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Prof. Dr. Thorsten Stumpf Dr. Harald Foerstendorf Dr. Frank Bok Dr. Anke Richter



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Cover pictures

How the Institute of Resource Ecology was kept running from home in 2020 due to pandemic restrictions — Impressions

Photos provided by: D. Abbasova, S. Bachmann, R. Berndt, S. Beutner, F. Bok, M. Demnitz, E. Dräger, B. Drobot, C. Fischer, H. Foerstendorf, C. Franzen, R. Gericke, A. Günther, S. Gürtler, J. Jessat, C. Kielar, E. Krawczyk-Bärsch, K. Kurde, H. Moll, K. Müller, K. Nebe, M. Patzschke, J. Raff, A. Richter, Q. Roode-Gutzmer, R. Steudtner, M. Stockmann, T. Stumpf, W. Webersinke.

Preface

The INSTITUTE OF RESOURCE ECOLOGY (IRE) IS ONE OF the eight institutes of the Helmholtz-Zentrum Dresden-Rossendorf (HZDR). Our research activities are mainly integrated into the program "Nuclear Waste Management, Safety and Radiation Research (NUSAFE)" of the Helmholtz Association (HGF) and focused on the topics "Safety of Nuclear Waste Disposal" and "Safety Research for Nuclear Reactors". The program NUSAFE, and therefore all work which is done at IRE, belong to the research field "Energy" of the HGF.

Our research objective is the protection of humans and the environment from the effects of radioactive radiation. To this end, we create a molecular process understanding using the most modern methods of spectroscopy, microscopy, diffraction, numerical simulation and theoretical chemistry. This takes place in a unique research environment at the HZDR. Our interdisciplinary approach combines radiochemistry, geosciences, and biosciences as well as materials science and reactor physics. This generates knowledge with application in particular for reactor and repository safety as well as in radioecology.

In 2020, the Corona pandemic made our work much more difficult. The laboratories were closed for several weeks and nearly no direct scientific exchange with collaborators, partners, and friends was possible. To my surprise, this had no influence on our scientific outcome. In 2020, 108 original papers were published in peer-reviewed international scientific journals with an average impact factor of 5.77. The extraordinary broadness of research topics and activities is illustrated below by some selected highlights.

Data of a start-up experiment at the Chinese Experimental Fast Reactor (CEFR) provided within an IAEA coordinated research project have been used for post-test calculations by means of the continuous energy Monte Carlo code SERPENT2. Beside the excellent agreement between calculation and measurement data, this activity demonstrates the feasibility to routinely conduct time-dependent simulations, an area which until recently was limited to deterministic methods. In general, due to the more accurate interaction physics combined with practically unlimited geometric complexity, Monte Carlo codes have a strong potential to produce high-fidelity reference solutions for the verification of deterministic codes, not just in static neutronic calculations as they are widely used so far, but also in time-dependent analyses (\rightarrow p.81).

We investigated relations between calculated primary displacement damage, the irradiated nanostructure, and the nanoindentation response of ion-irridiated Fe-Cr alloys. Cross-sectional scanning transmission electron microscopy indicated the appearance of irradiation-induced dislocation loops and revealed their number density as a function of the distance from the irradiated surface. A model aimed at integrating the calculated local hardening contributions over the indentation plastic zone was calibrated using the measured nanoindentation response. The model allowed the effect of the implanted interstitials to be separated from the total hardening. The result is an important step towards improved transferability of ion irradiations to the case of much more expensive neutron irradiations, where implanted interstitials do not appear (\rightarrow p.73). Within the framework of both the GRaZ joint project (BMWi) and the EURAD CORI project (EURATOM) investigations on the U(VI)-Citrate system were performed. Remaining unnoticed for over 70 years of intensive research until recently, complexes of a so-called sandwich structure were characterized combining NMR spectroscopy, especially of rarely applied ¹⁷O, ²³Na, and ¹³⁹La nuclei, and DFT calculation. As these species remain soluble in pH neutral media even at high U concentrations, the results are vital for understanding uranium's versatile coordination chemistry as well as environmental migration behavior (\rightarrow p. 16).

Molecular diffusion is the prevailing mechanism for radionuclide transport within clay-rich rocks. Current reactive transport modeling (RTM) applies a simplified parametrization of the effective diffusivity for simulating migration in host rocks such as the Opalinus clay. This prevents to model the heterogeneous diffusion behavior currently identified by using positron emission tomography (GeoPET) measurements, based on Na-22 as a conservative tracer for long-term transport experiments (\rightarrow p.43). To improve the predictive capabilities of RTM, we developed an upscaling workflow that integrates nm- and µm-scale simulations based on synthetic multiscale digital rocks for the prediction of the effective diffusivity of radionuclides (\rightarrow p. 40). Both the spatial variability of the pore network and the compositional heterogeneities at the pore scale are now available for studying heterogeneous diffusion patterns, in contrast to commonly assumed homogeneous behavior.

The mechanism of isotope-dependent hydrogen permeability of graphene membranes has been controversially discussed during the past few years. We confirm in a theoretical study (\rightarrow p. 46) that ideal graphene is quasi-impermeable to protons, yet the most common defect in sp2 carbons, the topological Stone–Wales defect, has an H⁺-over-D⁺ selectivity explaining all experimental results on graphene membranes that are available to date, thus challenging local hydrogenation as the competing explanation.

The investigations of basic properties of actinides were expanded to include a series of ligands with soft N-donor functions. (\rightarrow p.21). Particularly noteworthy is a work that combines theory and experiment in order to elucidate the binding relationships in actinide complexes with the PEBA ligand. As expected, the proportion of covalent bonds increases significantly with the occurrence of f-electrons in Pa(IV). Almost identical covalencies are found for U(IV), Np(IV), and Pu(IV), but the origin of this covalency shifts from a strong d-character for uranium to a stronger f-character for plutonium (\rightarrow p.22).

Proteomics studies were conducted this year to investigate the interaction of uranium with *Nicotiana tabacum* (tobacco) cells to understand how radionuclides migrate in the biosphere and through the food chain on a molecular level. First results elucidated an increase in the glutathione (GSH) metabolism of these cells in the presence of uranium. This confirmed the initiation of GSH-mediated sequestration upon uranium exposure to plant cells as part of the cellular response to the heavy metal stress (\rightarrow p.55).

The European Synchrotron Radiation Facility opened its new 4th generation electron storage ring for user operation in 2020, and with it our newly constructed Rossendorf Beamline (ROBL). Three experimental stations, the 5-crystal spectrometer for high-energy-resolution X-ray absorption and emission spectroscopy, the Pilatus 2M diffractometer for single-crystal and *in situ* diffraction, and the EXAFS station with the world's first multielement Ge detector with CUBE amplifiers and unprecedented detection limit, have accomplished the first successful user experiments; the 6-circle Huber diffractometer for surface and high-resolution powder diffraction has been delivered and is currently under commissioning. The combination of different synchrotronbased X-ray techniques, which are now available at ROBL, has proven vital to elucidate the plutonium oxidation state and structure of PuO₂ nanoparticles, which was hitherto heavily debated in the scientific literature (\rightarrow p.13).

Beside these highlights, we obtained many other new scientific results in the past year, which are presented in this annual report.

In the last year, more than 170 scientists, technicians, and students working on their Ph.D., diploma, master, or bachelor theses, were employed at the Institute of Resource Ecology. Thereof, 39 Ph.D. students worked at the institute in 2020.

As an important event in 2020, the international workshop entitled "How to integrate geochemistry at affordable costs into reactive transport for large-scale systems" was organized by the institute on February 5–7. Besides many interesting scientific talks and poster presentations, four excellent keynote talks were presented. In total, 40 scientists from 28 research institutes and eight countries extensively discussed advanced numerical concepts in geochemistry (\rightarrow p.103).

Furthermore, the annual meeting of HGF and BMBF funded interdisciplinary project iCross (Integrity of nuclear waste repository systems - Cross-scale system understanding and analysis) was organized fully virtual from November 25–26. More than 70 participants from HZDR, KIT, FZJ, UFZ and GFZ and invited guests from ministries, federal agencies and universities discussed in detail studies performed on laboratory scale and in the underground research laboratory Mont Terri (CH) as well as modern modeling approaches.

In 2020, funding to implement "Hover" – a decentralized research infrastructure was finally approved. Located at KIT, FZJ, and HZDR, it is dedicated to enable new and to support existing advanced scientific investigations and technical de-

velopments in the context of the phase out of nuclear energy within the German 'Energiewende' project. At HZDR, a radiotechnicum will permit comprehensive, host rockindependent investigations as necessary due to the longtime scales expected until a final decision on the siting of a nuclear waste repository in Germany will be made. The radiotechnicum will offer the facilities for experimental investigations of bio/geo/radiochemical processes on larger scales up to the mesoscale. Furthermore, Hover will be the seed crystal for a new controlled area building, which will replace the old radiochemical laboratories.

In 2020, the scientific director as well as the administrative director at HZDR changed. Prof. Dr. Sebastian M. Schmidt took over the position of the scientific director from Prof. Sauerbrey, who retired. We thank Prof. Sauerbrey for his strong support and wish him all the best for the new phase of life. Furthermore, at the end of the year, Dr. Diana Stiller took on the position of the administrative director. We are looking forward to a successful and trustful collaboration in the future.

In retrospect of a challenging but at the end successful year, I would like to thank the (unfortunately only a few) visitors, German and international ones, for their interest in our research and for their participation in the institute seminars. Although it was very limited in 2020, we will continue to strongly encourage the collaborations and visits by scientists in the future. Special thanks are due to the executive board of the HZDR, the Ministry of Science and Arts of the Free State Saxony (SMWK), the Federal Ministry of Education and Research (BMBF), the Federal Ministry of Economic Affairs and Energy (BMWi), the Deutsche Forschungsgemeinschaft (DFG), the European Commission, and other organizations for their support.

Thoush Stanget

Prof. Dr. Thorsten Stumpf Director of the Institute of Resource Ecology

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The Chemistry of LONG-LIVED RADIONUCLIDES

All roads lead to PuO₂ nanoparticles

E. Gerber, A. Y. Romanchuk,¹ I. Pidchenko, L. Amidani, A. Rossberg, C. Hennig, G. B. M. Vaughan,² A. Trigub,³ T. Egorova,¹ S. Bauters, T. V. Plakhova,³ M. O. J. Y. Hunault,⁴ S. Weiss, S. M. Butorin,⁵ A. C. Scheinost, S. N. Kalmykov,^{1,3} K. O. Kvashnina ¹Lomonosov Moscow State University, Department of Chemistry, Moscow, Russia; ²ESRF – The European Synchrotron, Grenoble, France; ³National Research Centre "Kurchatov Institute", Moscow, Russia; ⁴Synchrotron SOLEIL, France; ⁵Molecular and Condensed Matter Physics, Uppsala University, Uppsala, Sweden

Here, we report a full characterization of PuO_2 nanoparticles (NPs) at the atomic level and probe their local and electronic structure by a variety of synchrotron methods. Well dispersed, crystalline NPs with a size of ~2.5 nm in diameter are always formed in spite of varying synthesis conditions. Identical crystal structures and the presence of only the Pu(IV) oxidation state in all NPs were confirmed.^[1]

Plutonium can migrate in colloidal form in the subsurface environment across several kilometers.^[2,3] Recent spectroscopic and microscopic investigations showed that so called "colloidal Pu(IV) polymers" are in fact aggregates of PuO₂ nanoparticles,^[4,5] however, the exact stoichiometry and structure of such nanoparticles remain still questionable as well as the purity of the tetravalent oxidation state. In this study, the synthesized PuO₂ NPs are characterized by a variety of experimental methods: High-Resolution Transmission Electron Microscopy (HRTEM), Selected-Area Electron Diffraction (SAED), Extended X-Ray Absorption Fine Structure (EXAFS), X-ray Absorption Near Edge Structure (XANES) in High Energy Resolution Fluorescence Detection (HERFD) mode, High Energy X-ray Scattering (HEXS) and X-ray Diffraction (XRD).

EXPERIMENTAL. The nanoparticles were synthesized from precursors with different oxidation states of plutonium (III, IV, and V) under various environmentally and waste storage relevant conditions (pH8 and pH>10). The Pu L₃ HERFD and EXAFS experiments as well as XRD measurements were performed at The Rossendorf Beamline at ESRF.^[6] The HERFD spectra (except one) at the Pu M₄ edge were collected at Beamline ID26 of ESRF. High energy X-ray scattering data were collected at room temperature at the ID15A beamline of the ESRF.

RESULTS. It was found that similar non-amorphous NPs are formed, regardless of different precursors and pH conditions. The crystalline structure of the NPs is similar to that of bulk PuO2. The particle size varies in the range of 2.3-3.2 nm. Pu oxidation state in the PuO₂ NPs turns out to be exclusively Pu(IV), as proven here by HERFD methods at L₃ and M₄ edges in Fig.1. The EXAFS spectra (Fig.2) show effects of increased static disorder and reduced coordination numbers, which are clear influences of the nanometer-size of the particles, without the need to invoke different Pu oxidation states or a substantially different structure. Finally, the agreement among the results from the EXAFS shell fit, MC EXAFS simulations (Fig. 2), and from HEXS all unequivocally indicate that only one Pu(IV)-O interaction with a rather symmetric distance distribution (similar to the PuO₂ reference) is present for all six PuO₂ NPs samples.



^[2] Kersting, A. et al. (1999) Nature **397**, 56–59.



Fig. 1: Pu M₄ and L₃ HERFD experimental data. Data recorded for six PuO₂ NPs samples and compared with a PuO₂ reference and theoretical calculations, using AIM approximation (left) and FEFF code (right). The density of states is plotted at the bottom for clarity.



Fig. 2: Pu L₃-EXAFS spectra χ(R) fit results, Fourier transform (FT) magnitude of experimental EXAFS data (black) and a one Pu-O, one Pu-Pu shell fit (red), Monte-Carlo simulations fit (green). The area of the shell fit in R-space is marked with two dashed lines, MC simulation output: radial particle distribution function (n(r)) for Pu-O and Pu-Pu and 3D structural refinement based on bulk PuO₂ starting structure.

^[3] Novikov, A. et al. (206) Science 314, 638-642.

^[4] Powell, B. et al. (2011) Environ. Sci. Technol. 45, 2698–2703.

^[5] Romanchuk, A. et al. (2018) Dalton Trans. 47, 11239–11244.

^[6] Kvashnina, K. et al. (2016) J. Synchrotron Radiat. 23, 836–841.

Accurate structural characterization of ThO₂ nanoparticles with HEXS and HERFD-XANES

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An accurate characterization of the structure of nanoparticles (NPs) is often very hard to achieve, especially for non-homogeneous samples and NPs < 3 nm in diameter. By combining High-Energy X-ray Scattering (HEXS) and High-Energy-Resolution Fluorescence-Detected X-ray Absorption Near-Edge Structure (HERFD-XANES) spectroscopy, we characterized both the short- and mediumrange structure of ThO₂ NPs obtained by chemical precipitation. The initial stage of NPs formation shows a high concentration of Th units similar to Th hexamer clusters. With annealing, the average size of NPs grows and the Th clusters disappear.

The formation mechanism, reactivity, and environmental interaction of actinide nanoparticles is of substantial concern, but still poorly understood. Especially the accurate structural characterization of NPs is fundamental and not trivial, especially when NP size is < 3nm and samples are not homogeneous. The case of ThO₂ NP formation by chemical precipitation shows the importance of accurate structural characterization: many reports found that NP formation proceeds via an intermediate amorphous phase,[1] while other reported the direct formation of crystalline ThO2 NPs of size < 5nm.^[2] The key to correctly address the problem is to measure both short- and medium-range structural order to characterize the relevant length scales of the system. To this purpose, we measured HEXS and HERFD XANES at Th M₄ edge.^[3] The pair distribution function (PDF) extracted from HEXS probes both medium- and long-range order, while HERFD XANES probes the local environment of Th atoms.

EXPERIMENTAL. ThO₂ NPs of different sizes were synthesized by sequential heat treatment of freshly precipitated Th^{IV} samples. Characterization with XRD and HRTEM shows the presence of NPs <3nm in freshly precipitated samples and of NPs of increasing size with annealing at increasing temperature. HEXS and Th M₄ HERFD XANES were measured on ID15A and ID26 beamlines at the ESRF.

RESULTS. The PDF of bulk ThO₂, 150 °C annealed and 40 °C dried Th^{IV} samples shows a progressive decrease of size (Fig. 1a). A fit based on NPs of different size allowed determination of the NP distribution in each sample. The 40 °C dried sample contains a large amount of Th clusters similar to reported Th hexamers (60%), while mild annealing at 150 °C promotes recrystallization to bigger NPs and the disappearance of Th hexamers (see fit component legend in Fig. 1a). Th M₄ edge HERFD XANES of the 1200 °C sample (Fig. 1b) shows four characteristic peaks, labelled A-D. The 40 °C samples only show A and B and this change can be correlated to reduced size. Because M4 edge XANES probes the Th f-orbitals, the spectrum can be interpreted based on the Th f-density of states (DOS), as shown in the inset of Fig. 1b, where simulated f-DOS of bulk ThO2 and of Th hexamer from literature are shown below experimental data. The disappearance of feature C is well reproduced by the simulations, confirming that changes in the spectral shape are correlated to the presence of Th hexamers, for which Th local structure differs from that in bulk ThO₂.



Fig. 1: PDFs of bulk ThO₂, on 150 °C annealed and 40 °C dried Th[№] precipitates. Fit results are shown in red on top of data and the resulting size distribution indicated in the text box (a). Experimental Th M₄ HERFD XANES in the main panel. Inset reports the comparison of spectra on 1200 °C annealed and 40 °C dried samples with simulated f-DOS of bulk ThO₂ and Th hexamer (b).

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Single-step, high pressure, and two-step spark plasma sintering of UO₂ nanopowders

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Highly sinter-active UO_{2+x} nanopowders were densified using three different Spark Plasma Sintering (SPS) treatments. For the first time, a two-step method was adopted in SPS (2S-SPS), and its outcome was compared with the more established single-step (SPS) and highpressure SPS (HP-SPS). The 2S-SPS and HP-SPS treatments were optimized to limit coarsening while reaching the same density as in SPS. Such samples can be extremely helpful to study the effect of grain size on nuclear fuel properties (such as its oxidation behavior).

SPS is a field assisted sintering technique (FAST) that enables powder densification at relatively low sintering temperatures and hold times with respect to conventional sintering.^[1] The powder is loaded into a graphite assembly (die and pistons), through which pulsed direct current is passed, quickly heating up the sample due to the Joule effect.^[2] The graphite environment results in a reductive atmosphere above 600 °C that is beneficial for the processing of UO_{2+x} powders.^[3] HP-SPS and 2S-SPS are two approaches that can be adopted to try to reduce the final grain size, in order to synthesize surrogate materials reproducing the spent nuclear fuel (SNF) high-burnup structure (HBS), which are characterized, among other features, by 100-300 nm grains.^[4] In HP-SPS, SiC inserts are added to the assembly, allowing the application of higher pressures through the pistons (500 MPa instead of 100 MPa), and thus providing extra driving force for densification (i.e. lowering the sintering temperature).[5,6]

Two-step sintering instead involves the quick firing of the sample to a relatively high temperature (T_1) , and then a long hold time at a lower temperature (T₂) where densification is completed. Taking advantage of the different activation energies for grain boundary diffusion and grain boundary migration, with the correct choice of T₁ and T₂ densification can take place in the second step while grain growth is inhibited. For this to happen, the sample has to reach sufficient density in the first step (about 75 % TD).^[7]

EXPERIMENTAL. Batches of roughly 1 g of UO_{2+x} powder were produced by loading 2g of oxalate powder in an autoclave together with 5 mL of water and 0.5 mL of hydrazine, and heating up at 170 °C for 5h under autogenic pressure. The sintering treatments are summarized in Tab. 1. All samples were densified to 95% of the theoretical density (95% TD). The parameters reported for 2S-SPS and HP-SPS are already optimized to limit coarsening during densification.

Tab. 1: Processing parameters (sintering temperature, applied pressure, hold time, and heating rate) used in the SPS, 2S-SPS, and HP-SPS treatments.

SPS	2S-SPS	HP-SPS
1600°C	650–550°C	660 °C
70 MPa	70 MPa	500 MPa
10 min	3s-100 min	30 s
200°C/min	200°C/min	100°C/min



Fig. 1: SEM SE images of the microstructures obtained by applying SPS (a), 2S-SPS (b), and HP-SPS (c) treatments to the UO_{2+x} nanopowders. All samples have the same density (95 % TD) but differ significantly in grain size.

RESULTS. XRD characterization of the as-produced powder showed a lattice parameter of 5.464 ± 0.001 Å, indicating a stoichiometry of UO2.04, and Williamson-Hall analysis returned a crystallite size of 5 ± 1 nm.

The microstructures obtained by applying the treatments summarized in Tab.1 are shown in Fig.1. The grain sizes calculated with the ASTM E112-12 intercept method^[8] were: $3.08 \pm 0.06 \mu$ m for the single-step SPS, 478 ± 17 nm for the 2S-SPS, and 163 ± 9 nm for the HP-SPS. Due to the presence of SiC components, the sample produced by HP-SPS was slightly hyperstoichiometric, and therefore required a supplementary annealing step at 600 °C under Ar-4 %H₂ for 2 h to be reduced to $UO_{2,00}$.

The microstructure obtained by HP-SPS is of extreme interest for the study of SNF properties, as it settles in the range of grain size of the HBS. On the other hand, the much simpler 2S-SPS setup (not involving additional SiC components) also proved to be a suitable route to produce UO₂ with relatively small final grain size. Such samples will be used for a study on the effect of grain size on the oxidation behavior of UO2.

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Trimeric uranyl(VI)-citrate sandwiches Na⁺, Ca²⁺, and La³⁺ in aqueous solution

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The uranyl(VI) oxygen atoms (O_{yl}) of two 3:3 uranyl(VI)-citrate complex molecules sandwich a sodium ion in single-crystals.^[1] Nuclear magnetic resonance (NMR) spectroscopy supported by density functional (DFT) calculations provided unambiguous evidence for this U(VI)-O_{yl}-Na⁺ Lewis acid-base adduct, persisting in aqueous solution above a critical concentration of 3 mM uranyl(VI) citrate. Coordination to O_{yl} is not limited to Na⁺ only, but applies also to di- and even trivalent cations of comparable ionic *radii*. That is, Na⁺, Ca²⁺, and La³⁺ are sandwiched, while Li⁺, K⁺, and Rb⁺ are not, with uranyl(VI)-oxo Lewis base coordination to aqueous Ca²⁺ and La³⁺ being unprecedented.

EXPERIMENTAL. Details of both sample preparation and quantum chemical calculation can be found in our recent paper.^[2] Briefly, D₂O solutions of uranyl(VI) citrate, U(VI)-Cit, were prepared from stock solutions of uranyl(VI) nitrate and sodium citrate or citric acid. NaOD (for sodium containing samples, pD 6.9 ± 0.1) and LiOD were used for pD adjustment. The latter was used for sodium-free samples prepared in the presence of the following metal ions: Li⁺/K⁺/Rb⁺ (pD 6.9 ± 0.1), Ca²⁺ (pD 5.4) or La³⁺ (pD 4.8). NMR spectra were obtained at 298 K on a 14.1 T Agilent DD2-600 NMR system using 5 mm quartz NMR tubes.

RESULTS. Formation of the 6:6 U(VI)-Cit Na⁺ sandwich complex is evidenced by the observation of three distinct ²³Na NMR signals (Fig. 1A). The narrow signal at about 0 ppm corresponds to free Na⁺ aquo ions (*aq*). As can be seen, two further distinct Na⁺ sites are present, whose signals show remarkable upfield shifts. Upon coordination by six O_{y1} the sandwiched Na⁺ (*s*) senses the strongest shielding, hence resonating at –25 ppm. The third signal shows a chemical shift (–9 ppm) with magnitude between those of the free and sandwiched Na⁺. It is thus assigned to peripheral Na⁺ (*p*) strongly localized at the highly negatively charged complex [Na{\eta³-((UO₂)₃(µ₃-O)(Cit)₃)₂]¹⁵⁻.

The well-resolved ¹H NMR resonances (Fig. 1B and D) are likewise characteristic for the sandwich complex as its conformation is fixed. That is, the particular arrangement prevents intramolecular site exchange of U(VI)-bound and free –CH₂COO moieties whereupon the latter are kept in an all*cis* conformation, *i.e.*, for either 3:3 sub-unit being on the same side of the molecular plane (Fig. 1E and F).

Based on dilution series (Fig. 1C and D) we found that above a critical concentration of 3 mM uranyl(VI) citrate at neutral pD conditions two units of the 3:3 complex form the sandwich superstructure. The spectral features associated with the Na⁺ sandwich complex – ²³Na signals *p* and *s* as well as ¹H signals *a* through *d* – remain detectable until the sample concentration falls below 3 mM. As of this point, the sandwich complex falls apart in two 3:3 U(VI)-Cit units. Since the latter species reveals no steric demands on conformation, intramolecular dynamics facilitate easily, and the ¹H signals of the different –CH₂COO moieties average to the unresolved feature *e* in Fig. 1D.

Based on both NMR spectroscopy and DFT calculations, a corresponding aqueous solution molecular model was derived respecting these three types of Na⁺ (Fig. 1E).

