

Evidence for the Scissors Mode in ^{160}Tb from the Two-Step Gamma Cascades measurement



J. Kroll¹, M. Krtička¹, F. Bečvář¹ and I. Tomandl²

¹ Faculty of Mathematics and Physics, Charles University, Prague

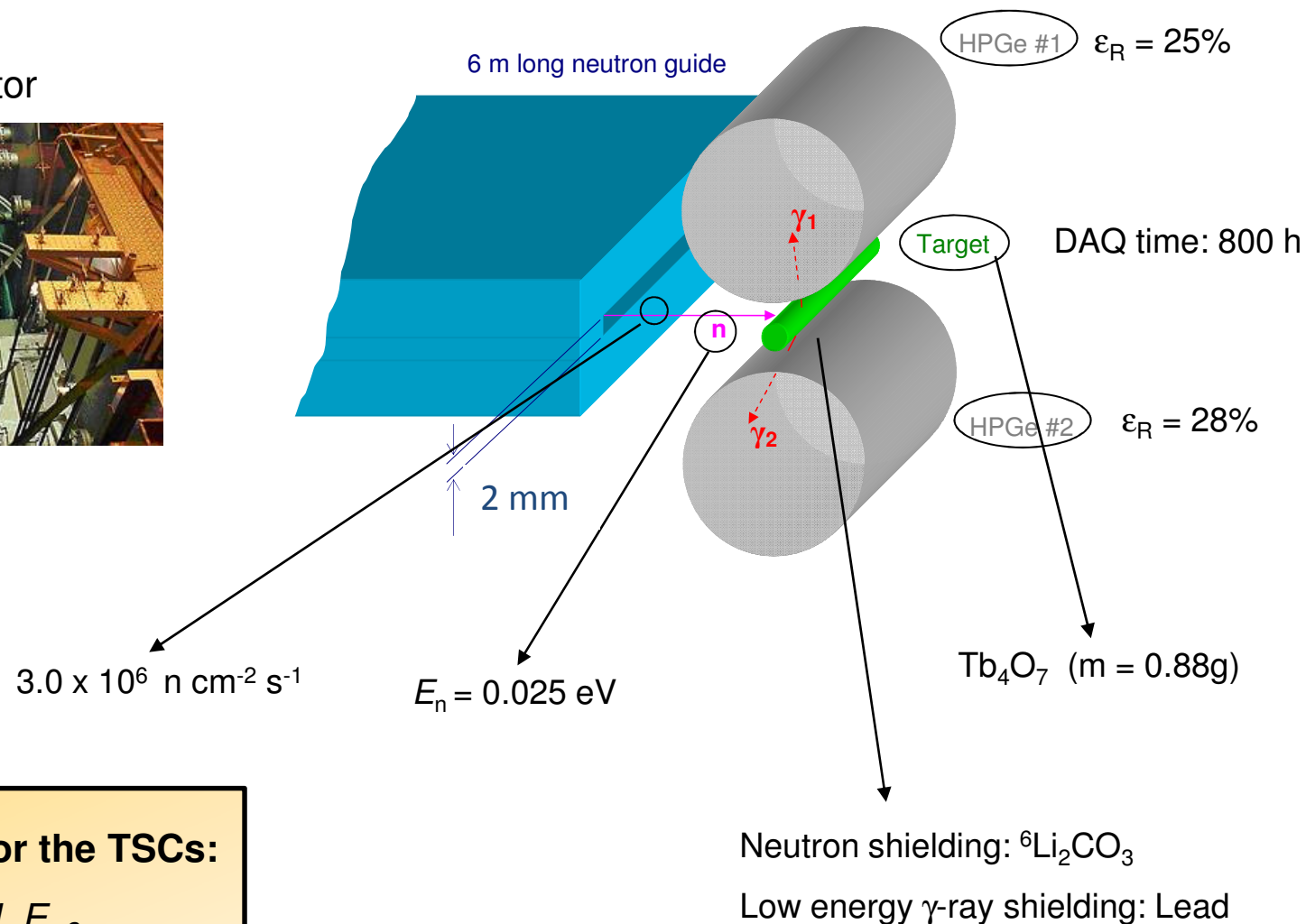
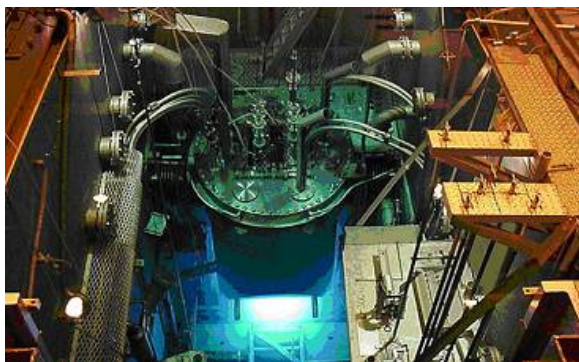
² Nuclear Physics Institute of the Czech Academy of Science, Řež

Outline

- **Experimental setup for the TSCs measurement**
- **Simulations of gamma decay**
- **Main results**
- **Conclusions**

Experimental setup for the TSCs measurement

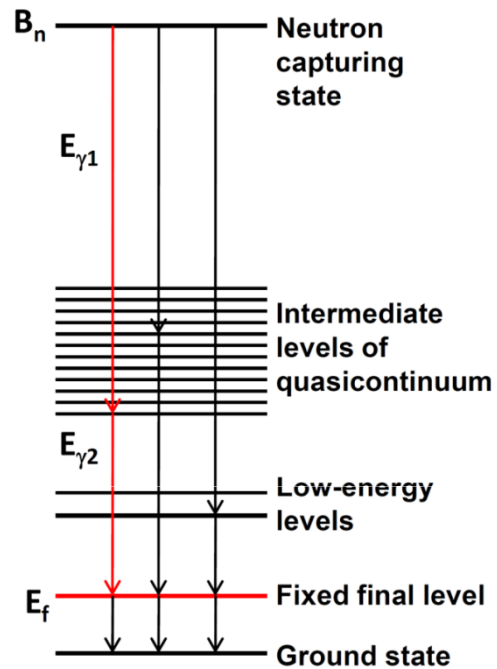
LWR-15 reactor



DAQ conditions for the TSCs:

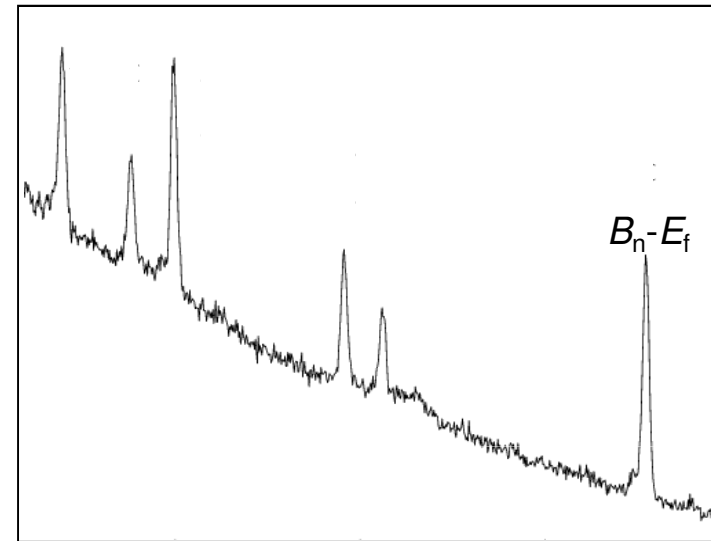
- energies $E_{\gamma 1}$ and $E_{\gamma 2}$
- time difference

Data processing



From information about $E_{\gamma 1}$ and $E_{\gamma 2}$ and **detection time difference**, one can retrieve virtually background-free TSC spectra.

