) + c.c.]}_{\text{axisymmetric oscillation}} \\
 & + \underbrace{\sum_{m=1}^M \sum_{n=1}^N \sum_{k=1}^{2K} \frac{1}{2} [A_{mnk} u_{mnk}(z, r, \varphi) + c.c.]}_{\text{inertial waves}} + u^{Bl}
 \end{aligned}
 \Rightarrow$$

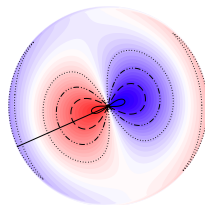


$$A_{mnk} = \int_V u_{DNS} \cdot u_{mnk}^* dV / N_{mnk}$$

(m, n, k)	$ A_{mnk} $
(1, 1, 1)	6.97×10^{-2}
(0, 0, 2)	2.02×10^{-3}
(0, 2, 1)	1.79×10^{-3}
(2, 0, 1)	1.42×10^{-3}
(1, 3, 2)	9.83×10^{-4}
(1, 1, 2)	6.58×10^{-4}
$(1, 1, 1)^-$	6.26×10^{-4}
(1, 5, 1)	5.89×10^{-4}
$Ek = 3.3 \times 10^{-5} \quad Po = 10^{-3}$	



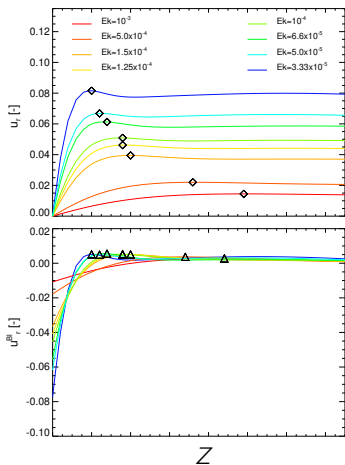
DNS



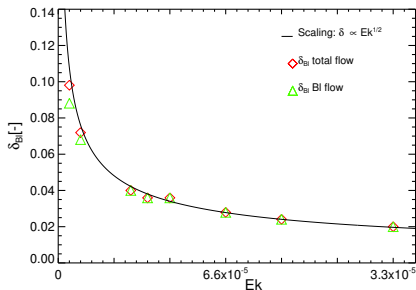
Linear analytical theory

Thickness in weak precession

Precession driven flows in cylinder present Ekman layers on the endwalls. (Meunier, Gans, Zhang, Kong...)

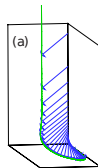


Checking the scale $Po/\sqrt{Ek} \ll 1$



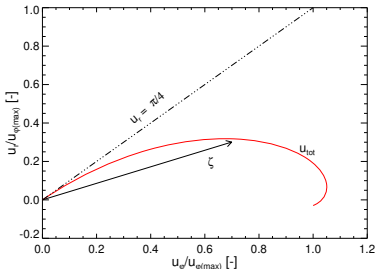
Particular features of precessing Ekman layer

Precession Ekman layer has the scaling of pure rotating flows but...
there are different behaviours: Ekman Spiral.



Spalart et al. *Phys. of Fluids*
(2008), 20, 101507

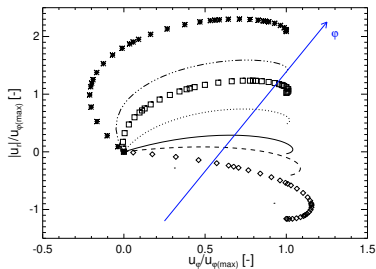
Spiral at $\varphi \approx 0$ for $Po/\sqrt{Ek} \ll 1$



$$m = 1$$

$$\Rightarrow$$

Poincaré force gives an azimuthal distribution.



The φ impacts on the axial vorticity of the flow (bulk and BL).

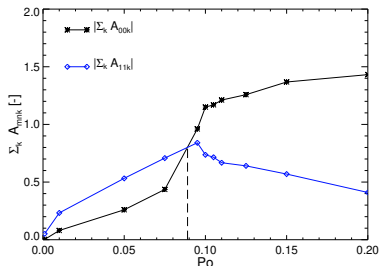
General study (up to Po)

Kong, et al., *Geo. e Astro. Fluid Dynamics* (2014)

$$\begin{aligned}
 \mathbf{u}_{DNS} = & \underbrace{\sum_{k=1}^K A_{00k} \mathbf{u}_{00k}(r)}_{\text{axisymmetric}} + \underbrace{\sum_{m=1}^M \sum_{k=1}^K \frac{1}{2} [A_{m0k} \mathbf{u}_{m0k}(r, \varphi)] + c.c.}_{\text{non-axisymmetric}} \\
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 \end{aligned}$$



$Re = 10000, \alpha = 90^\circ$

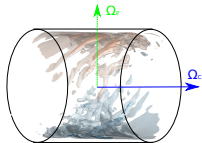


Emergence of an axisymmetric-geostrophic current for large Po .