Surprisingly, the sandwich superstructure forms also in the presence of other important Lewis acids such as Ca²⁺ and La³⁺ as inferred from the sharp ¹H signals obtained from so-



Fig. 1: Spectral deconvolution of the experimental ²³Na NMR spectrum obtained from the sample containing the 6:6 U(VI)-Cit Na⁺ sandwich complex (A). ¹H NMR spectra associated with the 6:6 U(VI)-Cit sandwich complexes of Na⁺, Ca²⁺, and La³⁺, respectively (B). ²³Na (C) and ¹H (D) NMR spectra of a dilution series; from top to bottom: [U] = [Cit] = 125, 25, 12.5, 6.2, 3.1, and 2.5 mM. Correspondingly derived aqueous solution molecular model of the Na⁺ sandwich complex, respecting three types of sodium: aquo ion (*aq*), peripheral ions (*p*), and the sandwiched ion (*s*) (E). Generic structure of one of the sandwich complex' two constituting trimeric U(VI)-Cit species together with proton labelling used for ¹H NMR signal assignment (F).

dium-free samples containing the nitrate salts of calcium or lanthanum (Fig. 1B). Interestingly, varying the nature of the present metal ions, only Na⁺, Ca²⁺, and La³⁺ with ionic radii of 1.02, 1.00, and 1.03 Å, respectively, match the sandwich cavity, while Li⁺ (0.76 Å) is too small, and K⁺ (1.38 Å) and Rb⁺ (1.52 Å) are too large.

Overall, these results are important for the development of an advanced understanding of environmental uranium chemistry, and once more demonstrate the value of NMR spectroscopy in actinide complex structure determination.

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Pu(III)-malate complexation study using isothermal micro-titration calorimetry

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Isothermal micro-titration calorimetric (ITC) experiments of the complexation reaction between Pu(III) and malate were performed. Preliminary values of the complexation constant (β) and reaction enthalpy (Δ H) were determined.

Up to now, there is a lack of thermodynamic data for the interaction of trivalent plutonium with organic ligands, which are relevant for the long-term assessment of a potential highlevel nuclear waste disposal. Malic acid (Mal), an α -hydroxycarboxylic acid, is a model ligand for complex organic matter like humic substances or concrete admixtures, that contains both hydroxyl and carboxyl groups. ITC permits the simultaneous determination of reaction enthalpies and stability constants in a single experiment at a given temperature and ionic strength. Micro-ITC is the method of choice for calorimetric experiments with actinides, as the radioactivity can be minimized owing to the very low required sample volume of ~200 µL.

EXPERIMENTAL. An aqueous Pu(III) solution (~3.6 mM , pH~0.7) in chloride medium was adjusted to pH3.95 by stepwise addition of aliquots of aq. NaOH (0.1, 1, and 5 M) under continuous electrolysis in an inert gas glove box. Ionic strength was adjusted to 0.6 m by addition of 67 μ L of 6 m NaCl solution (pH3.95) to 1 mL of the Pu(III) solution. Thus, the final Pu(III) concentration was 3.2 mM. For the ITC experiments (iTC₂₀₀, GE Healthcare), three aliquots (~300 μ L) of the Pu(III) sample were titrated in 1 μ L steps with solutions containing 73.1, 187.3, or 387.5 mM Mal and 6 m NaCl until pH3.95 was reached. The ITC method and the data analysis are described elsewhere.^[1] The oxidation state of Pu was controlled by UV-vis (Cary300, Agilent) after the sample preparation, and before and after the ITC experiments.

RESULTS. The Pu absorption spectra show pure Pu(III), described by the two characteristic bands between 575 and 625 nm and the weak band at 665 nm (Fig. 1).^[2] Neither dissolved Pu(IV), which is mainly characterized by a strong absorption band at 470 nm,^[2] nor Pu(IV) colloids were detected in the samples and no oxidation of Pu(III) to Pu(IV) occurred during the ITC experiments (Fig. 1).

The calorimetric heat curves of the Pu(III)-Mal titration (Fig.2) are very similar to those of analogous experiments with Nd3+ and to the Am³⁺/Nd³⁺-ITC measurements in Taube et al. 2019.^[1] All curves show two features with different slopes indicating more than one reaction. Based on the 1:1 and 1:2 complexation model, the heat curves were fitted using the program NITPIC-SEDPHAT.^[1,3] Different parameter sets were evaluated. The best fits and most plausible results were obtained when fixing the value for ΔH_2 . For this, we used the Nd(III) and Am(III) literature values (Tab. 1). The ΔH_1 of Nd(III)-Mal agrees with previous results, $^{[1]}\Delta H_1$ of Pu(III)-Mal is lower than the Am(III) and Nd(III) values. This could be explained by small differences of the inert salt concentration of the Mal titrant in comparison to the Pu(III) sample. At pH3.95, the values of $\log \beta_1$ and $\log \beta_2$ are one and two orders of magnitude smaller than at pH5.5, respectively, for both, Pu(III) and Nd(III) (Tab. 1). In analogy to previous ITC experiments of Am(III) complexation reactions, respective Pu(III) experiments are now feasible.



Fig. 1: Absorption spectra of Pu(III) samples at different times after sample preparation: 0 h (red), 4 h (after ITC-Exp. 1, green), 16 h (before ITC-Exp. 2, blue), and 19 h (after ITC-Exp. 2, black).



Fig. 2: Integrated calorimetric heat curves of the Pu(III)-Mal and analogue Nd(III)-Mal titration, [Pu³⁺] = 3.2 mM [Nd³⁺] = 3.1 mM; heat curves were corrected for the dilution heat of the Mal.

Tab. 1: Conditional thermodynamic data of $[Pu(Mal)_n]^{3-2n}$ (n = 1, 2) complexes at I = 0.6 m, compared with $[Am(Mal)_n]^{3-2n}$ and $[Nd(Mal)_n]^{3-2n}$.

Metal ion	$\log \beta_1$	ΔH1 (kJ/mol)	$log \beta_2$	ΔH ₂ (kJ/mol)
Pu ³⁺	$2.8\!\pm\!0.6$	-1.3 ± 0.6	$4.5\!\pm\!0.9$	-4 or -6 #
Nd ³⁺	3.0 ± 0.4	-0.3 ± 0.4	$4.7\!\pm\!0.2$	-4 #
$Am^{3+*[1]}$	3.9 ± 0.1	-0.4 ± 0.2	6.6 ± 0.1	-6 ± 1
Nd ^{3+*[1]}	3.8 ± 0.8	0.1 ± 0.2	6.3 ± 1.4	-4 ± 2

error: 2σ value, [#]: fixed value, ^{*}: I = 0.65 m, pH 5.5.

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Low-valent plutonium precursor compounds for metal-organic complexes

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The synthesis of a stabilized tetravalent plutonium chloride (*i. e.* [PuCl₄(dme)₂], dme = 1,2-dimethoxyethane) starting from PuO₂ under ambient conditions is reported. Spontaneous reduction of the desired product yielded an unprecedented trivalent plutonium chloride [PuCl₃(dme)]_n which holds potential for the synthesis of molecular trivalent plutonium complexes.

The availability of suitable anhydrous actinide starting materials is crucial for the synthesis of actinide complexes with organic or inorganic ligands. Based on previous work investigating the synthesis of tetravalent starting materials for thorium, uranium, and neptunium,^[1] we herein report a synthetic procedure to yield the tetravalent plutonium chloride [Pu^{IV}Cl₄(dme)₂] (1). The synthesis of this compound has been described previously starting from a solution of ²³⁹Pu^{IV} in hydrochloric acid.^[2] We have expanded this procedure to enable the usage of ²⁴²PuO₂. The group of Scott *et al.* already investigated the limited redox stability of 1 in tetrahydrofuran solution yielding a mixed Pu^{III}/Pu^{IV} complex.^[3] We have observed similar reactivity of 1 in dme yielding the trivalent plutonium complex [Pu^{III}Cl₃(dme)]_n (2).

EXPERIMENTAL. All preparations have been performed in dedicated glove boxes to handle milligram amounts of ²⁴²Pu. All manipulations have been performed under inert conditions to prevent oxidation or hydrolysis reactions. The used solvents have been dried over molecular sieve 3 Å prior to use.

RESULTS. PuO₂ has been dissolved in a mixture of 14 M nitric acid and 0.1 M hydrofluoric acid in 10:1 v/v ratio by stirring at 80 °C over two days to yield a dark green solution. The solvents have been subsequently evaporated in a stream of argon at elevated temperatures. The residue, a Pu^{IV} nitrate, is dissolved in concentrated hydrochloric acid generating a dark red solution accompanied by the evolution of nitrous gases. The solvent is again removed by evaporation. The procedure of dissolving in concentrated hydrochloric acid and evaporation of the solution is repeated three times to finally yield a dark red residue. This residue is dissolved in dimethoxyethane yielding a gold-colored solution. Residual moisture is removed by adding a few drops of trimethylsilylchloride. The supernatant is separated from a small amount of insoluble compounds and the volume is reduced by slow evaporation. The dark reddish brown residue is washed with diethylether to finally yield [Pu^{IV}Cl₄(dme)₂] (1). The oxidation state of 1 has been verified by UV/vis spectroscopy (see Fig. 1).^[2,4]

Over time, the formation of dark blue crystals was observed in a concentrated solution of **1** in dme. Analysis of these crystals by single-crystal X-ray diffraction revealed the formation of an unprecedented trivalent plutonium chloride which is stabilized by dimethoxyethane $[Pu(\mu_2-Cl)_3(dme)]_n$ (see Fig. 2). Each plutonium is coordinated by six chlorides and two oxygen atoms of dme forming a square antiprismatic coordination sphere. Three chlorides each are μ_2 -bridging to the next plutonium ion forming infinite chains along the [010] direction.

Such a polymeric structure is similar to $[AmCl(\mu_2-Cl)_2(thf)_2]_n$ as previously reported,^[5] although the same all- μ_2 -bridging chloride constitution has only been observed for a trivalent lanthanum chloride (*i.e.* $[La(\mu_2-Cl)_3(thf)_2]_n)$.^[6]



Fig. 1: UV/visible spectra of $[Pu^{U}Cl_4(dme)_2]$ (1, red) and of $[Pu^{UI}Cl_3(dme)]_n$ (2, black). The concentrations are approximately 10 mM.

The Pu^{III} compound **2** is soluble in pyridine yielding a bluecolored solution. UV/visible spectroscopy revealed a significant change of the absorption bands which is in accordance with the complete reduction of the tetravalent plutonium complex **1** to the trivalent compound **2** (see Fig. 1). As no additional reducing agent has been used, these results underline the redox sensitivity of tetravalent plutonium in organic solvents, even under inert atmosphere.^[3] Thus, the choice of the ligand system is crucial for the synthesis of oxidation state pure plutonium complexes. Both low-valent plutonium precursor compounds **1** and **2** will be used for the synthesis of metal-organic plutonium complexes in the near future.



Fig. 2: Molecular structure of [PuCl₃(dme)]_n (2). View along [100]. Hydrogen atoms have been omitted for clarity. Color code: plutonium (Pu, ochre), chlorine (Cl, green), oxygen (O, red), carbon (C, gray).

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Three series of tetravalent actinide (An = Th, U, Np, Pu) monosalen complexes [$AnCl_2(salen)(pyx)_2$] ($H_2salen = N,N'$ -bis(salicylidene)ethylenediamine; pyx = pyridine, 4-methylpyridine, 3,5-dimethylpyridine) were synthesized and studied regarding their affinity for neutral N-donor molecules. Single-crystal (SC) XRD as well as IR spectroscopy revealed isostructural series with conformational stability. Experimentally determined bond lengths An-Cl and An-N_{salen}/An-N_{pyx} point to an ionic binding behavior, whereas the trend of An-O distances clearly indicates covalent interactions for U, Np, and Pu. Quantum chemical (QC) calculations confirm these findings including that An-O bonds have the highest degree of covalency.^[1]

The question of covalency in early actinide binding can be approached by studying isostructural series in which the actinide (*An*) has one specific oxidation state, the longest such series being *An*(IV). Besides the actinide, the electronic and geometric properties of the ligand influence the complex properties. Schiff bases are known to be versatile ligands for almost every metal ion and oxidation state. They also allow partially saturated coordination spheres with labile complex positions and thus potentially reactive sites at transition metals or *An*, *e.g.* for biomolecules or as a basis for catalysis.^[2] Here, we investigate the solvent exchange on labile monosalen *An*(IV) complex positions using pyridine (py), 4-picoline (pic) and 3,5-lutidine (lut) using SC-XRD and quantum chemical calculations in order to analyze bonding trends with a focus on covalent contributions.

EXPERIMENTAL. All syntheses were carried out in nitrogen-filled glove boxes or using Schlenk technique for a strict exclusion of oxygen and moisture. The actinide chloride starting materials were synthesized according to literature procedures. Solvents were distilled and stored over 3 Å molecular sieve prior to use. The synthesis of $[AnCl_2(salen)(pyx)_2]$ (An = Th, U, Np, Pu; pyx = py (1-4), pic (**5–8**), lut (**9–12**)) was performed *via* $[An(salen)_2]$ and $[AnCl_2salen(MeOH)_2]$.^[1] A general synthesis route of **1–12** is as follows: The respective monosalen complex with methanol (*i. e.* $[AnCl_2salen(MeOH)_2]$) is dissolved in py, pic, or lut. Layering of these solutions with pentane yields single crystals.

RESULTS. Each An(IV) monosalen complex series with py (1-4), pic (5-8), and lut (9-12) isomorphously crystallizes in P-1, Aea2 or $P2_1/c$, respectively. Figure 1 exemplarily shows $[AnCl_2(salen)(pic)_2]$ (5-8).

In all complexes, the actinide is eightfold coordinated by a tetradentately binding salen molecule, two chloro ligands, and two pyridine-based molecules (*pyx*).

Among the *pyx* series, almost identical distances An–O, An–N_{salen}/An–N_{*pyx*} or An–Cl were found demonstrating a mi-



Fig. 1: Molecular structures of [AnCl₂(salen)(pic)₂] (An = Th, U, Np, Pu (from left to right)) with ellipsoids drawn at 50 % probability level and hydrogen atoms emitted for clarity. (C: dark grey, O: red, N: dark blue, Cl: light green, Th: light blue, U: dark green, Np: orange, Pu: red).

nor effect of the coordinating solvent *pyx* on the coordination environment. *An*–O bonds are the shortest with approx. 2.18 Å, compared to *An*–N_{salen} (\approx 2.59 Å) or *An*–Cl (\approx 2.72 Å). Interestingly, the *pyx* molecules occupy identical geometrical positions in **1**–**12** showing *An*–N_{*pyx*} distances (\approx 2.65 Å) that are about 7 pm longer than *An*–N_{salen} indicating weaker interactions. This explains the facile ligand exchange of *e.g.* MeOH and *pyx* at those positions.

For a detailed analysis of the bonding trends, structure optimizations and bonding were performed using Turbomole 7.3.1 and Orca 4.2.1.^[3,4] All optimized structures show only minor deviations (about 2 pm) from the SC-XRD structures. The trend of coordinating bond lengths is displayed in Fig. 2, showing a general decrease from Th^{IV} to Pu^{IV} with decreasing cation size. An-Cl and An-N_{pyx/salen} almost linearly follow the ionic radii of An(IV) with a coordination number (CN) of 8, indicating a purely ionic interaction.^[5] However, the An-O bond length suddenly decreases from Th to U followed by a plateau until Pu. This rather reflects the trend of covalent radii and suggests a significant contribution of covalent bonding. Calculated delocalization indices (DI) of 0.6 for An–O vs. 0.28 for An–N proof this hypothesis. Interestingly, the DI for the *An*–Cl bond of 0.56 is close to the DI for *An*–O. This points to strong interactions and a substantial degree of covalency in this bond which is usually assumed to be purely ionic.



Fig. 2: Coordinating bond lengths from SC-XRD in [*An*Cl₂(salen)(pic)₂] vs. the ionic radii (Shannon) for CN 8 with linear fits for *An*–N and *An*–Cl.^[5]

We conclude that the complexes $[AnCl_2(salen)(pyx)_2]$ possess two labile positions suitable for solvent exchange. The bonding analysis revealed major differences within the coordinating bonds, being mostly ionic for An–N but surprisingly covalent for both An–O and An–Cl.

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Mononuclear actinide (UO₂²⁺, U⁴⁺, Th⁴⁺, Np⁴⁺) complexes with a new Schiff base ligand

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The reaction of a new Schiff base ligand 2-methoxy-6-(((2-methyl-1-(pyridin-2-yl)propyl)imino)methyl)phenol (HL^{pr}) with actinide elements (UO₂²⁺, U⁴⁺, Th⁴⁺, and Np⁴⁺) results in the isolation of [UO₂Cl-L^{pr}·THF] (1), [ThCl₃-L^{pr}·2THF] (2), [UCl₃-L^{pr}·THF] (3), and [NpCl₃-L^{pr}·THF] (4). All crystal structures exhibit formation of mononuclear complexes. Binding studies confirm 1:1 stoichiometries and binding constants in the order Th⁴⁺ \approx U⁴⁺ > Np⁴⁺ > UO₂²⁺.

Coordination chemistry of actinide elements has gained widespread attention due to the actinides' underexplored coordinating properties as well as interesting redox and catalytic properties.^[1] Hence, efforts have been made to synthesize and isolate highly reactive, low-valent actinide compounds.^[2] Up to now, complexes with low valent actinides are largely dominated by cyclopentadienyl and related ligands.^[3] But also other organic ligands, such as di- and mono-anionic Schiff base ligands are explored due to the presence of strongly coordinating aryl-oxide donors along with a chelating motif as well as their non-innocent redox behavior.^[4,5] Here, we report the synthesis and isolation of mononuclear actinide complexes [UO₂Cl-L^{pr.}THF](1), [ThCl₃-L^{pr.}2THF](2), [UCl₃-L^{pr.}THF](3), and [NpCl₃-L^{pr.}THF](4).

EXPERIMENTAL. Condensation of o-vanillin with 2-methyl-1-(pyridine-2-yl)propane-1-amine hydrochloride results in the isolation of yellow oily product HL^{pr} in high yield. Complexes 1–4 were synthesized by treating HL^{pr} in its deprotonated form $(L^{pr})^-$ with $UO_2^{2+}(1)$ or M^{4+} (M = Th(2), U(3), Np(4)) in THF to afford mononuclear complexes.

RESULTS. Crystal structures of complexes 1, 3, and 4 exhibit seven-coordinated actinide complexes whereas Th⁴⁺ in 2 is eight coordinated. In all complexes, the metal center is present inside the ligand cavity comprising deprotonated phenoxide oxygen (O_{Ph}), imine nitrogen ($N_{C=N}$) and pyridyl nitrogen (N_{py}) while chloride atom(s) and THF molecule(s) complete the coordination sphere (Fig. 1).

The uranyl center in complex **1** is coordinated in a pentagonal bipyramidal geometry, whereas the geometry around U^{4+} and Np^{4+} center in **3** and **4** was found to be monocapped trigonal prismatic. Coordination from an additional THF molecule in **2**, results in the formation of an eightcoordinated trigonal dodecahedron geometry around Th⁴⁺.



Fig. 1: Ellipsoidal representation of complexes **1–4**. The Hydrogen atoms are omitted for clarity. Insets: polyhedral representation of coordination environment around central metal ion.



Fig.2: Capped stick representation of complexes 1–4, displaying the angle formed by phenoxide ring (pink line), pyridyl ring (red line) with the horizontal plane (green line) 1–4.

Importantly, coordination of metal ions to the ligand exhibits differently aligned aromatic rings around the horizontal plane containing An^{n+} , O_{ph} and (H)C = N (Fig. 2). In complex **1**, both phenoxide and pyridyl rings are inclined on the same side of the horizontal plane by 40.2° and 23.6°, respectively. Whereas, in complexes **2**–**4**, aromatic rings are present in a pseudo-trans manner to the horizontal plane. The phenoxide rings make angles of 9.0° in **2**, 15.6° in **3**, and 3.6° in **4** with the horizontal plane, whereas the pyridyl ring is inclined by 34.9°, 32.4° and 45.7° in **2**, **3**, and **4**, respectively. In contrasting observation, angle for phenoxide ring decreases and for pyridyl ring increases w.r.t horizontal plane, in the order $1 \rightarrow 3 \rightarrow 2 \rightarrow 4$.

Binding studies were performed by UV/vis spectroscopy by titrating $(L^{pr})^-$ with actinide ions. The changes in the absorbance reached maxima on addition of 1.0 equiv. metal ion indicating 1:1 stoichiometry. The binding coefficients were calculated (at 490 nm for **1** and 450 nm for **2**–**4**) using the Benesi–Hildebrand equation (1).^[6]

$$\frac{1}{(A-A_0)} = \frac{1}{K(A_{max}-A_0)[An^{n+}]^x} + \frac{1}{(A_{max}-A_0)}$$
(1)

The linear regression plot between absorption intensity, $1/[A-A_o] vs. 1/[An^{n+}]^x$ varied linearly as a function of $1/[An^{n+}](x = 1)$, confirming a 1:1 stoichiometry. The calculated binding constants (K) were found to be $2.3(\pm 0.2) \times 10^{-1} M^{-1}$ (1), $2.3(\pm 0.4) \times 10^{3} M^{-1}$ (2), $1.54(\pm 0.3) \times 10^{3} M^{-1}$ (3), and $2.01(\pm 0.4) \times 10^{2} M^{-1}$ (4). These numbers suggest that UO_2^{2+} binds weakly to $(L^{pr})^-$ compared to An^{4+} ions. Moreover, among An^{4+} , Th^{4+} and U^{4+} have similar binding affinities, whereas Np⁴⁺ has comparatively weaker binding affinity than Th⁴⁺ and U⁴⁺. These findings are in coalescence with ongoing quantum-chemical calculations to estimate binding and coordination properties of actinide elements with the new Schiff base ligand.

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Ion size-driven change of coordination environments in bis-pyrophen complexes

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The tetravalent thorium and neptunium complexes of the Schiff base ligand 1,2-benzenediamine-*N*,*N*'-bis(1*H*pyrrol-2-yl)methylene (pyrophenH₂) were synthesized via salt metathesis reactions using lithium diisopropylamine (LDA) as a base. Molecular structures were determined using SC-XRD and revealed distinct coordination environments for Th^{IV} and Np^{IV}, which presumably can be attributed to different ionic radii. Moreover, varying An–N bond lengths indicate different bonding situations to the pyrrol and imine N donor functionalities.

Understanding the complex chemistry of the early actinides is a huge challenge, going on for several decades now. For this purpose quite a lot of research work has been put into complex formation with salen-type ligands (salen = Bis(salicyliden)ethylendiamine), because of their capability to stabilize metal ions in different oxidation states and their easily changeable electronic and steric properties.^[1,2] Contrary to the extensively studied *N*,*O* donor ligand salen, analogous pure *N* donor systems have lagged behind. In order to characterize subtle differences in the actinide-nitrogen bonding, we have used the less studied salen relative 1,2-benzenediamine-*N*,*N'*-bis(1*H*-pyrrol-2-yl)methylene (pyrophenH₂) for complex formation with the tetravalent actinides Th and Np.

EXPERIMENTAL. Complexes were synthesized by adding the respective actinide chloride in thf to a stirring solution of the ligand with LDA base in thf. Color change occurred immediately from orange to dark red. After stirring for one day, volatiles were removed in vacuum and the red product was washed with pentane. Crystals were grown by diffusion experiments with pentane.

RESULTS. The synthesis route for the pyrophen complexes is displayed in Fig. 1. As the educts are used in a 1:2 metal to ligand ratio, the formation of a 1:2 complex is expected.



Fig. 1: Reaction scheme for the synthesis of 1 and 2.

SC-XRD measurements of **1** and **2** revealed the desired stoichiometry with pincer type coordination of both ligands (Fig. 2). The metal center is bound to eight nitrogen atoms of the two surrounding ligands, respectively. For **1** an additional thf molecule is coordinated to the metal, which is absent in the structure of **2**. This results in different coordination numbers (CN) for **1** and **2**, namely 9- and 8-fold. The resulting polyhedra can be termed as spherical capped square antiprism (in **1**) or triangular dodecahedron (in **2**). An explanation for the different coordination spheres could be the different ionic radii of the metal centers. Tetravalent Th with CN 9 exhibits a radius of 1.09 Å and Np with CN 8 has a radius of 0.98 Å.^[3] Since Np is 0.11 Å smaller, there is not enough space for an additional solvent molecule to coordinate to the metal.

A look at the mean bond lengths between the metal center and the donor atoms reveals remarkable differences. The ac-



Fig. 2: Molecular structures for [Th(pyrophen)₂(thf)] (1, left, CN = 9) and [Np(pyrophen)₂] (2, right, CN = 8). Protons and co-crystallized solvent molecules are omitted for clarity.

tinide bonds to the pyrrol nitrogen (An–N_{py}) are 2 or even 9 pm shorter in complexes 1 and 2 than those to the imine nitrogen (An-Nim). This indicates a stronger bond to the deprotonated nitrogen atoms, which has to be confirmed by theoretical bonding analysis. Furthermore, the metal ions are not located symmetrically within the cavity, the asymmetry being much more pronounced in 1 due to the additionally coordinating thf molecule. In 1, the ligand that is perpendicular to the thf molecule exhibits nearly identical distances between Th and the two different nitrogens, that is N_{py} (Th–N5: 2.534 Å, Th–N8: 2.533 Å) and N_{im} (Th–N6: 2.560 Å, Th–N7: 2.561 Å). In contrast, there are quite large differences in the Th–N_{py} (Th–N1: 2.524 Å, Th–N4: 2.555 Å) and Th-N_{im} (Th-N2: 2.552 Å, Th-N3: 2.572 Å) distances of the second ligand. Th is pushed towards N1, meaning the thf molecule is influencing the binding of the metal center to the pyrophen ligands. The asymmetrical coordination environment in 2 may also be attributed to the small ionic radius of Np^{IV}, making a symmetrical coordination energetically unfavored. Future work will focus on crystal growth and measurement of the U analogue to investigate whether there is an additional coordinating solvent molecule in its complex structure. On the one hand, the ionic radius for an 8-fold coordinated U center is 1.00 Å, which would point to a 2-like structure. On the other hand, a 9-fold coordinated U exhibits a radius of 1.05 Å, suggesting that coordination of an additional solvent molecule may be possible. Furthermore, the synthesis of the Pu complex ($r(Pu^{IV}) = 0.96 \text{ Å}$) shall aim to see if a 1:2 stoichiometry is still possible, when going to even smaller metal centers, and how Pu's bonding situation compares to the earlier actinides.[3]

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Understanding covalency in tetravalent actinide amidinates

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The emergence of covalency within the actinidenitrogen bond has been investigated in two series of isostructural tetravalent *f*-element amidinate complexes $[MX((S)-PEBA)_3]$ (M = Th, Pa, U, Np, Pu, Ce; X = Cl, N₃; (S)-PEBA: (S,S)-N,N'-bis(1-phenylethyl)benzamidinate). The combination of thorough characterization in solid state, in solution, and by quantum chemical bonding analysis reveals an increasing covalent character within the actinide series but a changing origin from mainly 6*d* to 5*f* participation.