Spectrum of energy sums



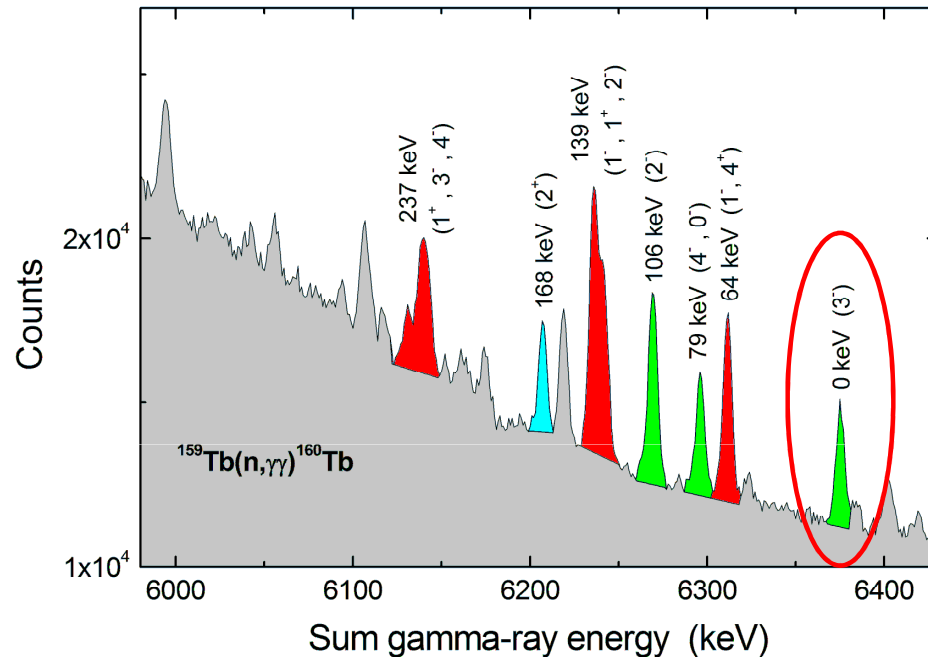
Energy sum $E_{\gamma 1} + E_{\gamma 2}$



List-mode data

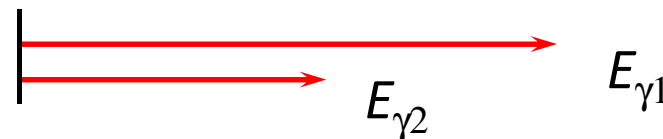
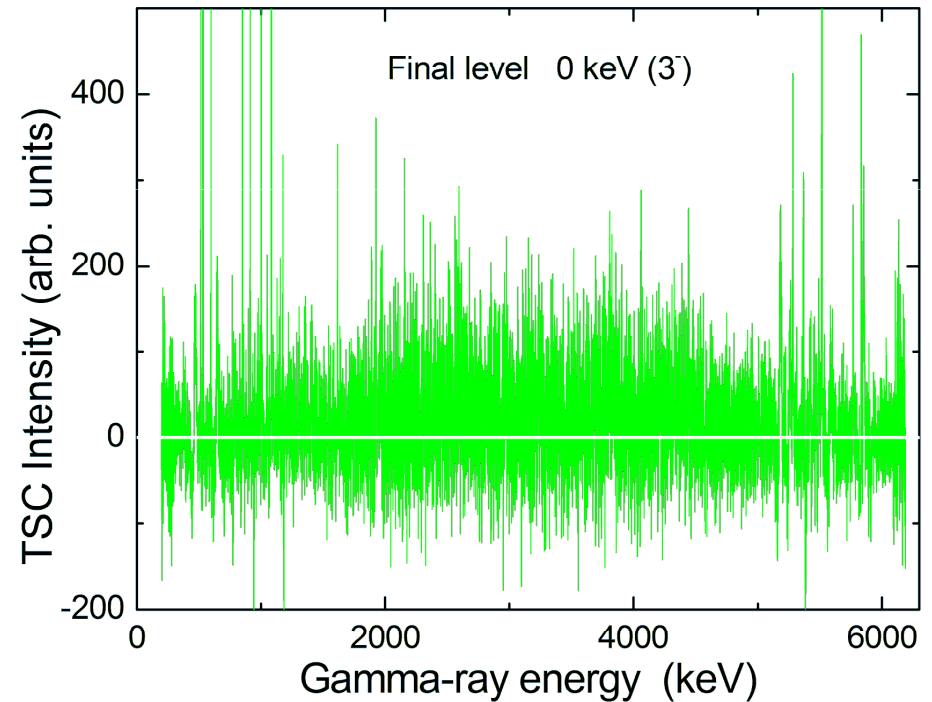
Experimental setup for the TSCs measurement (3)

Spectrum of energy sums



$$S_n = 6.375 \text{ MeV}$$

Spectrum of Two-Step Cascades



Simulations of gamma decay – DICEBOX (1)

1. Below a **critical energy** E_{crit} the energies E , spins J , parities π and the decay properties of all levels are taken from known data
2. Above the critical energy E_{crit} the energies E , spins J and parities π of levels are obtained by random discretization of an *a priori* known level density