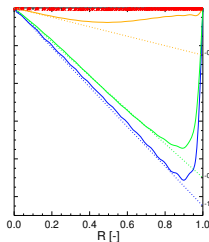
Amplitude A_{00} with azimuthal and axial wave number $(m, n) = (0, 0)$.

$$A_{mnk} = \int_V \mathbf{u}_{DNS} \cdot \mathbf{u}_{mnk}^* dV / N_{mnk}$$

$u_z \rightarrow 0$ in the bulk for large Po



$$U_g(r) \hat{\boldsymbol{\varphi}} = \sum_k A_{00k} \mathbf{u}_{00k}(r)$$



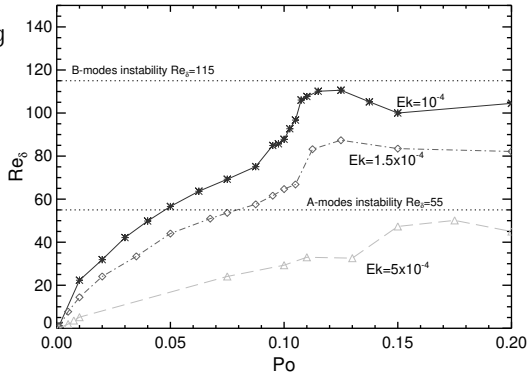
$$\Omega_g = \left. \frac{dU_g(r)}{dr} \right|_{r=0}$$

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Characterization of Ekman layer

Definition of Re_δ in precessing cylinder:

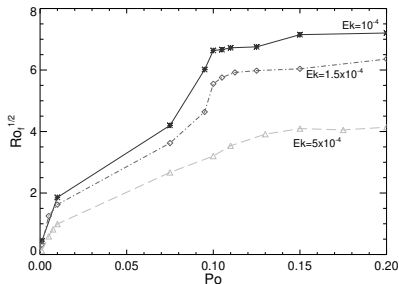
$$Re_\delta = \frac{(\Omega_g R^2) \sqrt{Ek}}{\nu}$$



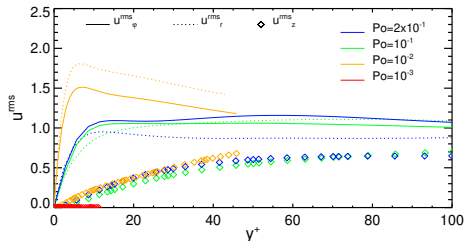
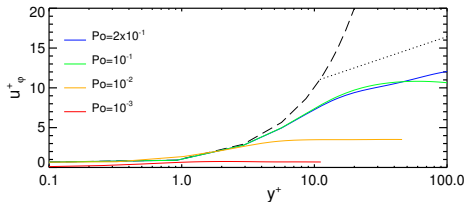
- Profiles depend by $Ek = \nu / (R^2 \Omega_c)$
- Peaks around $Po \approx 0.125$ for small Ek
- no cross of B-mode instability and no fully turbulent boundary layer (on the endwalls)

Other quantities

$$y^+ = \frac{zu_\tau}{\nu} \quad u_\tau = \sqrt{\partial u_\phi / \partial z|_{wall}}$$



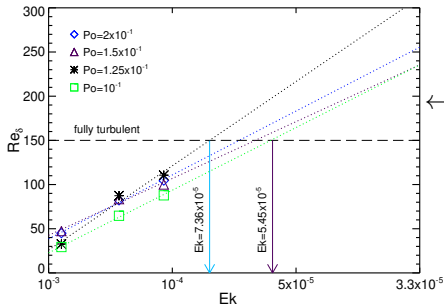
$$Ro_f = \frac{u_\tau^2}{f'\nu} \quad f' = 2\Omega_c + U_g/r + dU_g/dr$$



$$u^{rms} = \frac{\sqrt{\langle u_j' u_j' \rangle}}{u_\tau}$$

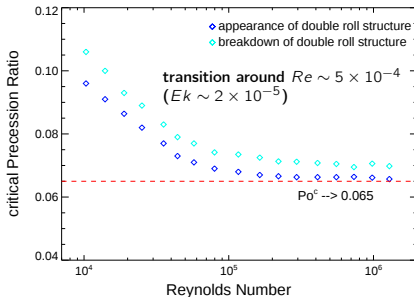
- NO turbulent phenomena.
- Flow follow the viscous sub-layer but not achievement of log-law
- fluctuations strenght small

Connection with previous experiments



← Extrapolation for turbulent Ekman layer

Relation with experiment



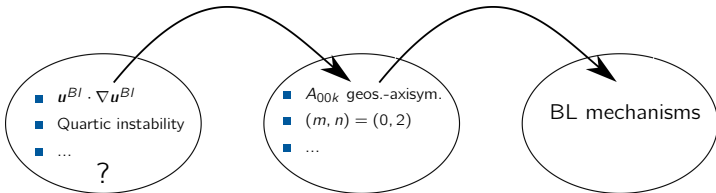
- critical precession ratio for transition to turbulence in dependence of Re .
- UDV measurements: up to $Re \approx 2 \times 10^6$ show **asymptotic behaviour** for $Re \geq 10^5$. $Po \rightarrow 0.065$
- width of instable region seems to be largely **independent** of Re

Summary

- characterization of the bulk flow structure through A_{mnk}
- Geostrophic-axisymmetric flow dominance for strong precession
- characterization of Ekman (endwalls) layer in precessing cylinder $f(Ek, Po)$
- Appearance of instabilities but no fully turbulent Ekman layer (in the range of our simulations)
- Reasonable prediction of turbulence onset (consistency with asymptotic profile shown by experiments)

Limitations of this study:

- results valid for large precession angle α and not extreme Ek .
- turbulent study requires finer mesh (e.g 10 grid points $0 < y^+ < 10$).
- no insights to decouple bulk mechanism and Ekman mechanism: what causes what?



Possible future study

Experiments on the Ekman layer (PIV ??).

Theoretical connection with bulk flows phenomena

Deeper analysis on the connection of $(m, n) = (0, 2)$ with the Ekman suction (average thickness could increase around $Po \approx 0.10$)

Ekman (endwall) layers phenomena in MHD context

The end

Thank you for your attention

Appendix