The evaluation of the covalent character of actinide heteroatomic bonds is an ongoing hot topic in actinide science. Especially ligands containing soft donor atoms like nitrogen are well suited for these studies. Thus, we have synthesized and characterized two series of tetravalent actinide amidinate complexes with the chiral benzamidinate ligand (S,S)-N,N'-bis(1-phenylethyl)benzamidinate ((S)-PEBA). This constitutes a rare study of series of pure nitrogen donor complexes of tetravalent actinides including a transuranium element (*i. e.* Np).

EXPERIMENTAL. The synthesis of the tetravalent *f*-element amidinates $[MX((S)-PEBA)_3]$ (M = Th, U, Np, Ce; X = Cl, N₃) has been reported recently.^[1,2]

RESULTS. The reported tetravalent actinide amidinate complexes are isostructurally crystallizing in the chiral space group type $P2_12_12_1$ with similar cell parameters. In all reported complexes, the actinide is 7-fold coordinated by three chelating amidinate ligands and one (pseudo)halide forming a monocapped octahedral coordination geometry. The amidinates are coordinating asymmetrically towards the metal center (see Fig. 1) with a shorter and a longer M–N distance as indicated in Tab. 1.

The comparison of the intramolecular bond lengths of the different tetravalent actinide complexes reveals a decreasing trend with decreasing ionic radius of the tetravalent actinide cation. This applies to all intramolecular actinide-donor distances in a similar manner and implies a mainly ionic bonding interaction. However, a clear deviation from this trend is observed for the tetravalent cerium complex as the intramolecular M–Cl and M–N_a distances are significantly longer than expected by the decreasing ionic radius.



Fig. 1: Molecular structure of [NpCl((S)-PEBA)₃] as representative of the whole series of tetravalent metal amidinates. Hydrogen atoms have been omitted for clarity.

Tab. 1: Averaged bond lengths in [MCl(L) ₃] (M = Th, U, Np, Ce; L = (S)-PEBA)
complexes. Intramolecular distances M–N1, N3, and N5 are averaged
to M–N _a and M–N2, N4, and N6 to M–N _b .

d (Å)	IR (CN = 7) ^[3]	M-Cl	M–N _a	$M-N_b$
[ThCl(L) ₃]	1.00	2.701(1)	2.560(12)	2.439(13)
[UCl(L) ₃]	0.95	2.641(1)	2.499(4)	2.379(19)
[NpCl(L) ₃]	0.93	2.630(1)	2.497(4)	2.368(15)
[CeCl(L) ₃]	0.92	2.644(2)	2.512(6)	2.365(20)

To gain a more detailed picture, a quantum chemical bonding analysis using Quantum Theory of Atoms in Molecules (QTAIM) and Natural Population Analysis (NPA) was performed (see Fig.2). The delocalization index, which is a measure for the number of shared electrons between two atoms is constantly increasing within the actinides series from Th to Pu which indicates an increasing covalent character. This is supported by NPA which indicates an increasing population of the actinides' 6d orbitals from Th to U but a decrease to Pu, while the 5f excess population is constantly increasing from Pa to Pu. The combination of both methods led us to conclude that the covalent character of the actinidenitrogen bond is increasing within the actinide series from Th to U and then remains on a similar level until Pu. The reason for this is a change of the origin of this covalent character from mainly 6d to stronger 5f participation.

Furthermore, the delocalization index shows a similar covalent character of the tetravalent cerium to tetravalent uranium which for Ce is driven by occupancy of 5d orbitals.





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Following up on ThSiO₄ synthesis by in situ SWAXS measurements

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Thorite, ThSiO₄, is one of the most abundant natural sources of thorium on earth. While actinides are known to form nanoparticles in silicate medium, no direct link between those colloids and the crystalline form of thorite has been evidenced until now. We observed that thorite can be produced with experimental conditions close to environmental pH and temperature. Thanks to in situ Small and Wide-Angle X-rays Scattering (SWAXS) measurements, colloids of a few nanometers have been first evidenced for low reaction times. These colloids exhibit elongated shapes and ultimately tend to aggregate after the size has reached 10 nm. Once aggregated, the system goes through a maturation step culminating in the emergence of nanocrystallites exhibiting thorite zircon structure. This maturation step is longer when the reaction temperature is decreased highlighting kinetic considerations. These results have potential implications in the paragenesis of Th minerals deposits, but also in the behavior of Th and, by analogy, tetravalent actinides in the environment.^[1]

Thorium chemistry may be considered as indicative of the reactivity of the actinide elements at the oxidation state +IV. Due to the abundance of silicate species in the environment, identifying the actinide-silicate interactions is a key point to evaluate the actinide behavior in the environment. On the one hand, thorium, uranium and neptunium are known to form nanometric particles in silicate reactive media.[2-7] The prevailing size of these colloids is less than 20 nm, they appear as virtually amorphous or poorly structured and exhibit low isoelectric point ($pH_{IEP} \leq 4.6$), corresponding to a negative surface charge.^[4,6] Thus, An^{IV} oxyhydroxy-silicate colloids may be able to stabilize significant amounts of An^{IV} in a waterborne form and may enhance their mobility in soils and aquifers. On the other hand, the zircon-type structure minerals thorite (ThSiO₄) and coffinite (USiO₄) are among the most abundant natural sources of thorium and uranium in ore deposits and their syntheses have been studied using high temperature dry chemistry and hydrothermal methods. However, although they could play an important role in the formation of ore deposits, to the best of our knowledge, no specific link between the colloidal species and the corresponding crystalline phases (i.e. thorite and coffinite) has been highlighted until now.

EXPERIMENTAL. Thorium nitrate and sodium metasilicate were mixed in aqueous nitric acid solution, the pH of the reactive media was then adjusted to a pH ranging between 3 and 6 with a concentrated NaOH solution.

Small volumes of these solution were taken and placed in sealed glass capillaries in order to perform *in situ* SWAXS measurements at fixed temperature, ranging between room temperature and 60 °C, in a specific oven. Complementary ex-situ experiments were performed by keeping the mother solution at the same temperature in order to confirm the formation of ThSiO₄.

RESULTS. This experiment allowed us to follow the formation of ThSiO₄ from the aggregation of colloidal particles, characterized by the SWAXS curves at low q values at the beginning of the experiment and the crystallization of thorite



Fig. 1: In situ SWAXS measurements kinetics at room temperature (RT), 40 °C and 60 °C for pH = 3. The inset graphics refer to a focus in the q range from 10 nm⁻¹ to 30 nm⁻¹.

for longer holding time characterized by thorite diffraction peaks at high q values (Fig. 1). It has been evidenced for the first time that the formation of thorite is possible at environmentally representative temperatures. Moreover, the results proved that the temperature has a strong impact on the kinetic of aggregation of thorium silicate particles and their crystallization.

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High-energy-resolution X-ray spectroscopy studies on Th compounds at the Th $L_{\rm 3}$ and $L_{\rm 1}$ edges

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We report here the X-ray Absorption Near Edge Structure (XANES) at the Th L₃ and Th L₁ edges, re-corded in high-energy-resolution fluorescence-detection (HERFD) mode of Th-based compounds. The capacity of this method to overcome limitations of conventional XANES and to push the limits for obtaining new information about the electronic and geometric structures of several Th compounds is discussed in detail.^[1-3] Experimental results are supported by ab-initio computational studies within the FDMNES code. Moreover, all investigations are corroborated by X-ray photoelectron spectroscopy (XPS) and Anderson impurity model calculations. We show that ThL₃ edge is much more sensitive to the electronic structural changes compared to the ThL1 edge. However, despite the large life-time broadening of the Th L₁ edge, the HERFD at the L₁ edge still shows the sensitivity towards changes in the electronic structure of the Th7p states, while the crystal field splitting effects of the Th 6d states are studied at the Th L_3 edge.

EXPERIMENTAL. Th(NO₃)₄·4.7H₂O was purchased from Merck. ThO₂ was synthesized at HZDR from Th(NO₃)₄·4.7H₂O by sintering in an oven at 1,000 °C for several hours. The structure was characterized as crystalline powder ThO₂ by X-ray diffraction. For the Th₂-(µ₂-OH)₂(H₂O)₁₂Cl₂]Cl₄·2H₂O synthesis, first, NH₄OH was added to 1 mL of a 0.7 M Th(NO₃)₄ solution to form ThO₂ (am., hyd.) precipitate. This precipitate was washed with distilled water until a neutral pH was reached. 1 mL of 1 M HCl was added to the precipitate. The formation of crystals was observed after a few days.

Th L₃ edge (ROBL): The incident energy was selected using the (111) reflection from a Si DCM. Rejection of higher harmonics was achieved by two Rh mirrors at an angle of 2.5 mrad relative to the incident beam. The Th L₃ HERFD spectra were recorded by X-ray emission spectrometer set to the maximum of the Th L_{a1} emission line. The emission energy was selected using the (880) reflection of one spherically bent striped Si crystal analyzer with 0.5 m bending radius aligned at 84.4° Bragg angle.^[4]

RESULTS. The computed spectra agree well with the experimental data (Fig 1). The inspection of DOS reveals an overlap of Th d- and f- states and O p-states for ThO2. Previously reported DFT calculations on ThO2 also showed a mixing of Th f- and d-states.^[5] $[Th_2-(\mu_2-OH)_2(H_2O)_{12}Cl_2]Cl_4 \cdot 2H_2O$ also shows Th d- and f-states at the same energy as Cl p- and O p-states. However, in the white line the O states show very little overlap with Thd-states in contrast to the Cl p-states which seem to overlap with Th d-states. The preedge seems to arise from a mixture of Th d- and f-states with Cl and O p-states. Cl p-states appear to have a better overlap with Th states, observable by the same energy position and shape of the DOS distribution. Also, Cl p-states are less localized than O p-states. The localization of O p-states suggest that the O-ligands behave more ionic than in ThO2. DOS overlap between Th and Cl offer an interesting starting point for future investigations either with a water free compound substitution of Cl with other halogens. The or



Fig. 1: Simulated L₃ HERFD spectra (grey) and correspondeing local density of states. Only DOS above the Fermi level are shown.

Th(NO₃)₄·4.7H₂O is described in the literature as an ionic compound with no covalent bond formation. Our DOS illustrates a strong superposition of Th d states and N p-states in correspondence of the WL. A slight ad-mixture of O p-states is also found for the same spectral feature. The d-DOS of Th is particularly sharp for this compound, confirming its ionic character. Simulated and experimental spectra are found in good agreement and show a different broadening of the absorption edge on various Th systems. Calculated DOS for the three Th-based compounds show clear evidence of the broader distribution of the Th 6d states of ThO₂ compared to Th(NO₃)₄·4.7H₂O and [Th₂-(μ_2 -OH)₂(H₂O)₁₂Cl₂]Cl₄·2H₂O.

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Combining *ab-initio* multiple scattering (FEFF) calculations with Fourier filtering: eased derivation and validation of EXAFS structural models

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The *ab-initio* multiple scattering code FEFF allows the calculation of theoretical scattering phase and amplitudes functions based on structural models stemming from XRD structural data bases, quantum chemical calculations or from intuitive chemical information. Since FEFF is able to predict the experimental scattering functions with high accuracy, one may think about a direct comparison between theory and experiment. However, the Debye-Waller damping terms, which are also needed for the calculation of the theoretical EXAFS spectrum (χ_t) , cannot be predicted precisely by the theory or are unknown at all, thus spoiling a direct validation of a structural model. Here we propose a Fourier filtering algorithm (FFA) enabling a direct comparison of χ_t with an experimental EXAFS spectrum (χ_e) without least square fitting.

EXPERIMENTAL. The Debye-Waller factor (σ^2) cannot be predicted in a considerable precision by FEFF calculations, hence making a direct comparison between $\chi_t(k)$ and $\chi_e(k)$ (k – photoelectron wave vector) impossible. Thus, the idea behind the FFA is to remove or at least to reduce the influence of the σ^2 on $\chi_t(k)$ by Fourier filtering, which will be described in the following. The Fourier transform (F) of a signal $\chi(k)$ in the k-interval $k_{min} \le k \le k_{max}$ is:

$$F(\chi(k)) = \frac{1}{\sqrt{2\pi}} \int_{k_{\min}}^{k_{\max}} \chi(k) e^{i(2kr)} dk, \qquad (1)$$

and the inverse *F*⁻¹, *i.e.* the backtransform of *F*, is given by:

$$z(k) = F^{-1}\left(F(\chi(k))\right) = \frac{1}{\sqrt{2\pi}} \int_{r_{\min}}^{r_{\max}} F(\chi(k)) e^{-i(2kr)} dr, \qquad (2)$$

with r_{min} and r_{max} being the lower and upper limits of the filtered range in the *r*-space.^[1] Note that due to the complex nature of *F F*⁻¹ contains an imaginary and a real part, *i.e.* $\Im[z(k)]$ and $\Re[z(k)]$, respectively, which allows the calculation of the amplitude *A*(*k*) and the phase $\phi(k)$ of $\chi(k)$,

$$A(k) = \{\Im[z(k)]^2 + \Re[z(k)]^2\}^{1/2}$$
(3)

$$\phi(k) = \arctan\left\{\frac{\Im[z(k)]}{\Re[z(k)]}\right\},\tag{4}$$

so that the filtered $\chi(k)$ can be formulated as:

$$\chi(k) = A(k)\sin(\phi(k)).^{[1]}$$
(5)

This holds for any spectrum $\chi(k)$, while for instance for two backscattering atoms:

$$\chi(k) = A(k)\sin(\phi(k)) = N_h e^{-2\sigma_h^2 k^2} B_h(k,r) \sin(2kr_h + \varphi_h(k,r) + \xi(k)) + N_j e^{-2\sigma_j^2 k^2} B_j(k,r) \sin(2kr_j + \varphi_j(k,r) + \xi(k)),$$
(6)

where *N*, *B* and φ is the coordination number, the backscattering amplitude function and the backscatter phase shift of the atom type *h* and *j*, respectively, and ξ is the phase shift of the central atom. In the case of a FEFF calculation eq. 6 yields $\chi_i(k)$ with $\sigma_h^2 = \sigma_j^2 = 0$, where $F^{-1}(F(\chi(k))) = A_t(k) \sin(\phi_t(k))$. If we assume for an experimental spectrum the same underlying structural arrangement as used for $\chi_i(k)$ and a global σ^2 for each atom, $\sigma_h^2 = \sigma_j^2 = \sigma^2$, then:

$$\chi_e(k) = e^{-2\sigma^2 k^2} A_t(k) \sin(\phi_t(k)), \tag{7}$$

and with the theorem of convolution (*):

$$F\left(e^{-2\sigma^2k^2}A_t(k)\sin(\phi_t(k))\right) = F\left(U(k)\right) * F\left(\sin(\phi_t(k))\right)$$
(8)

$$U(k)\sin(\phi_t(k)) = F^{-1}(F(U(k)) * F(\sin(\phi_t(k)))).^{[1]}$$
(9)

In consequence and only if $\chi_e(k) = U(k)\sin(\phi_t(k))$, the assumed structural model used for the FEFF calculation will fit the structural arrangement measured by the experiment. U(k) is gained by the Fourier backtransform of $F(\chi_e(k))$ (eq.2) and applying eq.3, while $\phi_t(k)$ is calculated from $F(\chi_t(k))$ (eq.2) and eq.4.

For the test of the FFA, a Tc K-edge $\chi_e(k)$ is calculated with eq. 6, where $N_h = 6$ (oxygen atoms), $N_j = 2$ (iron atoms), $\sigma_h^2 = 0.005 \text{ Å}^2$ and $\sigma_j^2 = 0.009 \text{ Å}^2$ (Fig. 1a)). With the same values and with $\sigma_h^2 = \sigma_j^2 = 0$, $\chi_t(k)$ is calculated (Fig. 1a)). Due to the influence of the σ^2 , $\chi_t(k)$ is not comparable with $\chi_e(k)$, hence, the structural model cannot be validated. With FFA and the correct FEFF model the best match is observed (Fig. 1b), while for just small deviations from the correct FEFF model much larger deviations occur (Fig. 1c and d). Notably and even by violating in our test the FFA condition $\sigma_h^2 = \sigma_j^2$, FFA minimizes the influence of the unknown σ^2 , hence enables a validation of FEFF structural models for real-world cases.^[2,3] Moreover, FFA allows also for the derivation of structural models, simply by adjusting FEFF structural models with for instance any molecular drawing software.



Fig. 1: Tc K-edge EXAFS spectra (left) and corresponding F (right). $\chi_t(k)$ – green, $\chi_e(k)$ – black, $U(k) \sin(\phi_t(k))$ calculated by FFA for different structural models (red, values on the right). Residuals – blue.

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Signatures of technetium oxidation states: a new approach

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A general strategy for the determination of the Tc oxidation state involving X-ray Absorption Near Edge Spectroscopy (XANES) at the Tc L₃ edge is shown. A comprehensive series of ⁹⁹Tc compounds, ranging from oxidation states I to VII, was measured and subsequently simulated within the framework of crystal-field multiplet theory. This approach could be extended to numerous studies of Tc systems as this method is one of the most sensitive methods for Tc oxidation state and ligand characterization.^[1]

⁹⁹Tc is primarily produced in bulk quantities in nuclear reactors and is considered as unwanted radioactive waste. The short-lived metastable nuclide ^{99m}Tc (γ -emitter, $t_{1/2} = 6.01$ h) is the dominating isotope in routine diagnostic nuclear medicine with about 40 million administrations annually. Nine known oxidation states exist, from –I to +VII, with +IV, +V, and +VII being the most common. Lower Tc oxidation states are commonly supported by π -acceptor ligands such as isocyanides, CO or NO⁺. It has been shown that Tc^V, Tc^{III} and Tc^I compounds exhibit a diverse chemistry and are frequently used in nuclear medicine. Investigating these types of compounds and their specific coordination chemistry is crucial to determine and improve both their stability and their biological activity.

EXPERIMENTAL. The syntheses of (NBu₄)[TcNCl₄],^[2] (NBu₄)[TcNBr₄],^[2] (NBu₄)[TcOCl₄]^[3], $(NBu_4)[TcOBr_4]$ ^[4] $[TcCl_4(PPh_3)_2], ^{[5]} [TcCl_3(PPh_3)_2(CH_3CN)], ^{[6]} [TcCl_3(PMe_2Ph)_3], ^{[7]}$ [Tc(thiourea)₆]Cl₃,^[8] (NBu₄)[Tc(NO)Br₄],^[9] (NBu₄)[Tc(NO)Cl₄(MeOH)],^[10][Tc(NO)Cl₂(PPh₃)₂(CH₃CN)],^[11] and (NBu₄)[Tc₂(CO)₆Cl₃]^[12] followed published procedures. For the syntheses of (NH₄)₂[TcCl₆] and (NH₄)₂[TcBr₆], solutions of (NH₄)TcO₄ in HCl (37%) or HBr (48%) were heated under reflux for 10 min. (NH₄)Cl or (NH₄)Br were added and the products precipitated upon cooling as yellow or red crystals. The Tc L₃ edge XANES spectra were recorded at the INE-Beamline of the KARA (Karlsruhe research accelerator) facility in Karlsruhe, Germany. The entire sample environment was encased in a He box to allow for the detection of the low energy Tc fluorescence without any modifications to the sample.

RESULTS. The spectral shape of the Tc L₃ edge white line arises from transitions from the $2p_{3/2}$ to the 4d level. The Tc ground state electron configuration is [Kr]4d55s2. When Tc ionizes, it loses its valence 5s electrons before losing 4d electrons. Tc has seven valence electrons that can be lost, but the number of 4d electrons will be different only for Tc with oxidation states higher than +II. Therefore, the expected rising trend in the L₃ edge shift might be visible only for compounds starting from oxidation states +III to +VII (see Fig. 1). The simulated spectra reproduce the main edge features well, demonstrating the strong influence of crystal-field interactions in the Tc 4d shell. As expected, the Tc L₃ edge is much more sensitive to the oxidation state than Tc K edge XANES because the L₃ white line position depends directly on the number of 4d electrons. Additionally, Tc L₃ XANES is found to be more sensitive to the Tc speciation in various



Fig. 1: Representative increasing oxidation state series of Tc L_3 edge XANES and the simulated spectra based on crystal-field multiplet theory. A representation of the local geometry is provided for each compound on the right hand side.

compounds. It can be explained by the fact that 4d orbitals have a rather large radius and thus overlap strongly with orbitals of neighboring atoms. We expect that this approach can be extended to numerous studies of Tc systems and can stimulate further Tc coordination chemistry research relevant to various fields of applications.

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Implementation of a 'Rolling Ball' algorithm for spectroscopy

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The 'Rolling ball' algorithm was applied to spectroscopic and spectrometric data for background subtraction obviating the restrictions of polynomial and spline functions.

The adequate subtraction of background contributions in optical spectra containing a high noise level is often challenging, in particular for quantitative analysis. This is mainly due to the impossibility to isolate the intrinsic background from the signal. Polynomial and spline functions are usually applied to simulate native background. As the degree of the polynomials as well as the number and position of interpolation points could not be determined automatically, the selection of these parameters might be significantly biased by the user. The 'Rolling Ball' algorithm inspired by Stanley Sternberg is applied to reduce subjective assessment during data evaluation processes. This algorithm is implemented for images in the ImageJ software package.^[1]

METHODOLOGY. The algorithm was originally developed to remove gradient of images, *e.g.* of protein gels. Its name based on the idea, that the contact points of a ball rolling over (or under) an image is representing the background. For dimensional reduced data sets, the dimension of the ball is reduced to a circle (ellipsoid). A graphical illustration of the algorithm is presented in Fig. 1.



Fig. 1: Illustration of the Rolling Ball algorithm for data vectors. A circle (blue) is rolling along the data (green) and the maximum of its trace is used as background (gray).

Thus, the only user defined parameter is the radius of the ball. Once the radius is chosen, it could be applied for entire datasets, even with backgrounds of different shapes. This makes the background subtraction routine less user biased.

RESULTS. The 'Rolling Ball' algorithm has been applied for infrared (IR) absorption, Raman and luminescence emission spectra, and HPLC chromatograms. The ball radius was set to the length of the datasets and thus big enough to prevent the ball from dipping into signal. As shown in Fig.2, an appropriate background is found for all data sets. For luminescence spectra, the shapes of the backgrounds is dramatically different. An accurate background subtraction is not feasible with spline or polynomial functions without an adjustment of the parameters during processing.

The algorithm is compatible with noisy data after smoothing procedure. It had to be mentioned, that kinked or discontinuous backgrounds cannot be estimated by the rolling ball, since those effects are too similar to the signal.



Fig. 2: Examples for the application of the 'Rolling Ball' algorithm. Shown are experimental data (green), background generated by the 'Rolling Ball' algorithm (gray) and final data after background correction (magenta) for IR (A), Raman (B), fluorescence spectra (C), and an HPLC chromatogram(D).

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Electrochemical synthesis of aqueous Pu stock solutions in chloride media

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Aqueous Pu solutions of oxidation states +III, +IV, and +VI were electrochemically synthesized in chloride media. The purity grade obtained was about 100 and 85 % for Pu(III/IV) and Pu(VI), respectively, as monitored by *in situ* photometry. The long-term stability of the Pu(III) solution and the harsh conditions necessary for the oxidation of Pu(IV) to obtain Pu(VI) suggests a strongly stabilizing effect of Pu(III/IV) chloro complexes.

According to demands at the IRE from biological tasks and certain analytical techniques, a preparation of Pu stock solutions containing selective oxidation states at high purity grade in chloride matrix was mandatory. Therefore, we set up an electrochemical experiment providing an *in situ* monitoring of the prevailing oxidation states by photometry.^[1]

EXPERIMENTAL. All electrochemical experiments (potentiostat Autolab PGSTAT302, Metrohm) were performed with a cell consisting of a quartz glass vessel (~20 mL) and a home-built Teflon lid providing drillings for the fixation of the electrodes. The vessel is mounted on a special rack equipped with apertures for mounted fiber optics (Fig. 1).



Fig. 1: Technical drawing of the electrochemical cell (left) and set up of the experiment (right) with mounted electrodes and fiber optics for *in situ* photometric measurements.

Photometric monitoring was performed with a spectrometer setup from Avantes (AvaSpec-2048x14). All experiments were performed under continuous stirring in a glove box under inert gas atmosphere except where noted.

The initial ²⁴²Pu(IV) solution was purified from ²⁴¹Am(III) by extraction chromatography (TEVA Resin) to reduce the overall dose.^[2] After filtration through a 0.45 µm polypropylene filter to remove potential column residues, an Am-free Pu-stock solution was obtained (~12.5 mM Pu/~0.1 M HCl) as verified by γ - and α -spectrometry.

RESULTS. The initial solution (9 mL of ~3 mM Pu(IV) in 4M HCl) contained nearly 100% Pu(IV) as verified by photometry. *Synthesis of aqueous Pu(III).* Preliminary electrochemical experiments showed that Pu(III) could be successfully obtained in 4M HCl. Further experiments were performed at lower ionic strength (~0.1 M) which lowers the reduction rate, but a pure Pu(III) solution was eventually obtained after several hours of reduction at a voltage of -0.6 or -0.4 V. For further sorption and complexation experiments, the Pu(III) solution was adjusted to pH4 as described in [3]. This solution was stored in a sealed quartz glass vessel in the dark at ambient conditions for about six months showing almost no formation of oxidation products after this period of storage.



Fig. 2: Photometric spectra of the Pu solution recorded before and after 15 and 120 min of electrolysis at pH 4 demonstrating the formation of colloidal Pu(IV) and a very small fraction of Pu(VI).