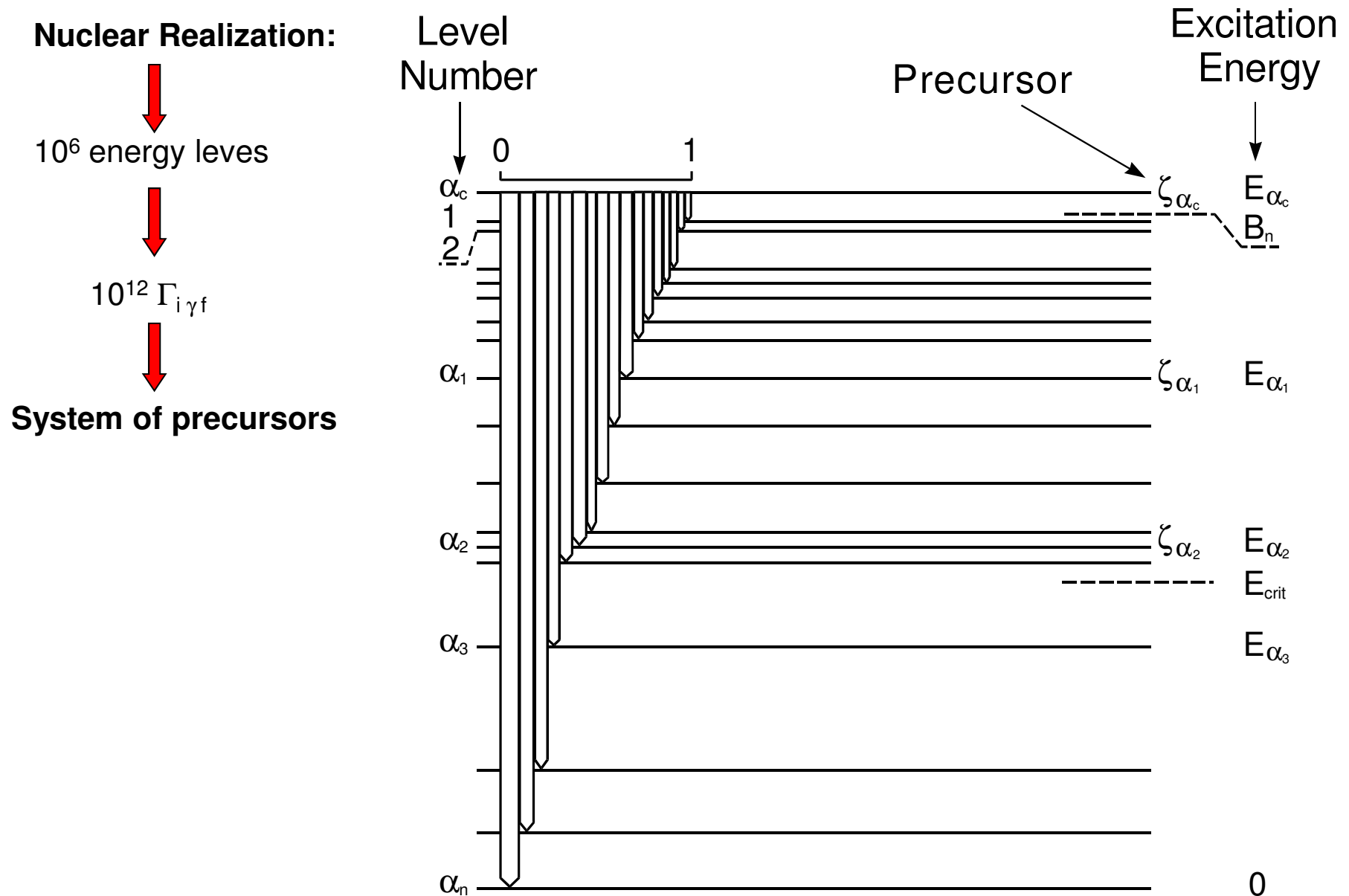
$$\rho(E, J, \pi)$$

3. **Partial radiation widths** $\Gamma_{i\gamma f}$ for transitions between initial (i) and final (f) levels are generated according to the formula:

$$\Gamma_{i\gamma f} = \sum_{XJ} y_{ifXJ}^2 (E_i - E_f)^{2J+1} \frac{f^{(XJ)}(E_i - E_f)}{\rho(E_i, J_i, \pi_i)}$$

4. Partial radiation widths $\Gamma_{i\gamma f}$ for different initial and/or final levels are statistically independent.

Simulations of gamma decay – DICEBOX (2)



Simulations of gamma decay – DICEBOX (3)

1. Below a **critical energy** E_{crit} the energies E , spins J , parities π and the decay properties of all levels are taken from known data
2. Above the critical energy E_{crit} the energies E , spins J and parities π of levels are obtained by random discretization of an *a priori* known level density

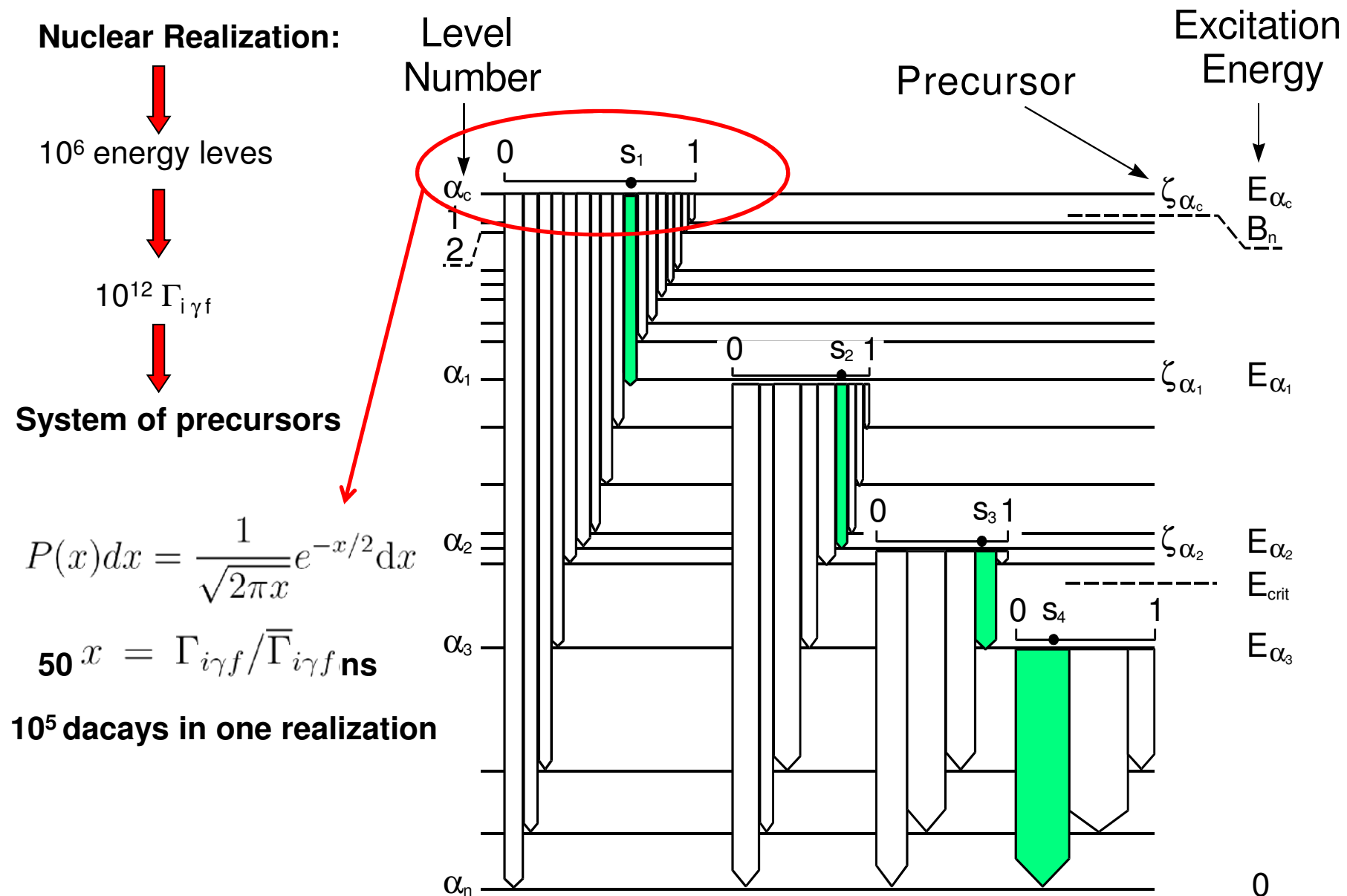
$$\rho(E, J, \pi)$$

3. **Partial radiation widths** $\Gamma_{i\gamma f}$ for transitions between initial (i) and final (f) levels are generated according to the formula:

$$\Gamma_{i\gamma f} = \sum_{XJ} \overset{\text{P-T fluctuations}}{y_{ifXJ}^2} (E_i - E_f)^{2J+1} \frac{f^{(XJ)}(E_i - E_f)}{\rho(E_i, J_i, \pi_i)}$$

4. Partial radiation widths $\Gamma_{i\gamma f}$ for different initial and/or final levels are statistically independent.

Simulations of gamma decay – DICEBOX (4)



Simulations of gamma decay – LD (1)

1. Below a **critical energy** E_{crit} the energies E , spins J , parities π and the decay properties of all levels are taken from known data
2. Above the critical energy E_{crit} the energies E , spins J and parities π of levels are obtained by random discretization of an *a priori* known level density

$$\rho(E, J, \pi)$$

Level density

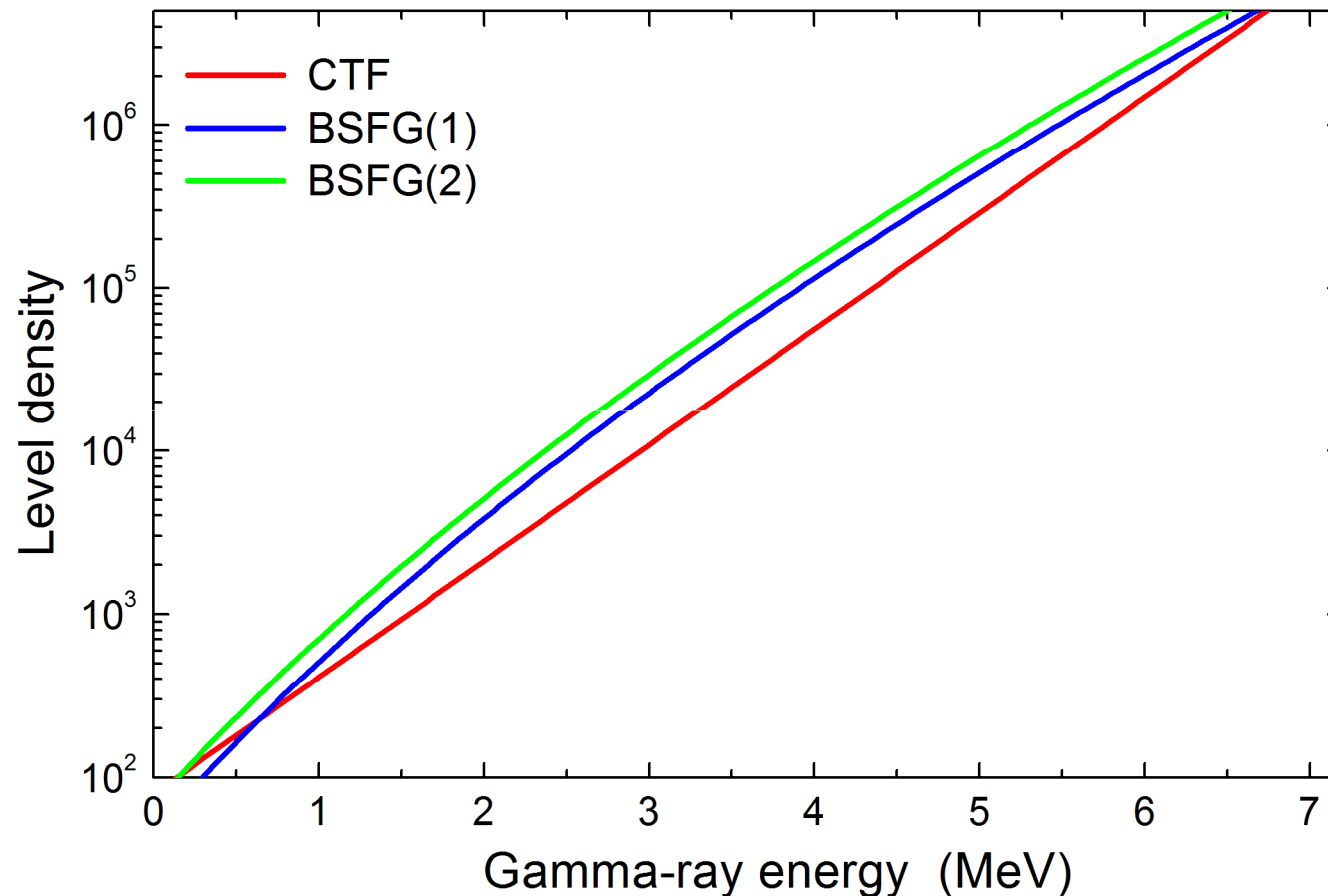
3. **Partial radiation widths** $\Gamma_{i\gamma f}$ for transitions between initial (i) and final (f) levels are generated according to the formula:

$$\Gamma_{i\gamma f} = \sum_{XJ} \overset{\text{P-T fluctuations}}{y_{ifXJ}^2} (E_i - E_f)^{2J+1} \frac{f^{(XJ)}(E_i - E_f)}{\rho(E_i, J_i, \pi_i)}$$

4. Partial radiation widths $\Gamma_{i\gamma f}$ for different initial and/or final levels are statistically independent.