Colloidal Pu(IV). The formation of colloidal Pu(IV) at pH 4 was observed by complete oxidation of Pu(III) at +1.7 V for 2 h (Fig. 2). The green-colored solution obtained remained suspended even without any applied stirring. After lowering the pH with 40 μ L 4M HCl and prolonged reduction overnight, interconvertible reduction to Pu(III) was completely achieved.

Oxidation to Pu(VI). As oxidation to Pu(VI) was observed neither at high ionic strengths (4 M HCl) nor in carbonate media in the alkaline pH range, the setup was placed in a fume hood where an influx of ozone into the reaction vessel was feasible. The evolution of Pu(VI) preferentially occurred at slightly acidic pH, reduced ionic strength, highly oxidizing conditions provided by O₃ and during prolonged electrolysis (Fig. 3). After 40 h of electrolysis, a Pu(VI) solution containing 85% Pu(VI) and 15% Pu(IV) was obtained.^[1]



Fig. 3: Photometric spectra of a Pu solution recorded at different times of electrolytic oxidation with influx of ozone. The arrows indicate the decreasing and increasing intensities of the predominant spectral characteristics of Pu(IV) (left panel) and Pu(VI) (right panel), respectively.

This technical paper demonstrates the feasibility of electrochemical preparation of aqueous Pu solutions showing a high selectivity of oxidation states in a chloride matrix. The harsh conditions necessary to generate Pu(VI) and the longterm stability of the Pu(III) solution generally suggests a high thermodynamic stability of the Pu-chloro-species.

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SCIENTIFIC CONTRIBUTIONS (PART II)

Long-Lived Radionuclides & Transport Phenomena in

GEOLOGICAL Systems

Sorption of curium(III) on crystalline rock material

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Cm(III) sorption on heterogeneous crystalline rocks was correlatively investigated using spatially-resolved techniques. Sorption uptake and speciation of Cm(III) between different and on the same mineral grains depends on mineralogy and surface roughness.

EXPERIMENTAL. Two thin sections were prepared from two crystalline rocks, a granite from Eibenstock, Germany, and a migmatised gneiss from Bukov, Czech Republic. Sorption experiments were performed over the course of a week using a 1 μ M solution of Cm(III) (pH8, 0.1 M NaCl). Using interferometry, the surface topography of the thin sections was characterized from which the surface roughness was calculated. The entire thin sections were further analyzed using calibrated autoradiography to quantify sorption uptake.

To evaluate the retention potential of the rocks, the speciation of Cm(III) on the mineral surface must be determined on the molecular level. To this end, smaller regions of interests (ROIs) in the range of 1×1 mm were investigated by µTRLFS (micro-focus time-resolved laser-induced fluorescence spectroscopy).^[1] Each luminescence map obtained consisted of 2,601 pixels with one pixel corresponding to one Cm(III) luminescence spectrum. Sorption uptake was then derived from each pixel's integral luminescence intensity (Fig. 1c). To obtain speciation information, all Cm(III) spectra were fitted using Gaussian functions. This provided the wavelength of the Cm(III) peak maxima (Fig.1d), where stronger red-shifts correlate to stronger bonds to the mineral grain's surface. To further analyze the speciation, the luminescence decay was measured for selected spots. From the determined luminescence lifetimes the number of species present at the surface and the number of water molecules in the first hydration shell of Cm(III) were derived. In addition, Raman microscopy revealed the underlying mineralogy (Fig. 1a).

RESULTS. Generally, the sorption uptake of Cm(III) was highest on mica-like minerals, intermediate to high on feld-spars, and smallest on quartz, which fits well with previous studies of Eu(III) sorption on the same crystalline rocks.^[1,2] A new finding is that this observation only reflects areas exhibiting a low surface roughness. With increasing surface roughness sorption uptake on quartz and feldspar increases by one to two orders of magnitude. Such an increased sorption uptake was also seen at surface features such as mineral grain boundaries, defects and cracks (Fig. 1b).

Highest luminescence intensities (LI) were measured on feldspar and epidote, while LIs on quartz were lower (Fig. 1c). Mica and amphibole did not show any luminescence since they are Fe-bearing minerals, therefore quenching the luminescence of Cm(III).

Speciation was investigated by examining the Cm(III) peak position. On quartz mainly strong inner-sphere sorption takes place (602–605nm), generally showing a homogeneous speciation of Cm(III) on the quartz surface (Fig. 1d).

On feldspar the speciation was found to be less homogeneous on both crystalline rocks (600-605 nm). On both, quartz and feldspar, the Cm(III) peak maximum was more red-shifted (603-605 nm) for regions with a higher surface roughness than for lower surface roughness areas (602-603 nm). This clearly shows that the surface roughness has a



Fig. 1: ROI analysis to determine mineralogy using Raman-microscopy (a), surface roughness using interferometry (b), sorption uptake using μ TRLFS (c), and peak maximum positions by Gaussian peak fitting and in red measured luminescence lifetime spots (d).

measurable effect on the speciation of Cm(III) on the mineral surface, and not just the sorption quantity.

To better characterize the speciation, lifetime analysis was performed, which showed that on quartz and feldspar with increasing surface roughness the number of water molecules surrounding Cm(III) decreases, meaning that Cm(III) surface complexation becomes stronger. This can be explained by areas of high surface roughness possessing a larger number of kink and defect sites.^[3] Sorption at these sites allows Cm(III) to form more bonds to the mineral surface, replacing water molecules from the hydration shell, while also inducing the previously described red-shift of the peak maxima. A similar explanation was hypothesized previously, regarding the inhomogeneous distribution of Eu(III) sorbed on quartz in Eibenstock granite.^[1]

In our studies we could show that sorption of Cm(III) and its speciation on the millimeter and centimeter scale is dependent on the mineralogy and the topography of the available surfaces on the nanometer scale. Surface roughness has an impact on the amount of Cm(III) sorbed on the surface as well as the speciation of Cm(III) on the mineral surface.

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Development of a generic surface complexation model for the sorption of trivalent actinides and lanthanides on orthoclase

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In this study a surface complexation model (SCM) for the sorption of various trivalent rare earth elements and actinides (unified to M³⁺) on orthoclase was developed. A generic approach was used to keep the model as simple as possible and simultaneously to cover many similar elements. Four surface complexes are responsible for the retention of M³⁺ on orthoclase, for which reliable formation constants were determined.

EXPERIMENTAL. The sorption of different trivalent lanthanides (Y, La, Nd, Eu, Lu) and Am on orthoclase was investigated under different experimental conditions ([M3+], ionic strength, grain size, solid-liquid ratio). The experimental procedure is described elsewhere.^[1] In total, 201 observations from 13 experimental series were used to develop an SCM. As the different trivalent metals showed a comparable behavior under the same experimental conditions (Fig. 1), only one generic trivalent metal, M³⁺, was considered. The thermodynamic data of Eu3+ in the ThermoChimie-TDB (10a, Version of Sept. 26, 2018) was used to calculate the aqueous speciation of M3+. The SCM development was a stepwise process and the model was evaluated after each addition of a surface complex (SC). To fit formation constants of considered SCs, the parameter estimation code UCODE 2014 (version 1.004)^[2] was coupled with the geochemical speciation code PhreeqC (version 3.5.2),^[3] while the diffuse layer model was applied.^[4] All observations were simultaneously used to fit surface complexation parameters (SCP). A bidentate coordination of M3+ to surface functional groups (≡S-OH) was assumed. A log K of −2.5 for the deprotonation of these groups was separately determined by column titration experiments.^[1]

RESULTS. Four SCs are necessary to adequately reproduce experimental batch data under the various experimental conditions. All formation reactions as well as their optimized formation constants are listed in Tab.1.

Tab. 1: Fo	ormation constants of the considered surface complexes with	ı their
no	onlinear confidence intervals.	

Species	Equation	Log K	Nonlinear conf. int.
SC-1	$\begin{array}{l} 2 \equiv S \text{-OH} + M^{3+} \leftrightarrow \\ (=S \text{-OH})_2 M^{3+} \end{array}$	-3.4	-3.1 / -3.8
SC-2	$\begin{array}{l} 2 \equiv S \text{-}OH + M^{3+} \leftrightarrow \\ (=S \text{-}O)_2 M^+ + 2H^+ \end{array}$	-8.0	-7.7 / -8.6
SC-3	$\begin{array}{l} 2 \equiv S\text{-}OH + M^{3+} + 2H_2O \leftrightarrow \\ (=S\text{-}OH)_2M(OH) + 3H^+ \end{array}$	-11.6	-11.3/-12.1
SC-4	$\begin{array}{l} 2 \equiv S\text{-}OH + M^{3+} + 2H_2O \leftrightarrow \\ (=S\text{-}OH)_2M(OH)_2^- + 4H^+ \end{array}$	-16.6	-16.1/-17.4

The species SC-1, in which M³⁺ is coordinated to two protonated surface groups, improves the fit for low M³⁺ concentrations at low pH values. Deprotonation of the surface groups leads to the formation of SC-2. Its two hydrolysis forms SC-3 and SC-4 are necessary to reproduce the immobilization at higher pH values. The existence of SC-2 to SC-4, which can



Fig. 1: Experimentally observed sorption of M³⁺ to orthoclase (symbols) and simulations for different selected experimental conditions (adapted and modified from [1]).

be seen as inner sphere complexes, was supported by timeresolved laser-induced fluorescence spectroscopy (TRLFS) measurements, using Cm³⁺ as luminophore.^[1]

Figure 1 shows experimentally determined values (symbols) for different experimental conditions and different M³⁺ in comparison with simulations (lines), using SC-1 to SC-4. The agreement between experimental and simulated values is very good for all conditions and all different M³⁺, which verifies the generic approach.

High robustness of the model is provided by the huge experimental data on which this model was built. Furthermore, the spectroscopic identification of the inner sphere complexes increases the reliability of the proposed SCM. The consideration of spectroscopic data during the fitting procedure, which is described in [1], leads to similar log K values. Furthermore, the proposed model yields reliable simulations, which were successfully validated for two individual and independent experimental literature data sets (detailed information is given in [1]).

The derived SCP values close an important gap in the thermodynamic database for the interaction of trivalent metals, relevant for nuclear waste disposal, and their interaction with mineral phases. Furthermore, they will be used to extend the database of the Smart K_d -concept,^[5] which is a modern approach for the simulation of radionuclide transport, which in turn contributes to a better safety assessment of repositories for radioactive waste.

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First sorption results of Np(V) on ZrO₂

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Batch sorption experiments of Np(V) onto zirconia (ZrO₂) revealed a fast, pH-dependent sorption process. The isoelectric point (pH_{IEP}) of the ZrO₂ was found at pH8.3 by zeta potential measurements. Np(V) concentrations in the lower micromolar range do not significantly impact the pH_{IEP}, whereas a shift to higher pH values was found with increasing Np(V) concentrations.

For the safety assessment of a nuclear waste repository, the interactions of long-lived radionuclides, such as neptunium, with corroded phases in the near field of the repository have to be considered. Zirconia is the corrosion product of the zircaloy cladding material of spent nuclear fuel rods and thus may represent one of the first barriers for radionuclides.

EXPERIMENTAL. All batch experiments (Tab. 1) were performed under inert gas (N₂) conditions with a ZrO₂ solid-toliquid ratio of $0.5 \, g \, L^{-1}$. Except for the time dependent investigations, samples were equilibrated for three days and subsequently centrifuged (1h; 4,020×g). Concentrations of Np-237 in the supernatants were derived from liquid scintillation counting. For experiments conducted at 0.01 M NaCl, zeta potential measurements were performed before centrifugation under exclusion of CO₂.

Tab. 1: Experimental conditions of the Np(V) batch sorption experiments. Concentrations and ionic strengths are given in mol L⁻¹.

Variable	pН	Time	[Np(V)]	I (NaCl)
Time	8	1h-7d	10-6	0.1
pH/I	3-10	3 d	10-6	0.01 or 0.1
[Np(V)]	9	3 d	$10^{-6} - 5 \times 10^{-5}$	0.01

The zeta potentials of the pure ZrO_2 suspension (0.5 gL⁻¹) were measured to determine the pH_{IEP} of the mineral. The samples were prepared with background electrolyte concentrations of 0.005 or 0.01 M NaCl and shaken for three days.

RESULTS. Time-dependent batch experiments revealed that complete sorption occurred after a few hours (data not shown). The pH series at 0.01 M NaCl showed a broad sorption edge reaching a maximum at pH \ge 7 (Fig. 1).

Under inert conditions, the predominant aqueous species is expected to be the fully hydrated neptunyl ion (NpO_2^{+}) in the investigated pH range.^[1,2] The presence of chloro species, such as NpO_2Cl (aq), can be neglected.^[3] Thus, the increased sorption at higher pH is due to interactions of the NpO_2^{+} species with a more negatively charged ZrO_2 surface. Two batch samples with a higher ionic strength (0.1 M NaCl)



Fig. 1: Sorption edges of Np(V) on ZrO₂ at different Np concentrations at (filled symbols) 0.1M NaCl and (open symbols) 0.01M NaCl (solid/liquid ratio: 0.5gL⁻¹).

showed no impact on the Np(V) sorption suggesting the predominant formation of inner-sphere complexes (Fig. 1). The pH_{IEP} of the solid phase was found at pH8.3 (Fig. 2) and is in agreement with literature values.^[4,5] Thus, complete sorption of Np(V) was already observed below the pH_{IEP} of the zirconia.

The zeta potential curves at 0.01M NaCl obtained in the presence and absence of 10^{-6} M Np(V) are shown in Fig. 2. All curves are in very good agreement, indicating no significant impact on the zeta potential by the sorption of Np(V). Further experiments were conducted at 0.01M NaCl and pH9 but with higher Np(V) concentrations of up to 5×10^{-5} M. The sorption curves and zeta potential plots are shown in Figs. 1 and 3, respectively.



Fig. 2: Zeta potential curves of ZrO_2 suspensions at different NaCl concentrations and in presence of 10^{-6} M Np(V) (solid/liquid ratio: 0.5 gL⁻¹).

Only the highest Np(V) concentration of 5×10^{-5} M showed that sorption is significantly reduced at alkaline pH values (Fig. 1). However, an increase of the zeta potential to nearly same extent could be observed for Np(V) concentrations higher than 10^{-6} M at pH9 (Fig. 3). These findings suggest a possible saturation effect occurring at higher Np(V) concentrations. Spectroscopic investigations (IR, XAS) will be conducted to gain a better understanding of the Np(V) sorption processes on the molecular scale.



Fig. 3: Zeta potential curves of ZrO₂ suspensions in presence of Np(V) at different concentrations (solid/liquid ratio: 0.5gL⁻¹; I = 0.01molL⁻¹ NaCl).

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Effect of background electrolyte composition on the formation of Th(IV) nanoparticles on the muscovite (001) basal plane

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We study the impact of background electrolyte composition on the sorption of Th(IV) on mica (001). Surface X-ray diffraction and *in situ* AFM show the formation of Th nanoparticles. Two main retention mechanisms are identified: At low [Th] = 0.1 mM, oligomerization occurs at the mineral-water interface (heterogeneous nucleation) and is strongly influenced by the electrolyte. A strong decrease in Th uptake is observed from Li to Cs in the alkali series, with Na as a clear outlier in this series, showing much lower uptake than expected according to the trend. At high [Th] = 3 mM, retention occurs by the formation of Th oligomers in solution (homogeneous nucleation) and their subsequent sorption onto the mineral surface. This second process appears to be largely independent of the electrolyte.

Actinide nanoparticle (NP) formation is known to have an influence on radionuclide mobility in the environment. For a reliable safety assessment of nuclear waste repositories, a fundamental understanding of the formation of actinide NPs as well as their interactions with charged mineral surfaces is thus important.

Previous studies investigated the retention of Th(IV) on muscovite using surface X-ray diffraction (SXD) and reported an unusual trend of Th(IV) uptake on muscovite (001), depending on the electrolyte composition (LiClO₄ > KClO₄ > NaClO₄).^[1,2] In case of LiClO₄, Th uptake (4.9 Th/A_{UC}, A_{UC}: area of mica (001) unit cell) was observed to be higher than required for compensation of the mica basal plane's surface charge (1 e⁻/A_{UC}), indicating particle formation. However, the reaction mechanisms remained unclear. Therefore, we applied SXD and *in situ* AFM to shed light onto the reactions causing the electrolyte effect on the retention of Th(IV) on muscovite (001).

EXPERIMENTAL. For SXD and *in situ* AFM measurements solutions of 0.1 and $3 \text{ mM Th}(\text{NO}_3)_4$ and 0.1 M ACl (A = Li, K, NH₄, Cs) at pH3.2 were prepared. The reaction occurred over night by dropping a freshly cleaved muscovite crystal into the sample solution.

SXD was performed at beamline 13-ID-C at the Advanced Photon Source. Experimental details can be found elsewhere.^[1] *In situ* AFM was performed using an Asylum Research Cypher ES instrument with Bruker MSNL-10 silicon tips in tapping mode with a resonance frequency of ~9kHz. SPIP software was used for image processing. For grain size analysis, particles were automatically detected by a height threshold of 1 nm.

RESULTS. At low [Th] = 0.1 mM, a rather broad Th electron density (ED) is observed for all investigated electrolytes ACI (A = Li, K, NH₄, Cs), reaching up to ~50 Å in case of LiCl (exemplarily shown in Fig. 1, top) and KCl. The observed Th uptakes decrease strongly within the alkali ion series and a linear correlation (R^2 = 0.9962) to the ionic radius of the electrolyte cation is observed (Fig. 2, left), suggesting competitive sorption of Th⁴⁺ and A⁺. With increasing ionic radius, the hydration energy of the electrolyte cation decreases, leading to an increased tendency to adsorb as inner-sphere sorption complexes,^[3] which seem to be preferable over the sorption



Fig. 1: Total electron density (lines) and Th contribution (filled area) of samples reacted with 0.1 and 3 mM Th(NO₃)₄ and 0.1 M LiCl at pH 3.2 ± 0.1.



Fig. 2: Th coverage for samples reacted with 0.1 mM (left) or 3 mM (right) Th(NO₃)₄ and 0.1 M ACl (A = Li, Na, K, NH₄, Cs) at pH = 3.2 ± 0.1 .

of Th⁴⁺ and inhibits the subsequent particle formation. However, the reported value for NaCl is a clear outlier in this series, showing lower Th uptake than expected according to the trend.^[2] The occupancies from SXD are in perfect agreement with the number of particles per area, obtained from grain size analysis of *in situ* AFM images. The particles have a median height of ~1–2nm and lateral size of ~10–20nm, typical for particles formed at mineral surfaces (heterogeneous nucleation), which is known to be influenced by the aqueous chemistry of the solution.

At high [Th] = 3 mM oligomerization of Th(IV) is known to occur in solution. The interfacial structures of all samples are dominated by a large amount of Th ED at a distinct height from the surface (~6.5 Å, see Fig.1, bottom). Here, LiCl (2.0 Th/A_{UC}), KCl (1.7 Th/A_{UC}), and NaCl (1.4 Th/A_{UC}) show comparable Th uptakes (Fig.2, right). Based on the results, it is suggested that Th oligomers form in solution (electrolyte-independent) by homogeneous nucleation, which sorb onto the surface.^[4]

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Tc removal by Sn and Mn pre-sorbed on alumina

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A well-known methodology to immobilize Tc consists of reducing the highly mobile $Tc^{VII}O_4^-$ to Tc^{IV} since its most stable species is a low soluble oxide. We have studied the aqueous scavenging of Tc by redox-active Mn^{II} and Sn^{II} pre-sorbed on alumina as a function of pH. It was shown that Tc was quantitatively removed ($K_d \approx 10^5$ L/g) from solution when using Sn pre-sorbed on alumina. K_d ranges from 50 L/g to 600 L/g when Tc is immobilized by Mn pre-sorbed on alumina.

Soluble Sn^{II} has been used for years as Tc reductant in solution (homoreduction),^[1] whereas the reductant behavior of Mn^{II} has not been analyzed up to date. Previous works reported that Tc reduction is highly enhanced when Fe is sorbed on a mineral phase (heteroreduction).^[2]

EXPERIMENTAL. All the experiments were performed in a N₂ glove box (O₂ < 0.5 ppm). The alumina used in the experiments has been previously characterized.^[3] For the batch sorption experiments, suspensions of alumina containing XCl₂ (60 μ M, where X = Mn or Sn) at different pH values have been in contact for seven days. Consecutively, the samples were centrifuged, the pH was measured and the supernatant sampled to determine the remaining X concentration in solution by ICP-MS. Then, the supernatant was discarded and the solid was dried inside the glove box. Once dried, the solid was re-suspended using a solution containing 5 μ M Tc^{VII}. After seven days in contact, the samples were centrifuged, the pH and redox potential (Eh) were measured and the supernatant was sampled for quantification of Tc concentration in solution by liquid scintillation counter.

The sorption yield of Y (where Y = Tc, Mn or Sn) was calculated as a distribution coefficient (K_d) using the equation:

$$K_{d_Y} = \left(\frac{\left[[Y]_0 - [Y]_t\right]}{[Y]_t} \cdot \frac{V}{m}\right) \tag{1}$$

where *V* is the volume of suspension in L, *m* is the mass in g, and $[Y]_0$ and $[Y]_t$ are the initial and final molar concentrations of Y, respectively.

RESULTS. The sorption of Mn^{2+} on alumina increases with increasing pH, being maximum for pH>9, whereas Sn^{2+} sorption on alumina is maximum for the entire range of pH evaluated and it only decreases at pH>9 (Fig. 1). Mn^{2+} follows the trend of cation adsorption on metal oxides, which is favored because of electrostatic attraction of the negatively charged alumina surface and the cation.^[3] Sn^{2+} is sorbed on the entire range of pH but the sorption mechanism is still to be determined as $Sn(OH)_2$ was not observed.

Once the solid pre-sorbed with X is dried, a solution containing Tc is added. Figure 2 shows K_{dTc} values when using alumina containing Sn pre-sorbed. They are one to two orders of magnitude higher than those by alumina with Mn presorbed. This is due to the fact that the oxidation reaction of Mn²⁺ to MnO₂ is not as favored as the oxidation of Sn²⁺ to Sn⁴⁺, thus a higher amount of Tc^{VII} remains in solution. However, at most of the evaluated pH values, the Tc sorption by Mn²⁺ pre-sorbed on alumina is enhanced compared to neat alumina, whose *Kd* maximum is 150 L/g.^[4] This suggests that when Mn²⁺ is sorbed on alumina, its redox potential might be modified to favor the reduction of Tc^{VII} to Tc^{IV}. Nevertheless, this hypothetized redox potential modification is not sufficient to produce a quantitative Tc removal from



Fig. 1: Sorption of Y (Y = Sn^{2+} or Mn^{2+}) on alumina in H₂O as a function of pH.



Fig. 2: Sorption of Tc promoted by Sn^{2+} or Mn^{2+} pre-sorbed on alumina in H_2O as a function of pH.

solution, as it happens when alumina containing Sn presorbed is used as sorbent.

With respect to our recent study of alumina pre-sorbed with $Fe^{2+}{}_{\prime}{}^{[4]}$ the overall Tc heteroreduction in the presence of different reductants follows this order: $Sn^{2+} > Fe^{2+} > Mn^{2+}$. Further investigations are in progress to evaluate the advantages of using Sn as Tc heteroreductant instead of homoreductant.

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New insights on the reductive immobilization of ⁹⁹Tc(VII) by iron sulfide

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The reductive immobilization of ⁹⁹Tc(VII) by pure pyrite and a synthetic mixture marcasite-pyrite 60:40 was studied by batch experiments, Raman microscopy and X-ray photoelectron spectroscopy (XPS). It was found that the presence of marcasite makes the retention slower and incomplete at pH10 due to the change of redox active species with respect to pyrite: S²⁻ instead of Fe²⁺.

Fe(II) minerals have shown a remarkable ability to immobilize ${}^{99}\text{Tc}(\text{VII})O_4^-$, a fission product of significant environmental concern due to its high solubility, by reducing it to Tc(IV).^[1-3] Among them, ubiquitous pyrite (orthorhombic FeS₂) was identified as an excellent Tc scavenger in a broad range of pH.^[4] However, even though there is strong evidence that polymorphism affects the retention of pollutants by minerals,^[5,6] no studies of the effect of marcasite (cubic FeS₂) of the Tc uptake by pyrite have been performed up to date.

EXPERIMENTAL. The pure pyrite was purchased from Alfa Aesar; the mixture marcasite-pyrite 60:40 (synthetic FeS₂) was synthetized following the procedure described by Hou *et al.*^[7] Both solids were characterized by X-ray diffraction. The batch sorption experiments were carried out inside a glovebox (< 1 ppm O₂) where Tc was left in contact with aqueous suspensions of either pure pyrite or synthetic FeS₂ (1.3 gL⁻¹) under constant agitation. Experiments were performed as a function of pH, time, ionic strength and Tc concentration. After certain time the solids were separated by ultracentrifugation ($2.4 \times 10^5 \times g$ for 1 h) and 0.25 mL aliquots of the supernatant were analyzed by liquid scintillation counting to determine the percentage of Tc removed. Solids loaded with 1,000 ppm Tc at pH6.0 and 10.0 were analyzed by XPS and Raman microscopy.

RESULTS. Figure 1 shows the pH effect on the technetium removal by pure pyrite and by the synthetic FeS₂ after 7 days in contact. While pyrite removes 100 % of Tc at pH \geq 5.5, in the presence of marcasite complete uptake takes place at 6.0 < pH \leq 9.0, being incomplete at pH10.0. At pH \leq 5.0, the Tc retention is incomplete in both solids due to the higher solubility of FeS₂ under acidic conditions.^[8]

The kinetics of the Tc-uptake by pure pyrite and by the synthetic FeS_2 at pH6.0 is presented in Fig.2. For pure pyrite complete Tc removal is reached within one day, whereas marcasite slows down the process to one week.

Figure 3 shows the Fe 2p XPS spectra of the Tc loaded iron sulfides at pH6.0 and 10.0. A peak around 711 eV demon-



Fig. 1: pH effect on the Tc removal by pyrite and synthetic FeS₂ after 7 days.