Simulations of gamma decay – LD (2)

Tested models of LD (total)



- (1) T. von Egidy, H.H. Schmidt and A.N. Behkami, Nucl. Phys., **A481** (1988) 189
(2) T. von Egidy and D. Bucurescu, Phys. Rev. **C72**, (2005) 044311

Simulations of gamma decay – PSFs (1)

1. Below a **critical energy** E_{crit} the energies E , spins J , parities π and the decay properties of all levels are taken from known data
2. Above the critical energy E_{crit} the energies E , spins J and parities π of levels are obtained by random discretization of an *a priori* known level density

$$\rho(E, J, \pi)$$

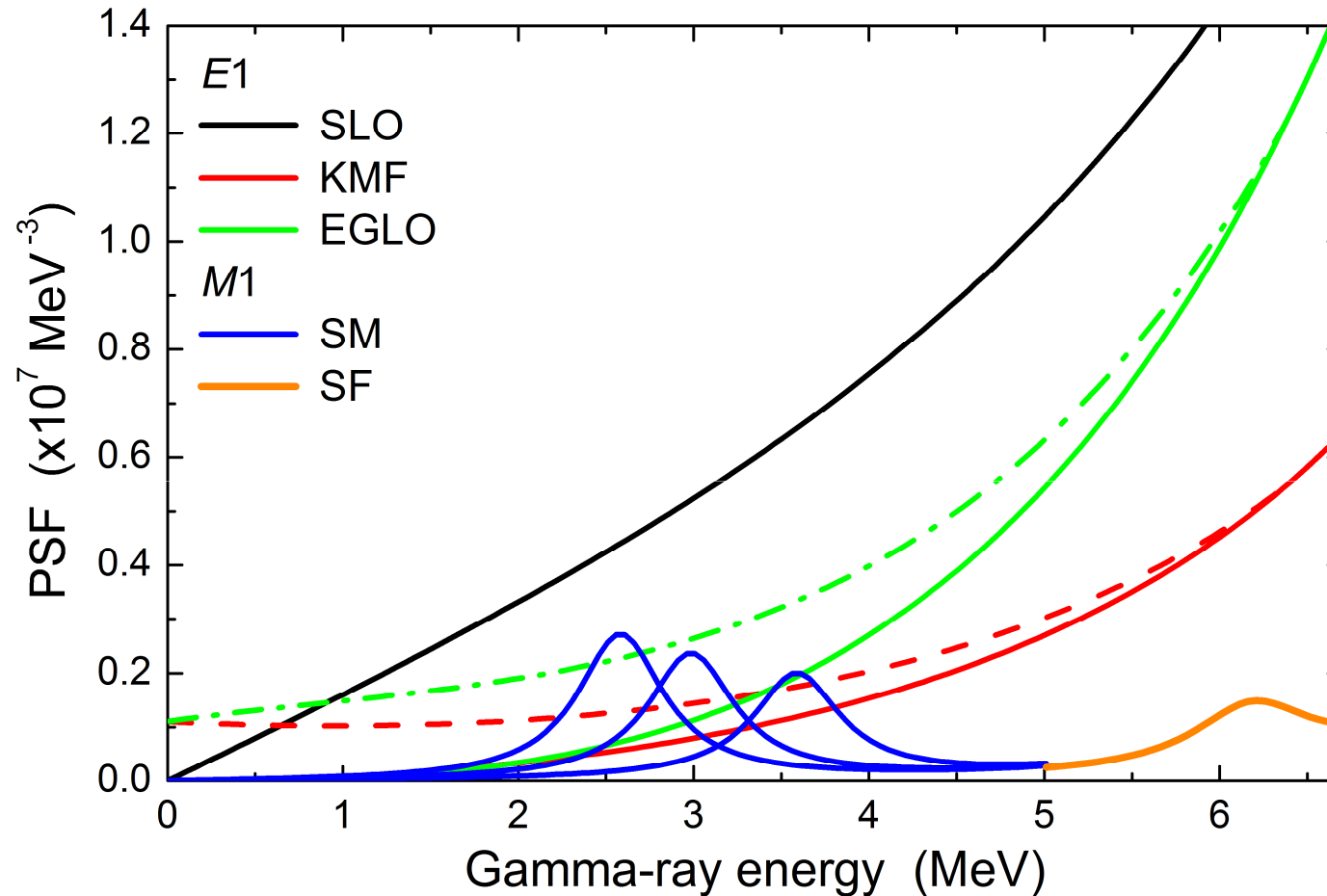
Level density

3. **Partial radiation widths** $\Gamma_{i\gamma f}$ for transitions between initial (i) and final (f) levels are generated according to the formula:

$$\Gamma_{i\gamma f} = \sum_{XJ} \underbrace{y_{if}^2}_{\text{P-T fluctuations}} (E_i - E_f)^{2J+1} \underbrace{\frac{f^{(XJ)}(E_i - E_f)}{\rho(E_i, J_i, \pi_i)}}_{\text{PSFs}}$$

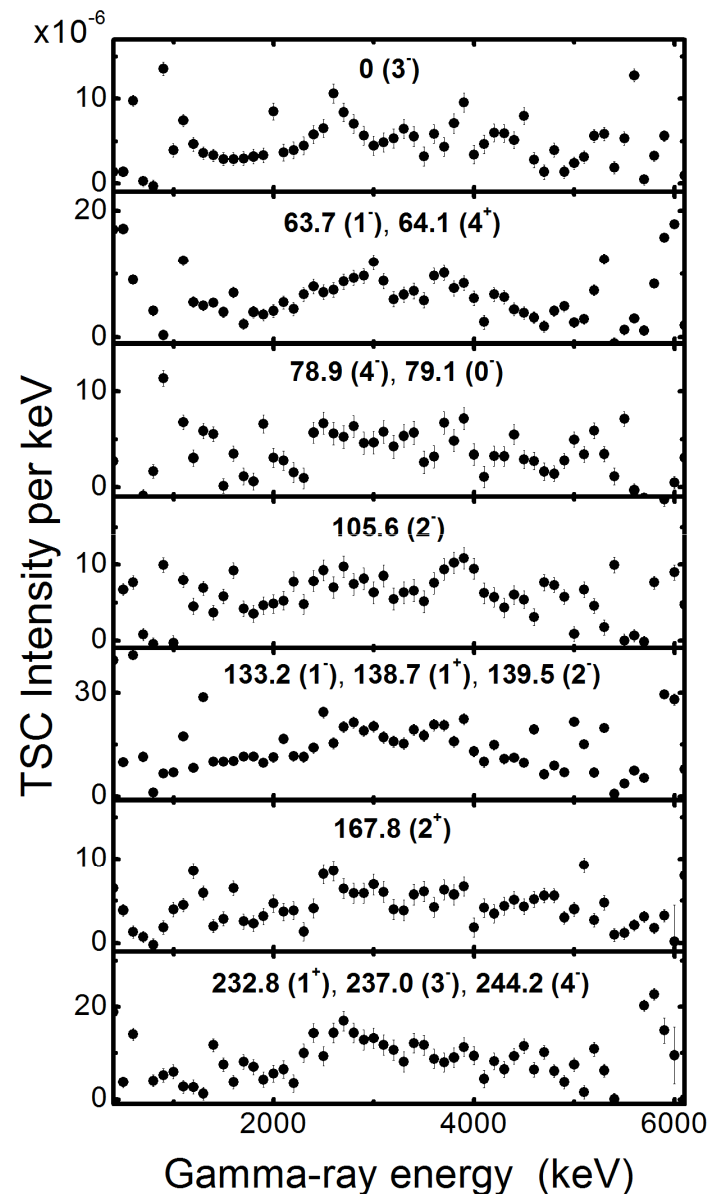
4. Partial radiation widths $\Gamma_{i\gamma f}$ for different initial and/or final levels are statistically independent.