Fig. 2: Kinetics the Tc removal by pure pyrite and synthetic FeS₂ at pH 6.0.



Fig. 3: Fe 2p XPS spectra of the Tc loaded iron sulfides. [Tc] = 1,000 ppm.

strates that at pH6.0 both solids trigger the formation of Fe(III) after the reduction of Tc(VII). This new Fe(III) phase was characterized as hematite by Raman microscopy (not shown). At pH10.0, the Fe 2p spectra show no formation of Fe(III) in the presence of marcasite, suggesting that it is not the reducing agent for Tc(VII). The analysis of the Tc loaded synthetic FeS₂ at pH10.0 by Raman microscopy revealed the formation of sulfate minerals on its surface, revealing that at this pH the reduction of Tc(VII) is a result of the oxidation of S²⁻ instead of Fe²⁺.

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Heterogeneous sorption of radionuclides predicted by crystal surface nanoroughness

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Reactive transport modeling (RTM) is an essential tool for the prediction of contaminants' behavior in the bioand geosphere. However, RTM of sorption reactions is constrained via two parameters, i.e., specific surface area (SSA) and specific site density (SSD). Their variability is responsible for heterogeneous sorption efficiency, not covered in current RTM approaches. Here, we study the spatially heterogeneous sorption behavior of Eu(III), as an analogue to trivalent actinides. We utilize polycrystalline calcite surfaces showing surface nanoroughness variability. We analyze the sorption efficiency as a function of surface nanoroughness. Based on experimental data from micro-focus time-resolved laser-induced luminescence spectroscopy (µTRLFS), vertical scanning interferometry (VSI), and electron back-scattering diffraction (EBSD), we re-parameterize a surface complexation model (SCM) based on nanotopographic data. Model validation of the quantitatively predicted spatial sorption heterogeneity suggests that retention reactions can be considerably controlled by nanotopographic surface features. Our study presents a way to implement heterogeneous surface reactivity into a predictive SCM for enhanced prediction of radionuclide retention.

EXPERIMENTAL. A polycrystalline marble sample surface was analyzed with EBSD to identify the crystallographic orientation of the crystal grains. The surface topography was then analyzed using VSI after a dissolution experiment. Subsequently, a batch europium sorption experiment was conducted by immersing the sample surface into a solution of 10^{-5} mol/L Eu(NO₃)₃ and 0.1 mol/L NaCl with a pH value of 8.5 for 24 hours. The sample was then rinsed with the background solution and dried. Spatially resolved information on the Eu(III) sorption capacity and the adsorbed species was gained using μ TRLFS mapping with a spatial resolution of 20 μ m.^[1]

MODELING. The surface complexation model (SCM) was developed to describe Eu(III) adsorption on the calcite surface. Published site density and specific surface area (SSA) are 5 sites/nm² and 0.262 m²/g, respectively.^[2] The chemical reactions were calculated using PHREEQC v3.5 with the Nagra/PSI Chemical Thermodynamic Database. The SCM results were validated using measured sorption efficiency data from the literature.^[2]

RESULTS. Figure 1 shows an overview of the investigated calcite crystal surface maps including (a) the surface topography, (b) crystal orientations of the polycrystalline calcite material, and (c) normalized sorption capacity including selected spots where fluorescence lifetimes were collected. The luminescence intensity reveals a heterogeneous distribution of the Eu(III) sorption, even within single crystal faces, varying by a factor of around 10 between lowest and highest concentrations. Several grains show statistically significant differences in sorption capacity, however, no correlation between crystal orientation and sorption capacity is found. Therefore, other factors must be responsible for the variable surface sorption capacity observed by µTRLFS. We explored surface nanotopography, represented by surface nanoroughness (Sq), as a potential controlling factor of the surface energy distribution.^[3] Based on analysis of the measured data from four different grains, a weak linear correlation



Fig. 1: Correlative microscopy of surface topography, field-of-view (FOV) size = 1×0.4 mm; ROIs I-II refer to two representative sub-regions for model validation (a). Crystal orientation mapping acquired by EBSD (b). Sorption capacity (normalized $F_1 + F_2 \mu$ TRLFS peak sum). Spatial resolution: $20 \times 20 \mu$ m (c).

 $(R^2 = 0.34)$ between the calculated Sq values (FOV: $2.5 \times 2.5 \mu m$) and sorption capacity is identified. The quantitative relationship between the Eu sorption efficiency and specific surface area (SSA) by SCM allows us to quantitatively link SSA and Sq pixel by pixel by cross-referencing the experimental sorption capacity data and the model result. Therefore, we utilize this linear trend between SSA and Sq to parameterize a predictive SCM. A quantitative comparison of model results (Fig. 2a) with the measured data (Fig. 2b) in ROIs I and II shows good agreement within a 20% discrepancy between model and experiment (Fig. 2c). In contrast to typical SCM approaches that employ a simple uniform SSA value, this modified SSA based on Sq analysis allows for the investigation of spatially resolved retention variability based on the variability of crystal surface energy.





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PET imaging of actinide analogue transport and retention through fractured crystalline rock

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The migration of actinides determines the long-term safety of a final deposit of nuclear waste. This migration is influenced by mineralogical, topographical and hydrodynamic heterogeneities. To investigate the influence of flow heterogeneities at the core scale we performed Positron Emission Tomography (PET) imaging studies using Y-86 as a PET-traceable actinide analogue radiotracer.

EXPERIMENTAL. A gneiss drill core was procured from the Bukov Underground Laboratory in the Czech Republic. Two column samples were prepared by encasing parts of the core in epoxy resin. The two samples consisted of a natural fracture and an artificial fracture that was induced by a mechanical stress test. X-ray tomography (CT) images were taken to characterize the geometry of the fractures. First, for PET imaging, a conservative tracer experiment was performed using 10 mM [F-18]KF in aqueous 0.1 M NaCl at pH7.5. This was followed by a reactive experiment using 10-5 M [Y-86]Y(NO₃)₃ in aqueous 0.1M NaCl at pH7.5. Y-86 and F-18 radiotracers (Tab. 1) were produced using the in-house cyclotron at Research Site Leipzig. For Y-86 a target of enriched 86SrCO3 was proton irradiated to trigger the nuclear reaction Sr-86(p,n)Y-86. The target material was dissolved and the radiotracer was separated by ion chromatography using LN-resin. F-18 was produced from O-18 enriched water.

Tab. 1: Radiotracer properties.

Radiotracer	Halflife [h]	β+ [%]
F-18	1.828	96.7
Y-86	14.7	33

RESULTS. The conservative PET imaging of the natural fracture showed a rather homogeneous flow.^[1] However, upon repetition of the experiment a pronounced heterogeneity of the flow path was observed. Changes in the fracture between the experiments now lead to two diverging flow paths right after the entrance to the fracture. The flow was diverted around a middle section that was not accessible anymore. In the following reactive experiment using Y-86 most of the actinide analogue was deposited along one of the two divergent paths of the second F-18 experiment (Fig. 1). The bulk of the Yttrium is sorbed on the fracture surface and only 1% is eluted. The distribution along the fracture can be fitted with a power law indicating an averaging of the influence of the mineralogical heterogeneity below the cm scale of the core.

For the artificial fracture, the F-18 experiment revealed a rather homogenous flow field, which nevertheless seems to be split in half due to a low-permeability surface topography feature in the fracture center right after the entrance. In the subsequent Y-86 transport experiment, it was observed that the bulk of the Yttrium got diverted to one side of the fracture taking only one of the divergent flow paths (Fig. 2). Furthermore, in the artificial fracture two distinct sorption hotspots could be observed that are related to mineralogical features. The mineralogical heterogeneity of the fracture surface is on a slightly higher scale than for the natural fracture, for which a levelling effect exists due to natural fracture filling material. In the artificial fracture, the influence of the



Fig. 1: Overlay of Y-86 sorption pattern on the gneiss fracture surface from PET (red-blue) and the conservative flow path measured with F-18 (blue-green-yellow). The Y-86 largely is deposited along one of the divergent flow paths seen in the F-18 experiment. The middle section of the fracture is not accessed by the flow.



Fig. 2: Conservative flow path from PET process tomography (left) and Y-86 sorption pattern from PET (right). The conservative flow field is split in two by a central low-permeability feature. The Y-86 sorption pattern reveals a diversion of the flow to one side of the fracture with two distinct hotspots observable.

mineralogical heterogeneity overlaps with the influence of the heterogeneous flow path.

Our experiments show the impact on actinide migration by heterogeneous flow patterns. Future work will focus on the impact of fracture filling materials on the flow path and more detailed investigations into mineralogy and actinide speciation.

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Site-specific crystal surface reactivity

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The variability of crystal surface reactivity under identical chemical conditions is explained by significant intrinsic rate variability. This is quantified using rate maps of reacting surfaces and rate spectra as a statistical concept. We discuss the density and distribution of surface kink sites as the crucial parameter controlling the observed rate variability. For the first time, we present surface rate data that quantify simultaneously the rate contributions of specific surface portions of crystals, *i. e.*, facets, edges, and corners.^[1] Resulting rate ranges could help improving the parametrization of reactive transport models at the pore scale.

Recent studies analyzed the rate variability and utilized rate maps and rate spectra as a means of quantifying multiple reaction rate portions.^[2] Specifically, surface-sensitive methods with high spatial resolution, such as Atomic Force Microscopy (AFM), Vertical Scanning Interferometry (VSI), and confocal microscopy are used as complementary techniques for studying reacting single crystal surfaces. The results suggested the impact of crystallographic orientation, defect density, and exposure time on measured dissolution rates. Such investigations complement basic experimental and theoretical studies focusing on the reactivity of surface steps and pits. Surface kinks are discussed to be the major contributor to elevated surface rates. A comparison of surface portions with contrasting kink site densities (KSD) would inform quantitatively about specific rate ranges as a function of KSD (facets: low, edges: enhanced, corners: highest KSD).[1]

EXPERIMENTAL. We utilized a Sensofar Plu S neox microscope combining interferometer and confocal surface analysis techniques. For the rate data measurement, we analyzed the calcite crystal surface using two confocal objectives (150 × and 20 × magnification). The sample surface topography was analyzed before and after each dissolution period. Details about the dissolution experiment are available in literature.^[11] Each measurement included a masked and a reacted surface portion. The measured datasets were processed with the software SPIP (version: 6.7.3; Image Metrology A/S, Hørsholm, Denmark). The velocity of height changes (dz/dt) equals the material flux at each point of the surface. Rate maps were calculated by dividing each dz/dt-value by the molar volume of calcite.

RESULTS. Figure 1 shows the topography of a calcite single crystal after five minutes of dissolution in a 2.2 mM Na₂CO₃ solution at pH8.8, see background height map with ambercolor scheme. The inset at the top of Fig. 1 illustrates the elevated dissolution rates at the crystal corner and edges, compared to facets. The analysis of rate data in the frequency domain (Fig. 2, left) shows a rate distribution covering about 2.5 orders of magnitude. The comparison to polycrystalline calcite dissolution rate data under identical chemical conditions illustrates the fundamental differences of the rate frequency distributions. Note the remarkably high rates at the corner of the single crystal, explained by the KSD scheme (Fig. 2, right side).

CONCLUSIONS. We categorized dissolution rate data based on kink site density, based on multiple literature data. The compilation suggests an explanation of the complex shape of multimodal rate spectra based on surface features



Fig. 1: Topography of a single calcite crystal with facets, edges and corner (background height map, amber-color scheme). Multiple measurements of the surface topography using interferometry microscopy are utilized to calculate a dissolution rate map of the specific surface building blocks (upper part).



Fig. 2: Dissolution rate distribution of calcite single crystal (purple) vs. polycrystalline material (red) (left side). Schematic illustration of KSD increase (bottom-top) and explanation of rate increase, central part. Categories of KSDs (i–iv) represent simple endmember situations while measured rate spectra result from a complex superimposition of such endmembers. Right: Simplified mechanistic explanation of the KSD variability of crystal surfaces.

and associated kink site densities. We suggest an application of this insight for the parametrization of pore scale models in reactive transport codes.

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Effective diffusivity prediction of radionuclides in clay formations using an integrated upscaling workflow

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Accurate estimation of radionuclides migration in the host rock using numerical tools plays a key role in the safety assessment of disposal concepts for nuclear waste. Current diffusive transport modeling applies a simplified parameterization of the effective diffusivity for simulating and predicting the radionuclides migration in the host rock such as the Opalinus clay (OPA). This prevents to model the heterogeneous diffusion behavior as indicated by the Positron Emission Tomography (GeoPET) measurements.^[1] In this study, we propose an upscaling workflow that integrates nm- and µmscale simulations based on synthetic multiscale digital rocks to predict the effective diffusivity of radionuclides in the OPA. The proposed numerical approach is validated with published experimental data confirming general applicability.^[2] In this approach, the spatial variabilities of the pore network and compositional heterogeneities at the pore scale are now addressed and available for studying heterogeneous diffusion patterns compared to commonly assumed homogeneous behavior.

MODELING. The effective diffusivities of radionuclides are determined by digital rock physics. The respective synthetic digital rocks include two types of materials for studying diffusive transport. Clay-rich rock sections are the starting point, commonly referred to as the clay matrix, which contains sheet silicate crystals and pores at the nm-scale (Fig. 1a). At this scale, we consider the surface diffusion for the cations and anions transport within the diffuse layer (DL), and the surface electrostatic potential is calculated using the Donnan approach in PHREEQC v3.5. Second, we consider non-clay minerals, such as quartz grains and diagenetic calcite crystals, as non-porous components, thus including additional compositional complexity at the µm-scale (Fig. 1b). At each length scale, the diffusion-sorption equation is numerically solved by a previously developed numerical simulator.[3]

UPSCALING. The suggested upscaling workflow integrates nm- and μ m-scale simulations, illustrated in Fig.2. Detailed workflow steps include: (1) The diffusivity of each tracer in



Fig. 1: Schematic representation (I) and digital rock model (II) of the partial rock structure built-up of sheet silicate crystals and pores (black). Black color represents pores and white color represents the clay crystals (a). (I) SEM image of the microstructure of OPA, indicating clay, quartz, calcite, and pyrite minerals. (II) Reconstructed digital rock with clay mineral aggregates (gray color) and other sedimentary components (red color) including quartz, calcite, and pyrite (b).

free water D_W and the geometrical model of clay matrix are used to determine the effective diffusivity D_e of the clay matrix at the nm scale. (2) D_e of clay matrix and geometrical model of microstructures of shales are used to determine the D_e of shale at the µm-scale. In this upscaling workflow, both small-scale information such as pore space and pore network geometry and larger-scale composition are considered.



Fig. 2: Flow chart of the proposed upscaling workflow from the nm- to μmscale. This includes diffusive flow in (1) porous clay matrix and (2) single laminae of varying composition.

RESULTS We determine the effective diffusivities of ²²Na⁺, HTO, and I⁻ in the shaly facies of OPA using workflow steps (1) and (2), as shown in Fig. 2. The synthetic digital rock of shale has a total porosity of 0.158 and is composed of 66% clay minerals,[4] containing elongated pores.[5] The calculated mean values of effective diffusivities of ²²Na⁺, HTO, and I⁻ by the model and the experimentally measured values from literature are listed in Tab.2. Good agreement between the diffusivity from the literature and our suggested upscaling workflow underscores the capability of predicting the diffusivity of the shales.^[2] Due to the considerations of both small-scale information such as pore space and pore network as well as larger-scale compositional information in the geometrical models, our concept of upscaling allows us to model heterogeneous diffusivity caused by the spatial variability of pore networks in complex rocks as well as the compositional heterogeneities in shales. The suggested upscaling strategy is applicable to provide insight and input data to simulations at the core scale and above, thus contributing to enhanced predictability of radionuclide migration.

Tab. 1: Effective diffusion coefficients of ²²Na⁺, HTO, and I[−] in shale from experimental data in the literature compared to the predicted mean values from the model results.^[2]

T	Diffusivity of sha	ale (10 ⁻¹¹ m²/s)	
Tracer	Experimental data ^[2]	Model results	
²² Na ⁺	7.2-9.4	7.73	
HTO	5.4-6.5	6.23	
I-	1.2-2.0	2.08	

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Combined µCT and MIP porosity measurements for identification of transport pathway distributions in salt rocks

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Potential fluid transport pathways through rock material in close proximity to a barrier material by combining Mercury Intrusion Porosimetry (MIP) and Micro-Computed Tomography (µCT) on unconfined samples were analyzed. The MIP data generally suggest the occurrence of a small portion of tight percolation pathways of the bulk material. However, apparent low-pressure intrusion raises concern that additional conductive structures with a size around 0.1 mm could be present as well. These large structures could be attributed to the surface roughness of the samples by μ CT data analysis. They are without consequences for the barrier function. In contrast, the largest detectable fracture structures of the bulk material are in the micrometer range and display low density contrast - either due to salt precipitates or residual brine.

Salt rocks are a prospective formation for hazardous and radioactive waste disposal owing to their extremely low permeability and their plasticity. Underground repositories require reliable technical barriers in close contact with the formation rock. Both elucidation of flow processes and assessment of injection concepts for technical improvements of the barrier are key issues.

EXPERIMENTAL. Small cubes with edge length of 20 mm were cut from salt rock drill cores by dry sawing. These samples were inspected with the Nikon XT-H-225 μ CT-scanner (U = 190 kV, I = 90 μ A, Filter: 0.5 mm Cu, voxel size: 12.4 μ m).

Elaborate background corrections and filtering were conducted prior to segmentation of detectable fractures with low contrast, because the images were prone to background variations caused by the square sample shape. After top-hat segmentation (Fig. 1), the local void sizes were computed with the thickness map algorithm of the Avizo software (Fisher Scientific). The outer surface roughness was estimated as the thickness map of the rough sample volume and its smooth external envelope.

MIP was conducted with AutoPore V (Micromeritics), which features an extended low-pressure range, in order to quanti-



Fig. 1: μCT, showing saw-cut surface roughness (bottom); exposed thickness map of fracture system, color-coded (top); histogram of the thickness map (inset).

fy large pores up to D = 0.8 mm. The smallest detectable pore size is D = 3 nm, according to a maximum injection pressure of 400 MPa.

RESULTS. The MIP data show a considerable portion of almost millimeter-scaled voids that are suspected to provide potential flow pathways. Using μ CT measurements, the existence of such pore structures is excluded. Instead, μ CT suggests roughness of the sample surface as a source of potentially misleading MIP data interpretation (Fig. 2, negative range). However, the lack of a well-defined smooth reference surface limits the quantitative comparability of MIP and μ CT with respect to large classes of the pore spectrum. We thus consider the size (D = 0.3 mm) of the largest pores detected by μ CT as the maximum pore size existing in the bulk material.



Fig. 2: Pore size distribution of a rock-salt sample. MIP in blueish colors (low pressure injection, high pressure injection, and decompression), μCT in reddish colors (1. surface roughness, 2. empty voids in the bulk volume, 3. voids with assumed porosity of 0.3). Negative MIP-volume values represent surface roughness, corrected by μCT.

In general, such voids in salt rock show the lowest density and are thus most probably air-filled. Interestingly, smaller voids show frequently considerably higher density compared to the larger ones. Most probably, such structures are not open, but contain portions of dense material. Consequently, these complex pore structures are considered as porous zones, either not detectable by MIP (solid fill, *e.g.*, brines) or correctly attributed to smaller pore classes (partly filled by microporous material, tentative porosity = 0.3, Fig. 2). Overall, the combined use of μ CT and MIP provides critical insight into the complexity of pore size distributions.

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PET transport experiments under confining pressure

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A new radiation-transparent vessel for the conduction of flow experiments under hydrostatic confining pressure particular suitable for Positron Emission Tomography (PET) and X-ray tomography (μ CT) was designed. First promising results are presented.

It is common best practice in core analysis for reservoir evaluation to conduct flow experiments under hydrostatic confining pressure in a pressure vessel. With this procedure, *in situ* pressure conditions can be investigated. Second, an improved sealing between sample and containment is ensured, thus avoiding flow artifacts. Nevertheless, current process-tomographic experiments were lacking this vital feature. Commonly used steel pressure vessels cause strong scatter and attenuation of radiation and thus interfere with both tomographic methods, PET and μ CT. To overcome these issues, we designed and successfully tested a plastic-material pressure vessel. This enables to conduct more realistic transport experiments. Additionally, this new setup allows for the analysis of reactive transport processes affecting shape and mechanical properties of the sample.

EXPERIMENTAL. The pressure vessel was lathed from polyoxymethylene copolymere (POM-C). It contains the mounted sample that is sealed with shrinking tube. (Fig. 1). The carrier medium for the confining pressure (p_c) is water, which can be pressurized with a pressure-controlled syringe pump or a gas pressure storage container (He). The thickness of the vessel wall is designed to withstand pressures up to 10 MPa, although the intended use is 2 MPa. The design of the sealing can be adapted to the sample size. The maximum diameter is limited to D = 120 mm, according to the maximum field-of-view of the scanners. The device has to fit entirely into the PET-scanner in order to include its scatter and attenuation effects into the image reconstruction process. Consequently, this limits the sample diameter to *ca.* 50 mm and the length to 100 mm.

Usually, our flow experiments are conducted with a slow constant flow rate and a driving pressure delimited to $p_d = P_c/2$, in order to avoid superficial flow and significant pore pressure effects. Within this pressure limit, the flow rate is adjusted to achieve a maximum flow velocity of *ca*. 60 mm/h. Although p_d can be reduced at will by decreasing the injection rate, a significant pressure gradient should be established to avoid too high sensibility to fluctuations.

RESULTS. Test experiments on a Rotliegend sandstone sample with a diameter of 50 mm were performed successfully (Fig. 2). Another important result is the capability of the applied PET data correction algorithms of removing attenuation and scatter artefacts and obtaining fair image quality. However, the first data analysis still shows increasing radial tracer concentration *versus* the sample surface (Fig. 2, left). The trough - shaped pattern indicates a disaggregated zone near the sample surface, but no superficial flow along a superficial fluid film. Additional experiments with varying confining pressure are planned, in order to study the flow behavior in such disaggregated zones.

In general, such pressure-dependent behavior of flow cells are usually not detectable without PET process tomography. Thus, we suggest PET measurements as the critical safeguard against misinterpretations of transport results collected by using pressure cells.



Fig. 1: Scheme of the radiation-transparent pressure vessel.



Fig. 2: First experiment on a Rotliegend sandstone sample in the new pressure vessel. Cross section of the PET-µCT overlay (µCT still without confining pressure) (left). Enlarged 3D-view of the flow vector field computed from the sequence of PET frames (right).

Production of ²²Na by the non-common nuclear reaction ^{nat}Mg(d, α) and purification of ²²Na for investigations with GeoPET

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A method for the production and purification of ²²Na ($T_{1/2}$ = 2.603 a, β^+ = 90.3 %), suitable for transport investigations in low-permeability systems, such as clay host rocks is introduced.

Opalinus clay is a possible host rock for nuclear waste repository systems. Currently, detailed studies are focusing on reactive transport properties of this complex geomaterial. GeoPET (Positron Emission Tomography optimized for geological samples) is an excellent method for the analysis of radionuclide transport and retention. Transport analysis requires a suitable non-reactive PET tracer, suitable for solutions mimicking the pore water chemistry of the Opalinus clay rock. Here, we present a method for the production and purification of ²²Na (T_{1/2} = 2.603 a, β^+ = 90.3%), suitable for transport investigations in low-permeability systems.

EXPERIMENTAL. ²²Na Opalinus water: A magnesium foil (diameter: 11 mm; weight: 79 mg) was irradiated at the cyclotron Cyclone 18/9 (IBA RadioPharma Solutions, Belgium) with 9MeV deuterons by using the nuclear reaction ^{nat}Mg(d, α)²²Na.^[1] As depicted in Fig. 1, the magnesium foil was dissolved in 6M HCl and the solution was evaporated to dryness. After dissolution in 0.1 M HCl the preconditioned cation exchanger Dowex 50W x 8 (0.1 M HCl) was used to fix both, magnesium and sodium on the column.^[2] The ²²Na was eluted by 0.3 M HCl, while magnesium remains on the column. The fractions were measured by γ -spectrometry. After evaporation of the combined ²²Na-frations, the ²²Na was dissolved in saturated NaCl solution or Opalinus water, depending on the flow experiment. The column was regenerated by washing two times with 6M HCl and preconditioned with 0.1M HCl for the next separation.

Diffusion experiment with GeoPET/CT: Sample material was a drillcore of Opalinus clay (sandy facies, length: 95 mm, diameter: 67 mm) with centered axial borehole (diameter: 3 mm). The axial borehole was filled with synthetic Opalinus



clay pore water, labelled with 0.55 MBq $^{22}Na^{\star}$. The GeoPET-image (Fig. 2) is the 9th from a suite of 32 frames (frame length: 30 min), acquired during a period of one month with a ClearPET scanner (Elysia-Raytest). The low resolution background μCT image was acquired with a Nikon XT-H-225 scanner for the purpose of PET attenuation and scatter correction.



Fig. 2: Overlay picture of GeoPET and CT scans visualizing the diffusion of synthetic Opalinus water in an Opalinus drillcore.

RESULTS. The magnesium target material was separated completely; no visible residue remained after evaporation of the ²²Na fractions. The mean radiochemical separation yield was $88 \pm 7\%$ for the stock solution. Using Opalinus water as stock solution no further sample manipulation is needed. The required chemical form of the tracer for the GeoPET studies with Opalinus clay could be precisely adjusted and is routinely provided for the GeoPET studies.

Fist applications were diffusion experiments with ²²Na Opalinus water in Opalinus clay. Due to its half-life of 2.603 a, ²²Na is an ideal PET tracer for such long-term diffusion studies. Figure 2 shows first results of diffusion experiments of the 0.5 mL sample within the water-saturated Opalinus clay drillcore using the combined GeoPET-CT approach. The overlay of diffusion patterns from GeoPET with the inhomogeneous fabric from μ CT provides new insights into the heterogeneous diffusion process as the dominant transport process to be considered in performance assessment of the barrier function of clay rocks.