Simulations of gamma decay – PSFs (2)



The energy of the SM is 2.6, 3.0 and 3.6 MeV, damping width is 0.6 MeV and the total $\Sigma B(M1)_{\uparrow} \approx 5 \mu_N^2$.

Results (1)

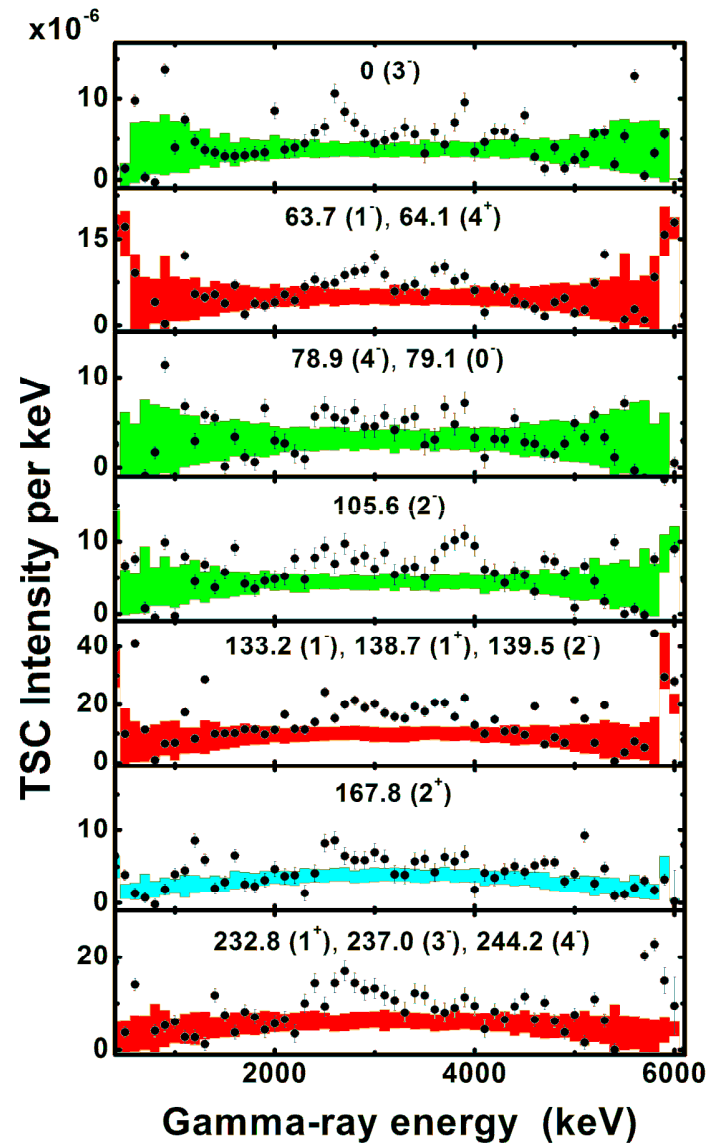


Experimental binned TSC spectra.

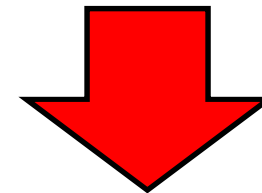
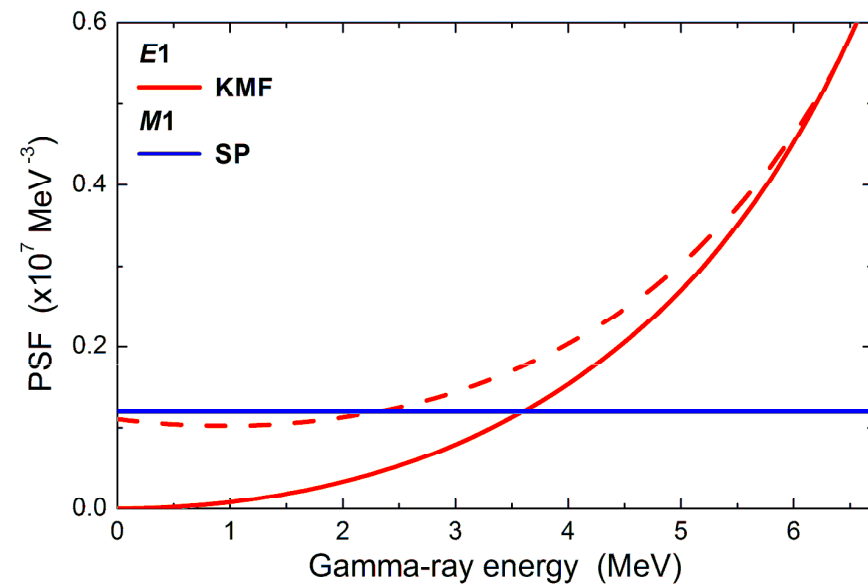
The bin width is 100 keV.

Individual experimental TSC spectra display distinct **resonance-like structures at 2.6 MeV and 3.6 MeV.**

Results (2)

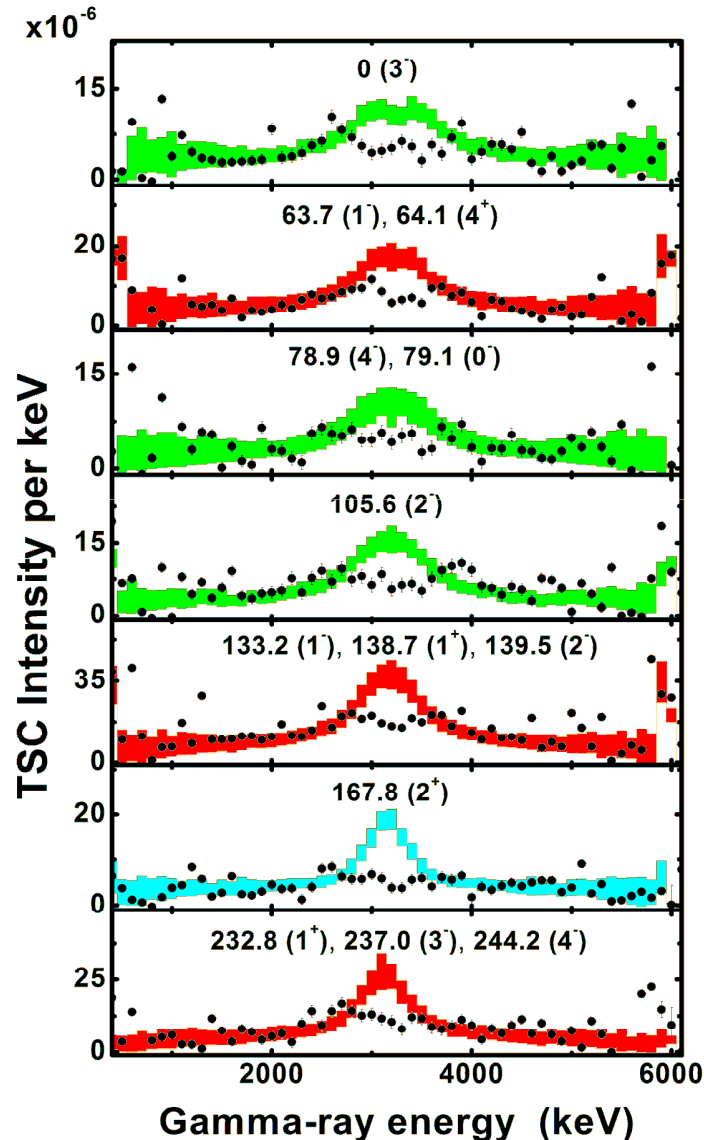


No resonance structure in *M1*.

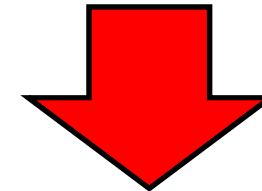
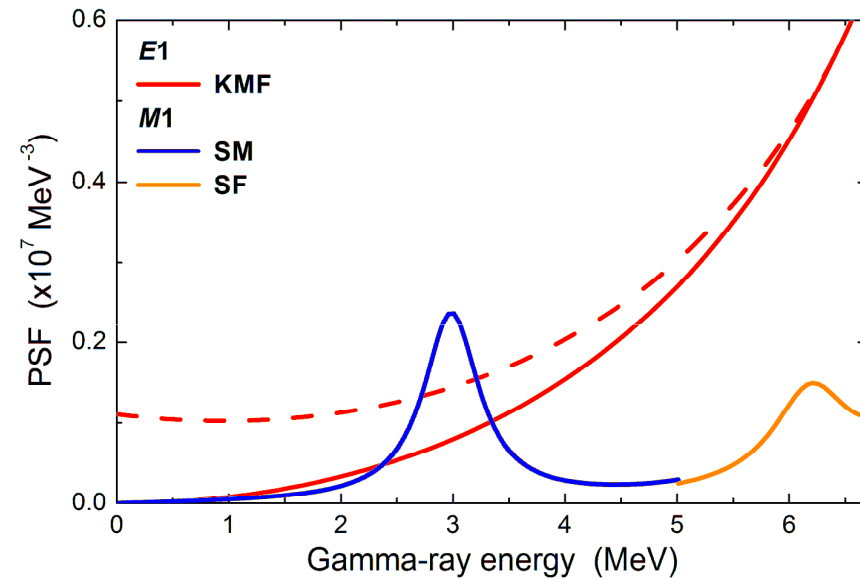


Evident disagreement !

Results (3)

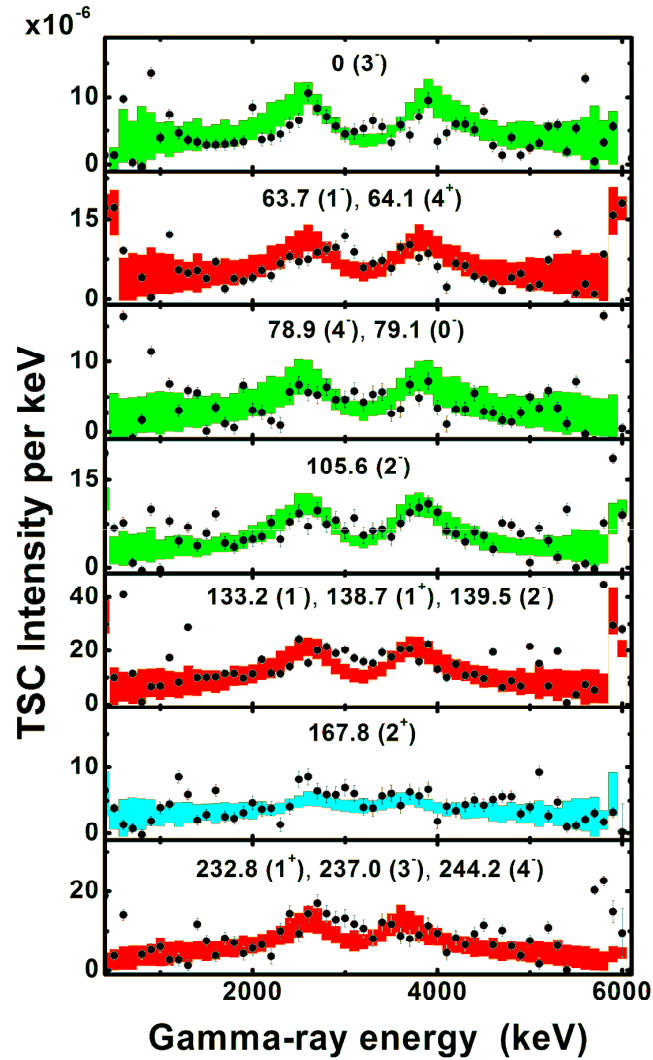


SM used for $M1$ PSF on the energy
 $E_{SM} = 3.0$ MeV, $\Gamma_{SM} = 0.6$ MeV and $\sigma_{SM} = 1.0$ mb.