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88 % ± 7 % radiochemical yield

Fig. 1: Separation scheme for the purification of $^{\rm 22}{\rm Na}$ from a deuteron irradiated $^{\rm nat}{\rm Mg}$ target.

First-time production of Nb-92m/95 at the cyclotron Cyclone 18/9 in Leipzig

K. Franke

The exit from nuclear energy implicates the dismantling of nuclear power plants in the near future. During this decommissioning, Nb-94 ($T_{1/2} = 112.03 \times 10^4$ a) has to be considered as a constituent of the radioactive waste in the dismantling process. This entails the need for suitable Nb-radiotracers to investigate experimentally the likely fate of Nb-94.

The half-life and decay mode suggest the isotope Nb-95 ($T_{1/2}$ = 34.991 d, E_{γ} = 765.80 keV/99.15%) as promising candidate.^[1] Nb-95 can be produced via a (d,x)-reaction with a Zr target with natural isotope abundance. The irradiation with 9 MeV deuterons results mainly in the generation of Y, Zr and Nb radionuclides.^[2] Y and Zr radionuclides can be easily separated from Nb-95 during target processing. The relevant cross sections for the formation of the Nb isotopes suggest the formation of Nb-90, Nb-91m, Nb-92m, Nb-95, Nb-95m and Nb-96 (Fig. 1).



Fig. 1: Cross section for the formation of Nb isotopes Nb-90, Nb-91m, Nb-92m, Nb-95, Nb-95m and Nb-96 via Zr-nat(d,x) reaction.^[2]

Most of the formed Nb isotopes can be neglected due to their half-life and their low γ -emission probability (Nb-90: $T_{1/2}$ = 14.6 h, Nb-91m: $T_{1/2}$ = 60.8 d, Nb-91: $T_{1/2}$ = 680 a, Nb-92: $T_{1/2}$ = 3.47 × 10⁷ a and Nb-94: $T_{1/2}$ = 2.03 × 10⁴ a). Only Nb-92m ($T_{1/2}$ = 10.15 d, E = 34.17 keV/99.15 %) has to be considered in the following experimental setup. $^{[1]}$ Luckily, Nb-92m and Nb-95 γ -emissions are easily distinguishable with common detection systems.

EXPERIMENTAL. *Cyclotron.* The cyclotron Cyclone 18/9 (IBA RadioPharma Solutions, Belgium) provides protons with an energy of 18 MeV and deuterons with an energy of 9 MeV. The maximal current on target is $100 \,\mu$ A for protons and $40 \,\mu$ A for deuterons.

Target station. Solid targets were irradiated at the Nirta[®] solid target system (IBA RadioPharma Solutions, Belgium) at port 3 of the cyclotron, mounted at a 2 m beam transfer line. The target is cooled with helium in front and water in the back-side.

Windows. $12.5 \,\mu$ m Ti foil was used as vacuum window. No target entrance window was used for the target to avoid further degradation of deuteron energy. The Zr-foil was cooled directly with He.

Target capsule. The target was hold in a target capsule (Fig. 2.). It consists of an Al-backing and open Al-cover.



Fig. 2: Irradiated target capsule. Beam spot indicates non-ideal focused beam in horizontal alignment.

Target. A metallic Zr foil (Alfa Aesar, diameter: 12 mm, thickness: $127 \mu \text{m}$) with natural isotopic composition was used as target material.

Irradiation. Irradiation was done with deuterons with an initial energy of 9 MeV and a current on target of $10 \,\mu$ A for $550 \,\mu$ Ah.

 γ -spectrometry. An HPGe-detector ORTEC GEM15180-P combined with signal processing module DSPEC 502 (ORTEC) was used for γ -spectrometry measurements.

RESULTS. The target capsule was opened 4d after end of bombardment to allow short-lived radionuclides to decay. The target foil showed an oval burning mark (Fig.2). Although this is not critical because of the low deuteron current and the high melting point of the Zr-foil further optimization of the beam focus is needed for a more homogeneous target irradiation. The irradiated Zr-foil was unloaded and transferred to γ -spectrometry. The spectrum confirms quantitatively the yield of Nb-92m (110 MBq), Nb-95 (27 MBq), Nb-95m (18MBq) and Zr-95 (21MBq) (Fig. 3). Subsequent chemical Zr/Nb separation combined with further decay of Nb-92m result in the required Nb-95. First applications at the Institute for Nuclear Waste Disposal, Karlsruhe Institute of Technology, Germany, include the investigation of the impact of isosaccharinic acid and chloride on the uptake of Nb(V) by cement (CEM I HCP).[3]



Fig. 3: Section of the γ-spectrum of the irradiated Zr foil. Prominent peaks of Y-88 (898.04 keV), Zr-95 (724.19 keV, 756.72 keV), Nb-92m (912.60 keV, 934.44 keV) and Nb-95 (765.80 keV).

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Metal dissociation from humic colloids: kinetics with time-dependent rate constants

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The mobility of radionuclides in aqueous subsurface environments is essentially governed by their interaction with humic or humic-like substances as colloidal constituents of Dissolved Organic Matter or highmolecular-weight organic components in clay. Dissociation is a slow kinetic process, which is even more impeded with increasing time of contact due to structural rearrangements. Based on results obtained in isotope exchange experiments, the convoluted time dependence of dissociation was fully described by integrating timedependent rate "constants" to address the phenomenon of kinetics within kinetics.

Growing inertness of complexes with Natural Organic Matter was observed for a number of multivalent metals, such as Eu(III), Am(III), Th(IV) and U(VI).^[1,2] Based on thermodynamic constants, conventional transport models are, however, assuming steady local equilibria. To include kinetically hindered processes, their time dependence must be understood and parameterized. Solving the kinetic problem of convoluted reaction rates was the objective of this study.^[3]

EXPERIMENTAL. The kinetic treatment refers to isotope exchange experiments performed with [160Tb]Tb(III) and humic acid, delineated in Fig. 1.[4]



Fig. 1: Schematic representation of isotope exchange experiments starting from a low metal load of humic acid (HA) with the radioisotope (M*) for variable times of pre-equilibration ($t_{\rm pre}$), followed by saturation with the stable isotope (M), which initiates partial desorption observed over variable times (t_{des}). Double circles indicate inertization.

RESULTS. As shown in Fig. 2, desorption of [160Tb]Tb as a radiotracer (M*) was found to be very slow, and moreover, the rate was dependent on the time of pre-equilibration with humic acid prior to addition of the stable isotope ¹⁵⁹Tb (M).Thus, an inertization takes place. Each data series is in turn affected by this process, since it is also active during the slow liberation of [160Tb]Tb.



Fig. 2: Release of [¹⁶⁰Tb]Tb from humic acid (HA) initiated by saturation with non-radioactive Tb (pH 4.0, $0.5\,g\,L^{-1}\,HA,\,1\,nM\,[^{160}Tb]Tb,\,2\,mM\,Tb,$ 0.1 M NaClO₄).

For an overall kinetic description of desorption superimposed by inertization, the first-order rate law for approach to equilibrium regarding the concentration of humic-bound $[^{160}\text{Tb}]\text{Tb}(c_{ads})$ was integrated including a time dependence of the rate constant *k*, given by an empirical function:

$$-\int_{c_{ads}^{c_{ads}}}^{c_{ads}} \frac{1}{c_{ads} - c_{ads}^{eq}}} dc_{ads,1} = \int_{t_{des}=0}^{t_{des}} k(t) dt_{des}$$
(1)

with
$$k(t) = a(t_{pre} + t_{des})^b + k_{\infty}$$
 (2)

The indices ini and eq refer to the initial and equilibrium state, respectively. The time t controlling k includes the preequilibration period. Beside the parameters a and b, an offset value k_{∞} for infinite time was added because k does not necessarily decline to zero. With respect to the very fast desorption at the initial stage, the total amount of sites occupied by [160Tb]Tb was split into a fast (index 1) and a slowly (index 2) desorbing part with the fractions x and (1-x), respectively. Integration yields:

$$c_{ads} = c_{ads}^{eq} + x(c_{ads}^{ini} - c_{ads}^{eq}) \exp\left\{\frac{a_1}{b_1 - 1} \left[\left(t_{pre} + t_{des}\right)^{1 - b_1} - t_{pre}^{1 - b_1} \right] \right\}$$
$$\exp\left(-k_{\infty,1} t_{des}\right) + (1 - x) \left(c_{ads}^{ini} - c_{ads}^{eq}\right)$$
$$\exp\left\{\frac{a_2}{b_2 - 1} \left[\left(t_{pre} + t_{des}\right)^{1 - b_2} - t_{pre}^{1 - b_2} \right] \right\} \exp\left(-k_{\infty,2} t_{des}\right). \tag{3}$$

With this equation, all time series can be well represented by one set of parameter values (line fits in Fig.2), without assuming a metal transfer from fast sites to slow sites. The extracted time functions for k_1 und k_2 are shown in Fig. 3.



Fig. 3: Inertization process as reflected by decreasing rate constants.

CONCLUSIONS. This analysis, breaking the rules of classical kinetics, showed that the inertization process is limited in time and extent. Thus, for modeling speciation or migration in nearly stagnant systems on a long-time scale, thermodynamic concepts using equilibrium constants remain applicable. However, the kinetic effect must be considered in any attempt to predict organically bound metal fractions from the composition of aquatic systems (and vice versa) if conditions change relatively fast.

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Stone-Wales defects cause high proton permeability and isotope selectivity of single-layer graphene

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While the isotope-dependent hydrogen permeability of graphene membranes at ambient condition has been demonstrated,^[1,2] the underlying mechanism has been controversially discussed during the past 5 years. The reported room-temperature proton-over-deuteron (H+over-D⁺) selectivity is 10, much higher than in any competing method. Yet, it has not been understood how protons can penetrate through graphene membranes - proposed hypotheses include carbon vacancies and local hydrogenation. However, neither can explain both the high permeability and high selectivity of the atomically thin membranes. Here, it is confirmed that ideal graphene is quasi-impermeable to protons, yet the most common defect in sp² carbons, the topological Stone-Wales defect (5- and 7-membered rings formed from 6-membered rings), has a calculated penetration barrier below 1 eV and H*-over-D* selectivity of 7 at room temperature and, thus, explains all experimental results on graphene membranes that are available to date. This is a topological defect and not a vacancy, thus, it would not be observed in typical experiments, e.g., using a nanoballoon. The competing explanation, local hydrogenation, which also reduces the penetration barrier, but shows significantly lower isotope selectivity, is challenged.

In this work,^[3] we calculated the flow of protons and their isotopes through graphene membranes and through the most common graphene defect sites to explain experimental results of graphene permeability and selectivity.^[1,2] The models include, for comparison, pristine *h*BN and, as principal target material, graphene as pristine membrane, with the topological SW 55–77 defect (with two 5- and two 7-membered rings - 5MR and 7MR) and with a locally hydrogenated ring (Fig. 1). Note that SW 55-77 defect always appear as pairs of the two ring types, which is formed from by rotation of one C-C bond by 90°. The flow of protons (and their isotopes) was estimated including quantum tunneling via the Wentzel–Kramers–Brillouin (WKB) approximation.

EXPERIMENTAL. All structures were optimized using DFT, employing the PBE0 functional and 6–31+G(d,p) basis set. The D3-(BJ) correction of London dispersion interactions following the approach of Grimme was used as implemented in Gaussian 09 package. The transfer energy barriers (U_{max}) of a single proton going through different types of rings in the cluster and perpendicular to the flake plane were calculated using single point energies on the optimized cluster structures. Quantum tunneling was considered by applying the WKB approximation. See [3] for detailed methods description.

RESULTS. While we provide data for all three isotopes of hydrogen, including tritium for comparison, only the selectivity for the separation of H⁺ and D⁺ can be compared to the available experimental data and is, thus, the focus here. We show that the permeability of protons through perfect graphene is too low to generate appreciable particle flow, even though it would result in a very high selectivity (much high-

er than reported in the experiment). Local hydrogenation significantly lowers U_{max} , but also reduces significantly the isotope selectivity to values much lower (around 2) than those reported from the experiment. On the other hand, we argue that the topological, vacancy-free SW 55-77 defect, with U_{max} below 1 eV and H⁺-over-D⁺ (H⁺/D⁺) selectivity of about 7, matches well the experimental findings of about 10 by Lozada-Hidalgo et al.^[2] The H⁺ flow through the 7MR is 10⁶ times larger than that through the 6MR, indicating that, even if the concentration of such a defect was as low as 1 ppm, it would still govern the transport process in the graphene flake (and yet neither be detectable in the Raman spectrum nor in the nanoballoon test). These findings indicate that increasing the concentration of 7MRs in vacancydefect-free sp² carbon structures are a promising route for the design of hydrogen-isotope-separation and potentially proton-exchange membranes. It is worth noting that the experimentally detected H+/D+ selectivity of 10 was also obtained for hBN monolayers.^[1,2] Since hBN has normally much less defects, we have calculated the energy barrier and selectivity using pristine model with 6MR. We again obtained H⁺/D⁺ selectivity of 7, in good agreement with the experiment. The summary of the finding of these studies is shown in Fig.1.



Fig. 1: Schematic representation and summary of proton (H⁺) and deuteron (D⁺) transfer and separation through graphene with different carbon ring sizes. The most plausible explanation of the high isotope selectivity is that seven-membered rings are present in the 2D membrane (green), as in the case of common SW 55–77 defects.

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Size dependence of lattice parameter and electronic structure in CeO₂

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Intrinsic properties of a compound, including electronic structure, crystallographic structure, optical and magnetic properties, define notably its chemical and physical behavior. In the case of nanomaterials, these fundamental properties depend on the occurrence of quantum mechanical size effects and on the considerable increase of the surface to bulk ratio. Here, we explore the sizedependence of both crystal and electronic properties of CeO₂ nanoparticles (NPs) with different sizes by stateof-the art spectroscopic techniques. X-Ray Diffraction (XRD), X-Ray Photoelectron Spectroscopy (XPS), and High-Energy-Resolution Fluorescence-Detected X-ray Absorption Near-Edge Structure (HERFD-XANES) demonstrate that the as-synthesized NPs crystallize in the fluorite structure and they are predominantly composed of Ce^{IV} ions. The strong dependence of the lattice parameter with the NPs size was attributed to the presence of adsorbed species at the NPs surface based on FTIR and TGA measurements.^[1]

EXPERIMENTAL. Ceria nanoparticles samples were synthesized at JRC Karlsruhe by alkaline precipitation of cerium ammonium sulfate precursor followed by hydrothermal condensation in a pressurized autoclave at different temperatures. HERDF-XANES measurements were conducted at the CAT-ACT beamline (ACT station) of the synchrotron light source (KIT, Karlsruhe, Germany).

RESULTS. The oxidation state and electronic structure of Ce was assessed using Ce L_{III} edge HERFD-XANES (Fig. 1) and XPS. No characteristic band feature of Ce^{III} was found for samples across all crystallite sizes, indicating that only tetravalent cerium is present.

XRD and transmission electron microscopy data show that the as-synthesized CeO₂ NPs crystallize in the Fm–3m fluorite structure (space group 225). Depending on the experimental conditions (and particularly the annealing temperature), the average crystallite diameters vary from 2.0 ± 0.1 to 91 ± 12 nm.



Fig. 1: Ce L₃ HERFD-XANES spectra of CeO₂ NPs with different sizes compared to bulk CeO₂ (> 100 nm).



Fig. 2: Decrease of lattice parameter with the particle size.

Figure 2 presents the size dependence of the lattice parameters of the as-synthesized CeO₂ NPs. When the particle size changes from 2.0 to 91 nm, the unit cell value varies from 5.456 (3) to 5.411 (1) Å.

This lattice parameter variation as a function of the particle size has been ascribed to two mechanisms: either to the formation of oxygen vacancies in the smaller NPs, which leads to a change in the oxidation state of the constituting cation, or to the variation of the surface stress resulting from the difference of coordination between atoms on the surface and in the bulk.^[2] This effect becomes more pronounced as the particle size reduces, *i.e.* as the contribution of the surface atoms to the structural characteristics increases. From our HERFD-XANES and XPS findings, the presence of Ce^{III} and hence of the oxygen vacancy, can be discarded, which means that the observed lattice expansion might then be only due to surface stress. To corroborate this assumption, the formation of species present at the NPs surface has been studied by FT-IR and TGA measurements. Indeed, our data confirms the presence of surface hydroxyl and carbonate groups that have a tensile effect on the crystalline lattice.

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Isotopic dual-radiolabeling of CeO₂ nanoparticles for uptake and dissolution studies

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Manufactured nanoparticles, such as CeO_2 , give rise to novel risks when released into the environment. To assess these risks it is important to quantify the nanoparticle mass flows, as well as their speciation and the mechanisms of their transformation. The radiolabeling of nanoparticles can be a valuable tool allowing detection at very low concentrations, regardless of any complex backgrounds. We have developed an innovative dualradiolabeling strategy for CeO₂ nanoparticles that will allow the tracking of these particles in uptake and dissolution studies with a sensitivity to dissolution effects.

The risks associated with manufactured nanoparticles (NPs) greatly depend on the transformation of said particles. Fast dissolving NPs can be treated as dissolved species while persistent NPs can be treated as particles. Slowly dissolving particles such as CeO₂ form a special case as transport and uptake can happen in particle form to deliver toxic ions *via* a Trojan horse effect. In particular for CeO₂, there are conflicting results on the role of dissolution effects in uptake by organisms.^[1,2] Based on our previous efforts we developed an isotopic dual-radiolabeling strategy to shed light on uptake pathways of anthropogenic cerium in organisms (Fig. 1).

EXPERIMENTAL. CeO₂ nanoparticles were procured from the JRC Nanomaterial Repository and activated *via* neutron irradiation at the TRIGA Mark II nuclear reactor of the Jozef Stefan Institute (JSI). 30 mg of CeO₂ NPs were sealed in a quartz ampule and irradiated at a neutron flux of 1×10^{13} for 63 hours.^[3]

Using the in-house IBA Cyclone 18/9 cyclotron at HZDR Research Site Leipzig Ce-139 radiotracer was produced by 10.9 MeV proton irradiation of a lanthanum metal foil at a current of 22 μ A for about 23 hours yielding a current integral of 500 μ Ah. According to the reaction cross-section, the optimum proton energy was set by the design of the target using an 850 μ m aluminum window.

The irradiated target was dissolved in conc. HNO_3 and the radiotracer was separated from the target material by extraction chromatography using UTEVA resin.



Fig. 1: Dual-radiolabelling strategy for manufactured CeO₂ NPs, illustrating the heterogeneous/isotopic distribution of the different radiotracers (particle volume vs. surface-near) in the nanoparticles and their release during dissolution. For dual-radiolabeling the Ce-139 solution was added dropwise to the neutron activated [Ce-141]CeO₂ NPs followed by a 5 hours tempering step at 300 °C prompting indiffusion of the Ce-139 into the NPs.

RESULTS. Irradiation of CeO₂ NPs with neutrons at the JSI nuclear reactor triggered the nuclear reaction Ce-140(n, γ)Ce-141, but also Ce-143 and Ce-139 were formed upon side reactions. Ce-143 was let decay, while Ce-139 was not expected to interfere with the further labeling strategy, since it was formed only two orders of magnitude less than Ce-141. Ce-139 was produced at the HZDR cyclotron with a yield of 0.1 MBq/µAh. The radiochemical yield of the tracer separation was 96 %.

The radiolabeling yield of the in-diffusion step was 99% thereby successfully generating dual-labeled NPs of [Ce-139/Ce-141]CeO₂ (Tab. 1). The different radiolabels are distributed differently in the particles. While the Ce-141 is equally distributed within the particles, the Ce-139 is located mainly at the particle surface. Therefore, the different properties of the radioisotopes enable us to track dissolution processes *via* the Ce-141/Ce-139 ratio (Tab. 2).

Tab. 1: Characterization data and specific activities reached for radiolabelled CeO₂ NPs.

Particles	Size [nm]*	Specific Activity [MBq/mg]
CeO2 NPs	221	-
[Ce-141]CeO ₂ NPs	223	0.255 Ce-141
[Ce-139/Ce-141]CeO2 NPs	219	0.272 Ce-139
		0.107 Ce-141

*: as measured by DLS.

Tab. 2: Cerium radiotracer properties.

Radioisotope	Half-life [d]	γ-peak [keV]
Ce-139	137.64	168.86
Ce-141	32.51	145.44

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Uncertainty explorations in the context of nuclear waste management

S. Pospiech, V. Brendler

The EURAD joint program deals with various aspects of nuclear waste management, covering the whole range from waste packages to geotectonics. Its work package 10 (UMAN) is dedicated to the management of uncertainties potentially relevant to the safety of different radioactive waste management stages and programs. The task of the Institute of Resource Ecology (IRE) is to contribute to the development of general guidelines to identify, categorize and process uncertainties.

The IRE contributes within UMAN concerning identification, classification and quantification of numerical uncertainties that might be of relevance in the various stages of radioactive waste management (RWM). Main results are summarized in a EURAD report.^[11] It provides definitions for frequently used terms, *e.g.* uncertainty in general, parameter uncertainty, uncertainty models, and aleatory *vs.* epistemic uncertainties. Sources used as input for this report were questionnaires, expert elicitation and literature survey. The following themes and stages of RWM were addressed:

- 1. National Program Management
- 2. Predisposal
- 3. Engineered Barrier System
- 4. Geoscience
- 5. Design and Optimization
- 6. Siting and Licensing
- 7. Safety Case

RESULTS. The consideration of urgency and potential benefit is the first step for a robust assessment of the priority for further investigation, established through expert judgement. It can be transferred into the following criteria:

- level of impact on safety,
- · level of impact on decision-making process,
- time span till utilization within the RWM stages,
- priority for further investigation.

A categorization of uncertainties can be based on very different features and is in principle depending on the purpose of application. Typical categories are:

- epistemic (more information will reduce it) *vs.* aleatoric (intrinsic properties, no reduction possible) uncertainties,
- time step of occurrence within RWM,
- RWM system component affected,
- management option (reduce/mitigate/ignore),



Fig. 1: Example for Fishbone (Ishikawa) diagram, highlighting the mutual relationship between uncertainties contributing to the overall uncertainty of distribution coefficients.

• scale of application (generic, national, site-specific),

• type of heterogeneity with regard to space or time. Two methodologies were applied for structuring uncertainties: fishbone diagrams and tables for uncertainty character-

istics. Examples are given in Fig. 1 and Tab. 1. Namely the fishbone diagrams are a flexible tool, regularly used to support the development of conceptual models required for uncertainty treatment. They can be used to combine and thus simplify uncertainty components, leading to model reduction. Also, they can carry information how the components are related to each other, e.g., if they are additive or multiplicative and/or repeatable or repeatedly arising. This is an important step for the uncertainty quantification / approximation by probability density functions. Such diagrams promote structuring individual sources of uncertainties in different systems. Numerical uncertainties must be evaluated and transformed in a way that a respective model can make directly use of them. The three most often used mathematical approaches are deterministic/probability statistics, worst case analysis and fuzzy set theory.^[2-4] Among the probability statistics, the two main branches are frequency statistics and Bayesian statistics.

Tab. 1: Disposal container uncertainty list applicable to high-level waste and spent nuclear fuel.

Component / process	Associated uncertainties	E/A*
Container corrosion (normal and specific aggressive conditions)	Interim storage times, container material composition & lifetime	E
Production of gases due to radiolysis and redox processes	Type and amount of gases produced, production rate	Ε
Leakages of gaseous compounds	Pressure build-up, release rates	Е
Overall size require- ment for a DGR**	Final volume of waste package per type of waste	Е
Heat generation with- in a DGR	Interim storage times, inventory	Е

*: (E)pistemic vs. (A)leatoric uncertainty type.

**: Deep Geological Repository.

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Implementation of the smart K_d-concept in the reactive transport code OpenGeoSys (OGS-6)

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Within the iCross project, the radionuclide transport through the far-field of a repository system is simulated with the transport code OpenGeoSys-6 (OGS).^[1,2] The big challenge in large-scale reactive transport modelling is to find an acceptable way to integrate geochemistry at affordable cost and to balance between geochemical realism and computational efforts. As an effective alternative to a direct coupling of OGS-6 with PHREEQC,^[3] the smart K_d-concept (www.smartkd-concept.de) has been implemented in OGS to consider variable geochemistry.^[4,5]

METHODOLOGY. The fundamental strategy of the smart K_d-concept, a mechanistic approach mainly based on surface complexation models, is to compute a priori multidimensional look-up tables with distribution coefficients (referred to as smart K_d-values) for a wide range of important environmental input parameters, e.g. pH, ionic strength, concentration of competing cations and complexing ligands. Thereby, the "Component Additivity" approach is used to describe a natural system close to reality. This bottom-up approach is based on the principle that the sorption of contaminants can be determined from the competitive mineralspecific sorption of each dissolved species on all surfaces. For each transport step, a smart K_d-value valid for the actual geochemical conditions is interpolated from the precalculated look-up table and used for the transport simulations. The schematic workflow is shown in Fig. 1.

t-dt t $t+dt$ timel	ine
Before onset of the simulation OpenGeoSys Op	

Fig. 1: Workflow of the look-up table approach in OGS-6.

MODEL SETUP. Based on a well-defined test case for a typical sedimentary rock system covering potential repository host rocks, namely salt and clay formations in Northern Germany,^[5] a 2D transport model was developed. Thereby, a reduced set of chemical components was considered to keep the model simple and to reduce the computational effort. The transport of H⁺, OH⁻, Na⁺, Cl⁻, Ca²⁺, dissolved inorganic carbon, and several radionuclides (RNs) was simulated



Fig. 2: Model domain with boundary conditions.

through an initially saturated column of 20 m in length and 3 m in width (Fig. 2).