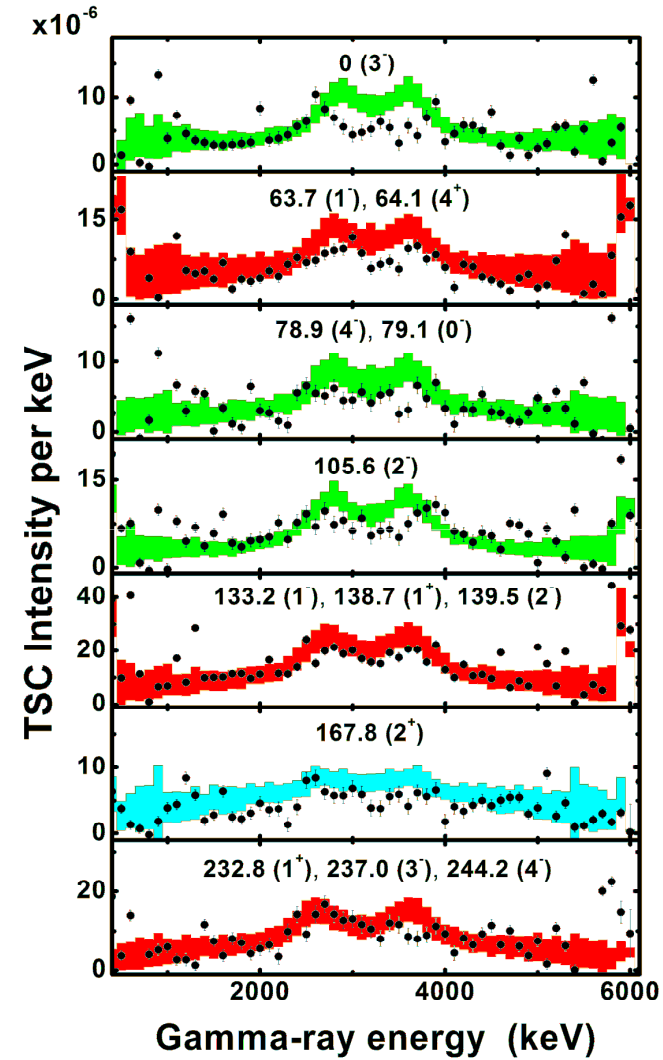


The double-humped structure still
not reproduced !

Results (4)

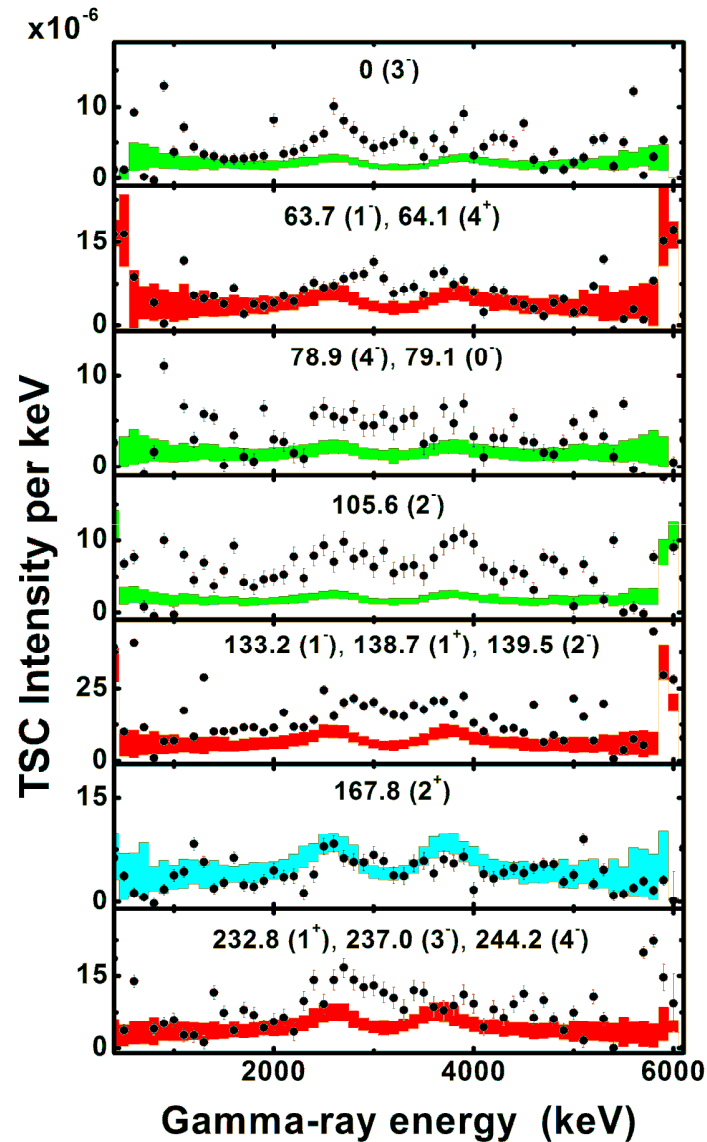


SM in *M1* PSF: $E_{\text{SM}} = 2.6 \text{ MeV}$,
 $\Gamma_{\text{SM}} = 0.6 \text{ MeV}$ and $\sigma_{\text{SM}} = 1.0 \text{ MeV}$

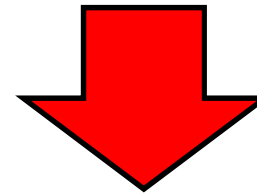


SM in *M1* PSF: $E_{\text{SM}} = 3.6 \text{ MeV}$,
 $\Gamma_{\text{SM}} = 0.6 \text{ MeV}$ and $\sigma_{\text{SM}} = 1.0 \text{ MeV}$

Results (5)



Lorentz-shape resonance structure with $E_R = 2.6$ MeV, $\Gamma_R = 0.6$ MeV and $\sigma_{SM} = 1.0$ mb is postulated in $E1$ PSF. For $M1$ it is supposed non-resonance shape.



The observed structures cannot be reproduced by the presence of local maximum of $E1$ PSF neither at 2.6 nor at 3.6 MeV !

Results (6)

Reaction	(γ, γ')	$(^3\text{He}, x \gamma)$	$(^3\text{He}, x \gamma)$	$(n, \gamma\gamma)$	$(n, \gamma\gamma)$
Nuclei	e-e ¹	^{160,161,162} Dy ²	^{163,164} Dy ³	¹⁶³ Dy ⁴	¹⁶⁰ Tb
E_{SM} (MeV)	~3.0	2.6 – 2.8	~2.8	~3.0	2.6 – 2.8 (3.6 – 3.8)
Γ_{SM} (MeV)	-----	1.2 – 1.6	0.8 - 0.9	0.6 (0.5 – 0.7)	0.4 – 0.9
σ_{SM} (mb)	-----	0.3 – 0.4	0.5 – 0.7	0.9 (0.8 – 1.0)	0.4 – 0.9
$\Sigma B(M1)$ (μ_N^2)	~3	~7	5 - 8	~6	~6 (3 – 9)

¹ Kneissl, Pitz and Zilges, Prog. Part. Nucl. Phys. 37 (1996) 349

² Guttormsen et al. Phys. Rev. C 68, 064306 (2003)

³ Nyhus et al., Phys. Rev. C 81, 024325 (2010)

⁴ Krticka, et al., Phys. Rev. Lett. 92, 172501 (2004)

Conclusions

- Double-humped structure in experimental TSC spectra clearly indicates that the $M1$ SM plays an important role in gamma deexcitation of ^{160}Tb .
- The $E1$ origin of the resonance-like structures in the TSC spectra is unambiguously excluded.
- The energy of the SM is very likely $E_{\text{SM}} = 2.6 \pm 0.1$ MeV but the value $E_{\text{SM}} = 3.6 \pm 0.1$ MeV cannot be completely excluded. The damping width of the SM has to be $\Gamma_{\text{SM}} = 0.5 - 0.9$ MeV. The best agreement is obtained with the strength of the mode $\Sigma B(M1)\uparrow = 6 \pm 1 \mu_N^2$.

Thank you for your attention !