RESULTS. To validate the concept, we compared the results of the look-up table approach with the results of the direct coupling of OGS-6#PhreeqC3.5.0. The results show a good agreement between both approaches for all considered repository-relevant RNs (Cs, Ra, Ni, Am, Th, Np, U). Figure 3 shows the spatial distribution exemplarily for neptunium(V) - Np(5) after 500,000 s simulation time. Np(5) has the weakest sorption of all considered RNs with the peak maximum at a distance of 11 m from the inflow compared to the other radionuclides. The curves of both simulations lie nearly perfectly one upon another. This implies that the look-up table approach works well. We are able to reflect the RN-specific retardation effect. This real application case serves as a comprehensive proof-of-concept and as a benchmark for field-scale transport simulations.



Fig. 3: Spatial distribution of neptunium after 500,000s simulation time.

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Thermodynamic reference database THEREDA: 10. The cumulative 2020 data release and major developments

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THEREDA^[1,2] is a joined database project to assess potential radionuclide migration scenarios accurately and to make well-founded predictions about the long-term safety up to one million years. The number of elements for which data is provided as well as the upgrades in supported geochemical codes necessitated a change in release strategy. Another innovation is the comparison of experimental data and THEREDA calculations, through which users can assure themselves of the quality of the data offered.

THEREDA provides data for actinides, fission products as well as matrix and building materials using the Pitzer ion interaction model^[3] for high saline solutions (ionic strength up to at least 7 molal) in the temperature range from 0 to 100 °C. In the fourth phase of the THEREDA project, the following relevant changes and new features have been introduced so far:

NEW RELEASE STRATEGY. In the earlier project phases, small, self-contained data packages on individual chemical systems were released. This should ensure data consistency and prevent cross-effects due to element combinations for which no validated data exists yet. This approach has become impracticable for updating purposes with increasing number of releases and supported geochemical codes. Therefore, all the released thermodynamic data and associated interaction parameters are now combined in one file per code. Furthermore, annual updates for the released data are intended. With the 2020 data release, THEREDA supports the codes ChemApp^[4], GEMS^[5], Geochemist's Workbench^[6], PHREEQC^[7] and ToughReact^[8].

DATA RELEASE 2020. In the present THEREDA data release, data for actinides and their chemical analogues (Th, U, Np, Pu, Am, Cm & Nd), fission products (Se, Sr, Tc & Cs) and matrix elements (Na, K, Mg, Ca, Al, Si | Cl, SO₄, CO₃) are combined. Many updates as well as a polythermal sele-



Fig. 1: Solubility of amorphous uranium(IV) hydroxide in 1 m NaCl solution at T = 25 °C. Symbols: experimental data,^[10–12] Lines: calculation using THEREDA (red), uncertainty range from solubility constants error^[13] (grey).



Fig. 2: Solubility of halite at T = 0–200 °C. Symbols: experimental data,^[14] Line: calculation using THEREDA.

nium(IV/VI) data set have been included. Also, the current version of CEMDATA (18.1) was integrated for the calculation of cementitious phases.^[9]

APPLICATION EXAMPLES. THEREDAs capabilities are now demonstrated using approx. 400 application case calculations, whose results were compared with experimental values published in literature. These demo calculations can be browsed on a new section of the website (<u>thereda.de</u>). Examples are shown in Figs.1 and 2.

ACKNOWLEDGEMENTS. Knowing that database work of this kind can only be accomplished within a great team, the authors thank all members of the THEREDA team at GRS, KIT-INE, TUBAF and PSI-LES. THEREDA is funded by the Federal Company for Radioactive Waste Disposal (BGE) with the contract number 45181017.

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SCIENTIFIC CONTRIBUTIONS (PART III)

Long-Lived Radionuclides in

BIOLOGICAL Systems

Proteomics confirms initiation of glutathione-mediated sequestration upon uranium exposure to tobacco (*Nicotiana tabacum*) BY-2 cells

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Increases in glutathione (GSH) and enzymes involved in the removal of reactive oxygen species (ROS) are part of the cellular response to heavy metal stress. Proteomics was used to confirm that uranium in culture causes an increase in GSH metabolism in tobacco BY-2 cells even under culture conditions of phosphate deficiency.

Plants have adapted various ways of dealing with heavy metal stress. This is of particular importance for hyperaccumulator plants, which are able to grow in soils with potentially toxic abundances of heavy metals. Uranium as a heavy metal has a high potential to lead to the formation of ROS through Fenton reactions, which could cause irreparable cell damage.[1] Plant cells typically respond by increasing GSH reserves and producing antioxidative enzymes to help in metal chelation and ROS removal. Glu-Cys ligase is one of the enzymes responsible for the production of GSH and the enzyme GSH S-transferase facilitates the conjugation of GSH to xenobiotic substances for detoxification.^[2] Superoxide dismutase, catalase and a variety of peroxidases are part of the enzyme repertoire involved in the inactivation of ROS. Previous research has exquisitely demonstrated the uranium-induced increase in GSH pools and ROS scavenging activity in Brassica napus using biochemical assays.[1] In this study, proteomics was used to contribute to a molecular process understanding of the interaction of uranium with plants by gauging the extent of the enzymatic response, using tobacco (Nicotiana tabacum) bright yellow 2 (BY-2) cells as a model for plant cells.

EXPERIMENTAL. BY-2 cells were grown in the dark in modified Murashige-Skoog (MS) medium (4.3 g/L basal salt, 200 mg/L KH₂PO₄, 30 g/L sucrose, 100 mg/L inositol, 1 mg/L thiamine and 200 µg/L 2,4D) for seven days.^[3] 1.5 mL of this culture were sub-cultivated into three flasks of 30 mL regular MS medium and into six flasks of 30 mL phosphate deficient MS medium (MS-red; containing only 1.7 mg/L phosphate) and grown for three days. Into three of the flasks of MS-red medium, uranyl nitrate (20 µM) was added. All flasks were incubated for 24 h. 15 mL of each culture were gently centrifuged to remove the culture medium and the remaining pellet was washed with 15 mL of 50 mM MOPS buffer (pH7.5). After another round of centrifugation, the cell pellet was resuspended in extraction buffer (50 mM MOPS, 2 mM dithiothreitol (DTT), 5% polyvinyl polypyrrolindone and protease inhibitors) and lysed by grinding with a mortar and pestle in liquid nitrogen. The crude extract was spun down at $16,000 \times g$ for 20 min at 4 °C and the sediment was discarded. Protein in the supernatant was subsequently quantified using a Pierce™ Bradford assay kit. 100 µg of the protein extract were centrifuged through cleaned 30 kDa Amicon® Ultra $0.5\,\text{mL}$ filters and buffer exchange with $8\,\text{M}$ Urea/50 mM Tris, reduction with 8 M Urea/50 mM Tris/100 mM DTT and alkylation with 8 M Urea/50 mM Tris/50 mM Iodoacetamide were all carried out inside the filters. A final buffer exchange step with 50 mM NH₄HCO₃ was performed prior to digestion with 0.5 µg of LysC enzyme and 2 µg of trypsin at 37 °C overnight. The resulting peptides were then centrifuged through the filters using extra NH4HCO3. The peptides were desalted using C18 STAGE-tips and subsequently resolved Tab. 1: Logarithmic (base 2) fold changes of differential protein expression showing effects of uranium treatment and phosphate deficiency on BY-2 cells. Fold change of 1 indicates the treatment caused a doubling of protein compared to the control; \neg 1 indicates that the protein quantity was half of the control; 0 means there was no difference in protein expression. Significance is denoted by **p* < 0.05, ***p* < 0.01, ****p* < 0.001.

	Log ₂ Fold Change		
Protein	Uranium treatment	Phosphate deficiency	
Glutathione S-transferase	0.82*	2.95***	
Glutathione peroxidase	0.80*	1.17***	
Glutamate-Cysteine ligase	-0.03	0.82*	
Superoxide dismutase	0.08	1.29**	
Catalase	0.07	0.67	

and analyzed by an EASY-nLC 1000 chromatography system coupled to an LTQ-Orbitrap Velos Pro mass spectrometer (Thermo Scientific). Protein and peptide identification were carried out as described elsewhere.^[4]

RESULTS. The experimental conditions not only exposed the plant cells to uranium but also introduced conditions of phosphate (Pi) deficiency in order to avoid the formation of insoluble complexes of uranyl(VI) phosphate. It was apparent that Pi deficiency was already inducing stress, based on protein expression patterns (Tab.1) and cell growth (not shown). All five proteins of interest were overexpressed with Pi deficiency, and all enzymes, except catalase, showed significant overexpression. This highlighted the importance of Pi to the culture and that three days of Pi deficient cultivation was sufficient to cause stress. Nevertheless, the presence of uranium elicited a specific response, whereby GSH S-transferase and GSH peroxidase were almost twofold overexpressed, suggesting the cells were sensitized to the presence of a foreign substance. GSH has been seen to form strong complexes with uranyl(VI) ions and other heavy metals.^[5] Since Pi deficiency had already led to elevated GSH levels, it can be deduced that in this instance, energy is being expended to sequester cytosolic (and possibly organellebound) uranium with the help of GSH S-transferase, thereby reducing the propensity toward cellular damage via ROS. At present, further investigation is underway to determine how similar the cellular response would be with the absence of stress induced by Pi deficiency.

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Do carotenoids quench uranium(VI) luminescence in TRLFS spectra of plant cells?

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Previous Time-Resolved Laser-Induced Fluorescence Spectroscopic (TRLFS) studies on the interaction of U(VI) with the two plant cell systems *Brassica napus* (rapeseed/canola) and *Daucus carota* (carrot) showed that spectra of significantly lower quality had been obtained for the latter. Therefore, experiments on carotenoid extraction from the cells were carried out. It was shown that both organisms contain considerable amounts of carotenoids, however, these substances are not responsible for the observed luminescence quenching effects of U(VI).

Time-Resolved Laser-Induced Fluorescence Spectroscopy (TRFLS) is applied to study the interaction of U(VI) with plant cells and to elucidate the U(VI) speciation on/in the cells and in solution on a cellular and molecular level. TRLFS measurements with the two plant cell species Brassica napus (B. napus) and Daucus carota (D. carota) revealed that a significantly lower luminescence intensity could be observed in the D. carota system for U(VI) in the initial nutrient media, in supernatants after cell contact as well as in the biomass. Since the U(VI) concentrations in both systems were similar, these observations were attributed to the quenching of the U(VI) luminescence in the D. carota system. A possible explanation for these effects could be the occurrence of carotenoids in the plant cells. For uranyl(VI) compounds and all trans- β carotenoids, the absorption and emission spectra lie in a comparable wavelength range.^[1-3] Thus, a significant reduction of the U(VI) luminescence can be expected when the absorption of excitation light by carotenoids present in the D. carota system occurs. The aim of this study was to extract carotenoids from *B. napus* and *D. carota* suspension cells and to examine the obtained carotenoid solution by UV-vis spectroscopy. Based on the findings, it was concluded whether the carotenoids contained in the cells could be responsible for the quenching effects in the D. carota system or not.

EXPERIMENTAL. B. napus cells were obtained as callus cell culture from DSMZ Braunschweig. D. carota callus cells were established as described by Sachs et al.[4] The cultivation of B. napus suspension cells was performed as described in the literature.^[5] D. carota cells were cultivated in the same way, however, using medium B5C2^[6] instead of medium R. Seven days after subcultivation, the suspension cells were separated from the culture media and frozen at -20 °C. The frozen cells were freeze-dried for three days (Alpha 1-4 LSC, Christ). Subsequently, the dried cells were crushed with a mortar and pestle while adding liquid nitrogen. The mortared cells were mixed in round bottom flasks with an acetone/petrol ether mixture (v/v = 1:1) with a solid/liquid ratio of 1g/30 mL according to Yan et al.^[7] The mixtures were heated to 42 °C in a water bath under reflux and stirring for 1 h (Fig. 1). After cooling the mixtures, the cell residues were separated by filtration and clear, bright yellow carotenoid solutions were obtained. The UV-vis measurement of the carotenoid solution was performed in quartz glass cuvettes (HELL117100F-10-40, Hellma Analytics) with a Silver-Nova spectrometer (Silver-Nova, StellarNet Inc.).

RESULTS. It was found that both systems, *B. napus* and *D. carota* cells, yield extracts that had a strong yellow coloration indicating the presence of carotenoids. The UV-vis



Fig. 1: Carotenoid extraction with *B. napus* suspension cells (left) and normalized absorption spectra of carotenoid extracts of *B. napus* and *D. carota* cells.

measurement resulted in characteristic carotenoid spectra (Fig. 1) for the extracts of both organisms. According to this, it was concluded that both cell cultures contain similar carotenoids. An estimation of the carotenoid content based on the UV-vis measurement and the applied cell mass was possible and allowed the conclusion that both organisms also contain comparable amounts of carotenoids. Accordingly, it could be excluded that carotenoids are responsible for the observed quenching of the U(VI) luminescence in the *D. carota* system, since otherwise the quenching would also have to occur for *B. napus*, which was not the case. Therefore, the quenching effects must have a different reason. Most probably, small differences in the composition of the used cell culture media for both organisms, which still have to be elucidated, are responsible for the quenching effects.

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Interactions of Desulfosporosinus hippei DSM8344 with Eu(III)

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Different *Desulfosporosinus* species were detected in bentonite and clay formations, which are suitable as barriers for nuclear waste disposal. Therefore, *Desulforsporosinus hippei* DSM8344^T was selected as a good representative to investigate its interactions with europium as a function of time. Time-resolved laser-induced fluorescence spectroscopy (TRLFS) revealed the presence of two Eu(III) species associated with the cells.^[1]

Microorganisms, which are common in bentonite or clay formations, can potentially impact the retardation or mobility of radionuclides.^[2] Diverse *Desulfosporosinus* species were verified to be present in bentonite,^[3] a potential geotechnical barrier material, and clay formations, a potential host rock for nuclear waste repository. *Desulfosporosinus hippei* DSM8344^T a mesophilic sulfate-reducing bacterium was selected to study its interactions with europium which is a non-radioactive analog for trivalent actinides, *e.g.* curium and americium.

EXPERIMENTAL. *D. hippei* DSM8344^T was purchased from the Leibniz Institute DSMZ-German Collection of Microorganisms and Cell Cultures (DSMZ, Braunschweig, Germany) and cultivated in DSM641 medium as already described.^[11] Europium bioassociation in batch systems was performed in a synthetic Opalinus Clay porewater solution (pH~5) with 10 mM sodium lactate and 50 or 100 µM Eu(III) as europium chloride.^[4] The Eu(III) concentration in the supernatant after incubation for different times was measured with inductively coupled plasma mass spectrometry (ICP-MS). Eu(III) was used as molecular probe to investigate its speciation on cells and in supernatants using luminescence spectroscopy (TRLFS). Therefore, supernatants, as well as washed and resuspended cells, were investigated separately in a quartz glass cuvette at room temperature as described earlier.^[5]

RESULTS. The association of Eu(III) on *D. hippei* DSM8344^T cells was investigated at 50 μ M and 100 μ M Eu(III) as a function of incubation time. The initial concentration of Eu(III) decreased only slightly under these conditions (data not shown). Approximately 17% and 13% of the Eu(III) was associated with the biomass at 50 μ M and 100 μ M Eu(III), respectively. During the incubation of up to 48 h an agglomeration of the cells was observed. The batch experiments could be done only with a relatively low amount of biomass. This might be one reason for the small effect observed in relation to the amount of bioassociated Eu(III).

Figure 1 summarizes the luminescence data measured for the Eu(III)-*D. hippei* system in the presence of 100 μ M Eu(III). The spectrum of the blank (red graph in Fig. 1) is characterized by a strong increase of the ⁷F₂-transition in comparison with the Eu³⁺-ion (black graph in Fig. 1). The ⁷F₂-transition is sensitive to changes of the chemical environment of the luminescent Eu(III).^[6] The luminescence decay follows a mono-exponential law with an average lifetime of 132 ± 5 μ s (N_{H2O} = 7.5). This value is longer compared with the Eu³⁺-ion (110 μ s, N_{H2O} = 9.0). Both information confirm a strong complexation of Eu(III) by inorganic components of the synthetic Opalinus Clay porewater at pH ~5.

The spectra showed no changes for Eu(III) present in the supernatants as a function of the incubation time. Hence, Fig. 1 presents the averaged data (olive graph in Fig. 1). The spectroscopic findings of Eu(III) in the supernatants in terms of



Fig. 1: Eu(III)-luminescence spectra measured after incubation of the *Desul*fosporosinus hippei DSM8344^T cells in synthetic Opalinus Clay porewater (pH ~ 5) and 100 μ M Eu(III).

intensity of the ⁷F₂-transition and the luminescence decay behavior were similar to the values of the blanks. A change in the luminescence behavior in the supernatants could be attributed to a release of metabolites or lyzed cells. In the present case the amount of cells was probably too small for the observation of measurable effects.

The shape of the spectra of the cell-bound Eu(III) (blue and dark cyan graphs in Fig. 1) differs from the blank and supernatant. The appearance of the symmetry-forbidden ⁷F₀-transition in combination with the intensity increase of the 7F2-transition points to a different Eu(III) coordination environment as found in Eu3+, supernatant, and blank. In contrast, a bi-exponential luminescence decay was measured: $61 \pm 27 \,\mu s$ and $367 \pm 191 \,\mu s$ (N_{H2O} = 2.3). The short lifetime is even shorter than the value of the Eu³⁺-ion. Additional ligand-specific quench processes other than from OH⁻-ligands could explain this finding.^[7] The long lifetime is a measure for significant changes in the first Eu(III) coordination sphere due to interactions with bacterial surface functionalities. In conclusion, the findings indicate a biosorption or bioaccumulation of Eu(III) by the cells in combination with a Eu(III) speciation change.

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Influence of fetal bovine serum and citrate on Eu speciation in cell culture medium

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The solubility and speciation of Eu in cell culture media with different contents of fetal bovine serum (FBS) and citrate were investigated by mass spectrometry (ICP-MS) and luminescence spectroscopy (TRLFS).^[1,2] Results demonstrate that FBS is the key factor affecting Eu solubility and speciation, while citrate has no significant effect in the presence of FBS. Eu speciation is dominated by (a) soluble protein complexe(s) in FBS containing medium. In FBS free medium, insoluble Eu complexes with inorganic anions, mostly carbonate and phosphate, dominate. Up to now, thermodynamic modeling of the Eu speciation largely diverges from experimental results.

The knowledge of the lanthanides' speciation in cell culture media is crucial for the elucidation of the chemical affinities of these metals to the prevailing biomolecules and the resulting biological effects.

EXPERIMENTAL. Four different cell culture media, *i.e.* standard medium with 10% FBS (st), serum-reduced medium with 1 % FBS (sr), serum-free medium (sf), and st-medium with 0.1 mM citrate (st-cit), were incubated with 10⁻⁵ or 10⁻³M Eu. Eu solubility was derived from ICP-MS measurements of the medium supernatants after sterile filtration. TRLFS spectra in cell free culture media were recorded at room temperature after 24 and 48 h incubation time. The theoretical Eu speciation in each medium was calculated with thermodynamic modeling using current data.

RESULTS. After 24 h, Eu is completely soluble up to 1 mM in medium containing 10% FBS (Fig. 1). Lowering the FBS content leads to significant decrease in Eu solubility indicating that proteins of FBS stabilize Eu in solution, thus making it bioavailable for cells. Identical results were obtained after incubation for 48 h. Citrate seems to have no additional effect on Eu solubility.



Fig. 1: Solubility of Eu in different cell culture media after incubation for 24 h. The dotted line equals 100% solubility.

At 10^{-3} M Eu, visible precipitation occurs in sr- and sf-medium. Analysis of the settled deposits revealed elemental ratios of Eu:P ~ 1:1, indicating the formation of insoluble Eu-phosphate, and Eu:Ca ~ 1:0.2. TOC measurements indicate that also inorganic carbonate is contained in the deposits from both cell culture media, whereas the content of organic carbon is significantly higher in deposits from sr-medium. The latter points to organic ligands like proteins to be present in deposits from sr-medium, too. In general, TRLFS spectra recorded after 24 and 48 h incubation time are very similar and indicate the coexistence of two Eu species in all four cell culture media (Fig. 2). In medium with 10 % FBS, the long lifetimes suggest the presence of only one or none water molecule left in the Eu³⁺ complexation shell. Hence, the complexation of Eu by (a) protein(s) from FBS can be assumed which concurs with solubility results. Citrate addition has no significant effect on TRLFS spectra indicating the same speciation in st- and st-cit-media. In FBS deficient media, especially at 10-3 M Eu, TRLFS spectra demonstrate the predominant complexation of Eu by (a) protein(s) as long as some FBS is present. However, also Eu complexes with phosphate, carbonate, and organic ligand(s) occur. This becomes obvious in the supernatant of sf-medium by completely different steady-state spectra and a significantly shorter luminescence lifetime (see Fig. 2C and D). In summary, we propose four different Eu species in dependence on the composition of the cell culture medium: Eu bound to (i) protein(s), (ii) phosphate, (iii) carbonate, and (iv) organic low molecular weight ligand(s).



Fig. 2: Steady-state spectra (A, C) and luminescence decay curves (B, D) of 10^{-5} (A, B) and 10^{-3} M Eu (C, D) in different cell culture media after incubation for 24 h. Data for the Eu³⁺ aquo ion is given in black.

Independent from the medium composition, calculated Eu speciation diagrams are very similar and predict a predominant carbonate and, to some extent, phosphate complexation, whereas organic constituents are predicted to be negligible. The main reason for this discrepancy with our experimental results is most probably the lack of thermodynamic data for Eu complexes with proteins and other organic ligands as well as the variable and largely unknown composition of FBS. Hence, there is an urgent need to extend the current thermodynamic data sets of Eu complexes with organic ligands and proteins, to improve thermodynamic modeling for speciation predictions in complex biological fluids.

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Lanthanide-dependent methanol dehydrogenase studied by computational chemistry

S. Tsushima

Molecular dynamics (MD) simulations were employed to rationalize enzymatic activities of methanol dehydrogenase (MDH) carrying different lanthanides.^[1]

Lanthanides were long believed to have no biological function until a catalytic role of La³⁺ ion for MDH was discovered.^[2] This fact expanded the importance of Ln³⁺ ions in biochemistry and in bacterial metabolism. Moreover, it was revealed that bacteria growth depends significantly on the size of lanthanide and that the activity of MDH decreases as the size of Ln³⁺ ion decreases. This fact has significant impact to chemical engineering and could possibly provide a breakthrough in Ln³⁺ separation.^[3–5] Here, MD simulations were employed to study the active sites of Ln³⁺–MDH.

CALCULATION. MD simulations were performed using AMBER 15. Protonation state of the protein was adjusted to model pH7.4. Na⁺ ions were added to make the system electrostatically neutral. TIP3P waters were added with minimum water thickness of 6 Å. 1 ns preconditioning run was performed after which 50 ns MD run was performed in a periodic boundary condition in NPT ensemble at 320 K. For each Mⁿ⁺–MDH, 50 ns MD simulation was performed, from which only the last 30 ns MD trajectory was used for statistical sampling.

RESULTS. MD simulations were performed on Mⁿ⁺–MDH from strain SoIV (XoxF-type) for Mⁿ⁺ = La³⁺, Eu³⁺, Yb³⁺ and also that isolated from Methylorubrum extorquens AM1 (MxaF-type) for $M^{n+} = Ca^{2+}$, Eu^{3+} . Initial structure of the protein including cation and cofactor was taken from the crystal structure. Three carboxylic groups of PQQ were deprotonated. In all Mn+-MDH carrying different metals, the residues associated with the metal centers are identical, namely Glu¹⁷², Asn²⁵⁶, Asp²⁹⁹, and Asp³⁰¹. These metal-binding residues are also identical to those in the crystal structure of Eu³⁺-MDH. Among different Ln³⁺-MDH (La³⁺, Eu³⁺, Yb³⁺), all metal-binding residues are in the same binding modes except for Glu¹⁷² which switches from bidentate to unidentate in going from lighter (La, Eu) to heavier (Yb) lanthanide (Fig. 1). In addition, coordination of PQQ to Ln3+ switches from chelate binding mode involving pyridine nitrogen to unidentate fashion with single dangling carboxylic oxygen in going from lighter (La, Eu) to heavier (Yb) lanthanide. To fill the remaining coordination site around Ln center, maximum of one water is attached to the metal. There are discreet differences in the residues associated to PQQ amongst La³⁺-, Eu³⁺-, and Yb³⁺-MHD. In Eu³⁺-MHD, PQQ interacts with (besides Eu³⁺) Trp²³⁸ (through π - π stacking), Cys¹⁰⁴-Cys¹⁰⁵ (to disulfide bond), and Arg³²⁶ (through H–bond), whereas no interaction with disulfide bond is found in La3+-MDH and no interaction with $\mathrm{Arg^{326}}/\mathrm{Trp^{238}}$ in Yb³⁺–MDH.

Basically, as it shifts from lighter to heavier lanthanides, lanthanide contraction triggers closer contact between Ln with protein or with PQQ, favoring unidentate coordination. Coordination mode of PQQ to Ln³⁺ has critical importance to the activity of Ln³⁺–MDH because nucleophilic attack of methanol oxygen on PQQ occurs on its carbonyl oxygens. Upon switching of the Ln–to–PQQ binding from chelate binding to unidentate, carbonyl oxygens become loosely associated to the metal. Thereby PQQ exhibits less tendency towards reduction when associated with heavier lanthanides. The difference between the active sites of La³⁺– and



Fig. 1: Representative structures in the vicinities of the active sites of Eu³⁺– (upper) and Yb³⁺–bound MDH. Ln³⁺ ions and PQQ molecules are highlighted with red and blue.

Eu³⁺–MDH is overall minor and the largest difference is that there is additional water coordinated to the metal in the case of the La³⁺–MDH. It is unclear at this moment if and how additional water coordination correlates to the observed higher activity of La³⁺–MDH. It is possible that solvent– accessibility of the metal in La³⁺–MDH stimulates facile contact between PQQ and MeOH and eventually leads to the higher activity of La³⁺–MDH.

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Interactions of Brassica napus callus cells with Eu(III) and U(VI)

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Experiments using *Brassica napus* callus cells grown *in vitro* in the presence of Eu(III) or U(VI) showed that *B. napus* cells were able to bioassociate both potentially toxic metals (PTM). Most of the Eu(III) and U(VI) was found to be enriched in the cell wall fraction. Under high metal stress (200μ M), cells responded with reduced cell vitality and growth. Speciation analyses using both metals as luminescence probes confirmed that *B. napus* callus cells provided multiple-binding environments for Eu(III) and U(VI).^[1]

The accumulation of radionuclides, such as U(VI), and other potentially toxic metals (PTM), such as Eu(III), into plants, and, thus, into the food chain, represents a potential pathway for human exposure. Hence, studies exploring the interaction of PTMs with plants at the cellular and molecular level are necessary. Accordingly, our study was designed to determine the tolerance of canola (*B. napus*) callus cells to U(VI) and Eu(III) at two different metal concentrations. The effects of both PTMs on cell growth and vitality, as well as on the total phenolic content of the cells, were studied. Furthermore, this investigation also focused on the speciation of bioassociated U(VI) and Eu(III) and their distribution in various fractions of *B. napus* cells, since *B. napus* is known to be able to accumulate PTMs in higher quantities than many other species.^[2]

EXPERIMENTAL. *Brassica napus* callus cells (DSMZ PC-1113) were grown on a solid medium R with a reduced phosphate concentration spiked with 20 or 200 μ M UO₂(NO₃)₂ or 30 or 200 μ M EuCl₃ (99.999%, Aldrich, Taufkirchen, Germany) for about 6 weeks. All details concerning cell cultivation, vitality measurements, estimation of phenolic compounds, statistical analyses, cell fractionation, and laser-induced luminescence spectroscopy experiments can be found in literature.^[1]

RESULTS. Brassica napus cells had a strong capacity to bioassociate both PTMs, $628 \text{ nmol Eu}/g_{\text{fresh cells}}$ and 995 nmolU/g_{fresh cells} in the presence of 200 µM Eu(III) and U(VI), respectively. Most of the Eu(III) and U(VI) ions were bound on the cell wall fraction. This finding indicates that the cell wall might serve as an effective protective barrier of metabolically active cells against the threat of PTMs. In particular, under high Eu(III)/U(VI) stress, both metals were found in the cytosol fraction indicating an uptake of Eu(III)/U(VI) into the cells. Moreover, this study confirmed that cell growth was reduced in combination with a decrease in cell vitality (cf. Fig. 1). The total cellular phenolic content, which was expected to increase due to PTM stress, was similar for cells that were exposed to 30 or 200 µM Eu(III), 200 µM U(VI), as well as for our controls. Only a slight tendency for a slightly higher phenolic content was found for cells grown in the presence of 20 µM U(VI).

The appearance of the symmetry-forbidden $^{7}F_{0}$ peak in the luminescence emission spectra of Eu(III) bound to *B. napus* cells points to the formation of Eu(III) complexes. The occurrence of a bi-exponential luminescence decay confirmed the existence of two Eu(III) coordination environments. The strong intensity of the $^{7}F_{2}$ peak, as a measure of changes in the Eu(III) speciation, indicates the formation of strong Eu(III) complexes. Further analysis of the Eu(III) coordination environment revealed strong inner-sphere Eu(III) spe-



Fig. 1: Schematic overview of our study. TRLFS: Time-Resolved Laser-Induced Fluorescence Spectroscopy.

cies, possibly with organic phosphate and carboxyl groups provided by the B. napus cells. From this we conclude, that B. napus cells show multiple-binding environments for Eu³⁺. Due to the bioassociation of U(VI) onto B. napus cells, we observed a significant bathochromic shift of the U(VI) emission bands of the plant cell species compared to the spectra of the reference solution (U(VI) in 0.154 M NaNO₃) (cf. Fig. 1). Further, by comparing our findings with model compounds and other biological systems, there is a clear argument to be made for the predominant binding of U(VI) by organic and/or inorganic phosphate groups of the plant biomass. In conclusion, B. napus cells are able to accumulate and tolerate PTMs like Eu(III) and U(VI). As a result, these metals have the potential to enter the food chain and may become a severe health risk for humans. The results of this study were obtained through a combination of biological, biochemical, and spectroscopic methods. Hence, this integrative approach will contribute to an enhanced understanding of the interaction processes between actinides/lanthanides and plants at the molecular level, which is important for modeling the transfer of these elements in the environment.

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Evidence of uranium immobilization by *Magnetospirillum magneticum* AMB-1 under ambient conditions

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The α -proteobacterium *Magnetospirillum magneticum* AMB-1 is known as a motile, magnetotactic, gram-negative bacterium. It is facultative anaerobe and usually found in oxic-anoxic transition zones. In kinetic batch experiments at pH 6.5, cells of *Magnetospirillum magneticum* AMB-1 were contacted with 0.01, 0.05, and 0.1 mM U(VI) for 24 h. The results show that the cells are able to remove most of the U from the solution during the first hours of incubation, independently on the initial U concentration. For TEM-based studies, the cells were contacted with 0.1 mM U(VI) for 24 h. Element distribution analysis reveal that U is bound predominantly in the outer membrane of the cells.

EXPERIMENTAL. 100 mL of a cell suspension of Magnetospirillum magneticum AMB-1 were grown in 200 mL flasks in 1.5 mM MagMin medium, modified by CEA Cardarache with the addition of iron malate (1:2,000) and vitamin elixir (1:2,000).^[1] The flasks were kept at 30 °C on a rotary shaker at 80 rpm. An inoculation of the bacteria was carried out every 2-3 days. For each U kinetic batch experiment, 8 flasks were taken, each with an OD_{600nm} of 0.3–0.4, and washed twice with sterilized tap water with a pH adjusted to 6.5. The final OD_{600nm} in each experiment was 1.66 ± 0.04 which corresponds to 76.28 ± 0.66 mg dry biomass. The washed cells were suspended in a 200 mL flask with 100 mL sterilized tap water (pH6.5), and a 0.1 M stock solution of UO₂(NO₃)₂ was added to a final concentration of 0.01, 0.05, and 0.1 mM, respectively. The cell suspensions were incubated at 30 °C and 80 rpm on a rotary shaker. After distinct times, 1 mL samples were taken to determine the residual U concentration in the supernatant by inductively coupled plasma mass spectrometry (ICP-MS) measurements. For transmission electron microscopy (TEM) studies, cells of Magnetospirillum magneticum AMB-1 previously contacted for 24 h with UO₂(NO)₃ solution (0.1 mM, pH6.5) under sterile conditions, were recovered by centri-fugation (5500 rpm, 10 min). The supernatant was removed, and the cell pellet was washed twice with sterile tap water (pH6.5). Subsequently, the pellet was fixed at 1% (vol/vol) glutardialdehyde from 25% (vol/vol) stock solution. Further preparation was performed at the Center of Regenerative Therapies (TU Dresden). About 100-nm-thick sections of the samples were analyzed by TEM and spectrum imaging analysis based on energy-dispersive X-ray spectroscopy (EDXS) for element distribution determination.

RESULTS. The results show a substantial removal of U by *Magnetospirillum magneticum* AMB-1 cells during the first hours of incubation with emphasis on the first hour, where even 39.55%, 63.71%, and 77.51% of U was removed from the suspensions with an initial U concentration of 0.01, 0.05, and 0.1 mM U, respectively (Fig. 1). Furthermore, the U removal capacity of *Magnetospirillum magneticum* AMB-1 cells at 0.01, 0.05, and 0.1 mM U after 24 hours was determined to be 0.87, 12.54 and 30.17 mg U/g DBM, which correspond to 29.94%, 85.99% and 95.51%, respectively. TEM imaging of *Magnetospirillum magneticum* AMB-1 cells contacted with 0.1 mM U for 24h (Fig. 2, left) shows an area of high massabsorption contrast that can be interpreted as a detaching



Fig. 1: Comparison of U sorption by *Magnetospirillum magneticum* AMB-1 during the first 6 h of incubation using initial uranium concentrations of 0.01, 0.05, and 0.1 mM under aerobic conditions at pH 6.5.



Fig. 2: TEM image (left) and EDXS-based U element distribution analysis (right) of *Magnetospirillum magneticum* AMB-1 cells. Initial U concentration: 0.1 mM, 24 h incubation.

cell membrane. Element distribution analysis (Fig. 2, right) performed for this bacterium clearly indicates that U is entirely present in this cell membrane. The fast immobilization of U by *Magnetospirillum magneticum* AMB-1 in the membrane within the first hours can be interpreted as a passive process. Live-dead staining experiments performed with Syto®9 and Propidiumiodide show that the cells are mostly alive during the batch sorption experiments. Consequently, we conclude that living cells of *Magnetospirillum magneticum* AMB-1 are responsible for the immobilization of U by passive biosorption during the first hours of incubation.

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Microbially influenced corrosion of cast iron in different types of bentonite: microscopic investigation

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Cast iron is used as a potential container material for radioactive waste in a deep geologic repository (DGR) concept. Therefore, mesocosms containing different types of bentonite, cast iron coupons and a synthetic Opalinus Clay porewater solution were setup. After an incubation of 271 days, the surface damage associated with the corrosion was characterized by Scanning Electron Microscopy with Energy Dispersive X-ray (SEM-EDX) analysis to verify the impact of indigenous microorganisms on the corrosion process in the various bentonites. The results showed that selecting the appropriate bentonite for a DGR with the goal to control the microbial activity and thus corrosion of cast iron is a crucial step.

EXPERIMENTAL. The experimental mesocosm-setup was described elsewhere.^[1] For the current setup three different bentonites were used: B25, Calcigel[®] and MX-80. Aliquots of bentonites (15g) were dispersed into glass bottles together with 30 mL of synthetic Opalinus Clay porewater solution and GGG40 cast iron plates. Also, some of the mesocosms were supplemented with 5 mM sodium lactate and hydrogen (to 0.5 bar of total pressure) to stimulate microbial activity. The mesocosms were incubated for 271 days at 30 °C.

The mesocosm material was split after the incubation time and subjected to various analyses. First, metal coupons were removed and washed with deionized water, a portion of the slurry was used for DNA extraction, and the remainder was centrifuged. The filtrated supernatant was subjected to geochemical analysis to determine the concentration of various metals, sulfate, lactate, acetate and pyruvate (ICP-MS, IC, HPLC). To evaluate the changes in surface morphology and identify the corrosion components, metal coupon surfaces were examined by scanning electron microscopy coupled with energy-dispersive X-ray spectroscopy (SEM-EDX).

RESULTS. The presence of black precipitates was visible for the mesocosms containing Calcigel with lactate as substrate and for all the substrate-containing samples with MX-80. The geochemical analysis of these two bentonites showed a decrease in the sulfate concentration, which could be the electron acceptor for the respiration of sulfate-reducing bacteria (SRB). The strongest effect appeared in MX-80 bentonites.

With the exception of the B25 sample without a substrate, SEM data showed significant surface damage for all samples. Furthermore, two types of crystalline structures were found, covering to some degree the surface area of the coupons. One of these structures was iron and/or calcium carbonate, which was present in all the samples based on EDX (Fig. 1) and Raman spectroscopy (data not shown).

Iron carbonate in this case is most likely the product of the CO_2 -corrosion. It forms a protective layer which reduces the rate of corrosion. It has been argued that the formation of this kind of protective iron carbonate films decreases the corrosion rate by blocking the surface of the metal rather than reducing the diffusion of the corrosive material.^[2]

Iron sulfide was found on the surface of cast iron coupons incubated in B25 mesocosms supplemented with hydrogen, in MX-80 mesocosms with no substrate and with hydrogen (Fig. 2), and in Calcigel mesocosms without any substrate.



Fig. 1: SEM-micrograph and selected EDX-analysis of the surface of GGG40 coupon incubated with Calcigel and 5 mM lactate in synthetic Opalinus Clay porewater solution. The square inset shows the EDX result of the marked position.



Fig. 2: SEM-micrograph and selected EDX-analysis of the surface of GGG40 coupon incubated with MX-80 in synthetic Opalinus Clay porewater solution and hydrogen as a substrate. The square inset shows the EDX result of the marked position.

Iron sulfides can occur as part of a mixture of anaerobically corroded iron found in mineral and organic deposits. But iron sulfides are also characteristic products of SRB-induced corrosion.^[3] So, the presence of such products could be an indication for SRB activity. The analysis of the microbial communities is still ongoing. But a different microbial community structure was observed by intergenic spacer analysis (RISA) depending on the conditions applied and used bentonite (data not shown). So, different bentonites show different reactivity and a distinct microbial community. Therefore, the selection of the bentonite seems to have a big influence on the integrity of the container material and, thus, on the safe long-term storage of nuclear waste.

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Corrosion at the bentonite-cast iron interface under high-salt conditions with H₂

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The metal canister that contains the high-level radioactive waste (HLW) in a deep geological repository (DGR) should act as an effective shielding against external disturbances like the influx of pore water through the bentonite barrier. In order to study the resistance of ductile iron against corrosive conditions in a DGR located in saline clay rock formations, microcosms were set up. First results indicate the formation of deep pits that expand into the metal, leading to a high porosity of the protective metal barrier.

Corrosion is the destructive attack of a metal by its reaction with the environment.^[11] The safe retention of HLW in a DGR is realized by applying a multi barrier system where the metal canister is in direct contact to the bentonite barrier, which acts as buffering and sealing material when clay or granite serves as host rock. Cast iron offers excellent machinability combined with good impact resistance which makes it a suitable canister material.^[2] When selecting clay rock in the northern parts of Germany as final disposal site, saline pore water can enter the DGR. Anaerobic corrosion can take place at the interface between bentonite and cast iron, subsequently leading to the generation of H₂ that can promote microbial activity.^[3]

In order to study the influence of saline pore water and H_2 on the corrosion process of cast iron as well as on the microbial community structure, microcosms were set up.

EXPERIMENTAL. Microcosms were set up as described previously^[2] with the additional application of three ductile iron plates (cast iron type GGG40; L×W×H:10×10×5mm) in each microcosm. For simulating saline pore water, diluted cap-rock solution (CAP) was used.^[4] The microcosms were incubated for 30 days at 37 °C. Sampling was performed as described.^[3] After opening the flasks, cast iron plates were taken out with sterile plastic forcipes and stored under anaerobic conditions until further analytics. Surface analyses were performed by low-voltage Scanning Electron Microscopy (SEM) and local cross-sections were prepared using a 30 keV Ga+ Focused Ion Beam (FIB) and subsequently analyzed with SEM in a FEI Helios Nanolab 660. Extraction of DNA and sequencing was performed as described elsewhere.^[3]

RESULTS AND DISCUSSION. The influx of naturally formed saline pore water into the DGR is a crucial issue for the long-term stability of the cast iron canister harboring the HLW. Usually, corrosion of ductile iron in soils is decelerated, but it has been shown that chloride ions participate in corrosion reactions.^[5] Higher chloride ion concentrations lead to higher corrosion rates, larger pit depths and suppress the initial decreasing of corrosion rates.^[5]

The CAP solution applied contains 14 wt.-% of NaCl and is therefore much higher concentrated compared to previous experiments that dealt with cast iron corrosion in solutions ranging between 1 wt.-% and 3.5 wt.-% NaCl.^[5,6] Indeed, the analysis of the cast iron surfaces reveals a more perforated surface compared to cast iron plates that were incubated in bentonite microcosms containing Opalinus Clay pore water (OPA) and H₂ (data not shown). The observed pits are very



Fig. 1: FIB-SEM cross-section through an incubated cast iron surface. The respective cast iron plate was incubated for 30 days at 37 °C in anaerobic bentonite microcosms containing CAP rock solution and H₂. The cross-section is part of a sequence of cross-sections that allows viewing the 3D-shape of the sub-surface voids.

deep and "grow" into the metal surface, even beneath an intact surface (Fig. 1).

At first glance, the results of this work argue against an application of the cast iron as a technical barrier for a DGR. However, it has been shown, that Ni-alloyed cast iron is more corrosion resistant than unalloyed.^[6] Thus, testing different cast iron alloys might identify more suitable metals. Furthermore, a longer exposure time may thicken the rust layers, thereby retarding corrosion conversely.^[5] Indeed, Raman spectroscopy and SEM-EDX analyses of the here analyzed plates, revealed the presence of siderite (FeCO₃, data not shown), which could potentially form a passivating film on the cast iron surface.

In former experiments, dealing with the microbial influence on corrosion of cast iron, sulfate reducing bacteria (SRB) have been detected.^[7] Additionally, it has been shown that the bentonite applied here harbors SRB, which are active in OPA-containing microcosms supplemented with H₂.^[3] The microbial diversity analysis is still ongoing and might reveal a distinct role of detected microbial genera during the corrosion processes.

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When the biofilm on materials is needed: biofilm formation after electrostatically attachment of films of living bacterial cells

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The modification of the surfaces of cells or material by polyelectrolytes (PEs) allows the deposition of bacterial cells, biofilm forming and non-forming ones, in order to construct artificial biofilms with prolonged stability, causing the selective pressure on further colonization of environmentally occurring bacteria.

Bacterial attachment is crucial in many biotechnological applications. However, biofilms have deteriorative effects on materials. That is why the current strategies to manage biofilms are focused on biofilm inhibition. Knowing that biofilm formation is an inevitable process, instead of biofilm prevention, we propose to manage its microbial composition by formation of biofilms with beneficial microbes. Using an artificial biofilm prepared with the potentially probiotic strain of *Bacillus* sp., we were able to reduce the attachment and completely inhibit the growth of *E. coli* cells on the surface of the *Bacillus*-based artificially induced biofilm.

EXPERIMENTAL. Different bacterial strains were selected based on their intrinsic physiological properties. The strains were used for (*i*) testing the toxicity of the modified PEs, (*ii*) observing the effect of PE deposition on the formation and stability of biofilms and (iii) proving the concept of the protective nature of artificial biofilms. In preliminary experiments polyethyleneimine (PEI) demonstrated toxicity when exposed to bacterial cells. Consequently, it was tested if reducing the number of amino groups in the PE molecule would decrease its toxicity. We used acetic anhydride to acetylate PEI. Further, we tested the efficiency of biofilm formation for biofilm-forming and non-forming bacteria, coated with acetylated and non-acetylated PEI, incubating cells upside down over 48 hours. To observe how the artificial biofilm of Bacillus 25.2.M inhibited the attachment of a potentially pathogenic bacterium we firstly deposited acetylated PEI coated Bacillus 25.2.M on the glass surface and secondly E. coli on the top of Bacillus 25.2.M and left to incubate.

RESULTS. As acetylated PEI allowed to manipulate the cell surface charge without being toxic to the cell, we determined how the electrostatic deposition of cells to the glass surface enabled both the biofilm forming and non-forming bacteria to occupy the freshly exposed surfaces (Fig. 1). In contrast to the controls, where almost all the cells of the biofilm forming bacteria were detached after 48 hours, surface modifications by PEs resulted in 4.66 times larger occupied surface areas. Furthermore, the occupied surface areas were 26.02 times larger for the non-biofilm-forming strains. To test the effects of the artificial biofilm in an in vitro system containing two different bacterial species, we constructed the biofilm with Bacillus 25.2.M, a potential oral probiotic strain, and deposited the cells of E. coli over its surface (Fig. 2). In the presence of the artificial biofilm the cell number of E. coli TOP10 emGFP was constant over time and no significant growth was observed within the 24 hours exposure period. Here, we have demonstrated a new solution which does not prevent the inevitable establishment of the biofilm, but carefully manages its development by attachment of live bacteria. The modification of PEs widens the repertoire of strains that can be used in the purification systems used for drinking and



Fig. 1: The modification of the cell's surface charge or the charge of the material surface affects the number of cells that attach onto the surface and the following growth of the biofilm. The surface area covered by cells was examined for 8 different bacterial strains, biofilm forming or biofilm non-forming, over a 50 h period.



Fig. 2: The colonization and growth of potentially pathogenic *E. coli* are inhibited by the artificial biofilm made from the probiotic candidate strain *Bacillus* sp. 25.2.M. The number of attached cells of *E. coli* TOP10 emGFP, expressing the green fluorescent protein (gfp) (green), was determined at 0, 6 and 24 hours of cultivation when the artificial biofilm was made from the antimicrobial-producing *Bacillus* sp. 25.2.M, being deposited onto the glass surface (A–C) or from *E. coli* TOP10 (red) not producing gfp (positive control) (D–F), or no preapplied biofilm on the surface (negative control) (G–I). Scale bar equals 10 µm.

waste water treatment, and consequently the proposed approach can be applied to different surfaces in a variety of industries such as biotechnology, medical and other sectors, where prevention and control of biologically induced corrosion is important.

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The lab lysimeter - an application for biogeochemical research with radionuclides

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This report summarizes technical details and application of a lab lysimeter, which was developed and built by ecoTech Umwelt-Meßsysteme GmbH, Bonn, Germany, for a research project according to our requirements.^[1] The goal of the project is the identification of the underlying processes determining the migration behavior of radionuclides in a complex environment composed of soil, organic matter, microorganisms, and plants.

TECHNOLOGY. Lysimeters are measuring devices for water and material balance parameters and are used for the investigation of the interaction and accordingly mass transport between atmosphere, different organisms (such as plants, fungi, animals), soil and water in environmental research. Our lab lysimeter consists of a cylinder made of polypropylene, which consists of an upper transparent and a lower nontransparent part. The total height is 105 cm and the diameter is 30 cm (Fig. 1).



Fig. 1: Lab lysimeter system in the control area of the institute.

A sprinkling head equipped with light emitting diodes controls the water content of the soil matrix and the illumination during the growing of plants and fungi with an adjustable intensity in the range of 3,000 to 12,000 Lumen (Fig. 2). Furthermore, it is possible to connect up to 54 redox electrodes, pH electrodes, and suction candles arranged in 9 levels about each other at a distance of 5 cm between each level. Thus, 6 sensors can be installed at every level. Additionally, different tensiometers can be used for the determination of the soil water or soil moisture.

In the lower part of the lysimeter a suction bottom consisting of 7 segments and connected to a bottle system allows the





Fig. 3: Sevenfold segmented suction bottom (left) and bottle system (right).

collection of seepage water (Fig. 3). The controlling, data registration modules and a data transfer technology are integral parts of the lab lysimeter system.

APPLICATION. The main attention of the scientific work with the lysimeter is paid on the decent investigation of the interactions of radionuclides and metals with components of the soil and of the soil solutions. Furthermore, information on the horizontal and vertical distribution of the heavy metals in soils and their uptake by organisms under near natural conditions can be gained from the experiments.

With respect to the soil, measurements determining soil parameters and their change upon gradual adding of selected heavy metals and radionuclides as well as complexing organic agents will be carried out. Here, the focus is on the sorption of the metal ions onto the soil particles, the determination of the corresponding metal speciation, and the examination of the immobilization effects of the chelating agents on the radionuclides. The studies to the metal- and radionuclide behavior in soils and other porous materials will be connected with the performance of flow analyses, determination of flow paths and the hydrodynamics regarding the assessment of groundwater contamination by heavy metals.

Considering biological aspects, the influence of plants, fungi, and bacteria on the migration behavior and transport of metals and radionuclides in soils will be studied. The characterization of radionuclide uptake by the organisms at welldefined soil conditions is feasible, for instance by monitoring the biome composition using molecular biological methods. In analogy to performed batch experiments, the aim is to obtain information on radionuclide and metal speciation in the biomass cultivated in the lysimeter at molecular level using various spectroscopic and microscopic methods.

The outcome of the experiments with the lab lysimeter system is crucial for the transferability of data gained from small-scale experiments to larger scales and or the modeling of the radionuclide's behavior in the environment.

[1] https://www.ecotech.de.

Fig. 2: Lighting and sprinkling unit.

Liquid light guide technology augments use of small economical detectors in liquid scintillation counting

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In pursuit of smaller and more economic detectors in liquid scintillation counting (LSC), silicon photomultipliers (SiPM) seem to be promising candidates, as they are robust, can be produced in a small size, at low costs and require much lower operating voltages than conventional photomultiplier tubes.^[1] Here, we show that, for small scale SiPM detectors, liquid light guide (LLG) technology will enhance the collection efficiency for the scintillating light, giving rise to the analysis of sub-mL LSC samples or samples with low activity.

EXPERIMENTAL. 0.3 mL of 14C citric acid dissolved in ultima gold (~100Bq/mL) was measured in fluoro-ethylenepropylene (FEP) polymer hosing and in Eppendorf microvials made of polypropylene (PP). The scintillating pulses were detected with a 3×3 mm SiPM detector,^[2] attached to a printed circuit board with stabilized 30 V power supply. The SiPM is an avalanche photodiode in Geiger mode and was operated 5V above its breakdown voltage. The produced voltage pulses were amplified by a pre-amplifier to give ~60 ns voltage pulses of up to 500 mV. In the results shown, analog digital conversion (ADC) and mean pulse voltage analysis were performed by a 2.5 GHz oscilloscope with 8 bit resolution. Pulse width was measured additionally to act as a quality monitoring quantity, *i.e.* only pulses with widths between 40 and 150 ns were used for analysis. Due to the slow process of oscilloscope data acquisition and analysis (~5s per pulse) counting time was usually several hours. Hereby, low energy scintillation pulses were discriminated by setting the trigger level of the oscilloscope as well as setting the vertical voltage resolution. For all measurements, both quantities were selected to be well above the amplitude of the dark current pulses coming from the SiPM (< 10 mV). Typically, the selected range for the detected voltage pulses was between 50–450 mV and the trigger level was at 20 mV. The mean data of the pulses collected was then converted to a histogram and normalized with respect to counting time using OriginTM.

RESULTS. The enhancement of collecting scintillation light using FEP hosing is shown in Fig. 1. Herein, ¹⁴C is used as the radioactive source, which has been dissolved in the scintillator and inserted into the FEP hosing. Due to the high differences in refractive indices between FEP (1.34) and the scintillator (1.59) the system is highly LLG active. Therefore, a nearly tenfold increase in collection efficiency is observed compared to conventional, only slightly LLG active polypropylene tubes (refractive index ~1.48) with additional diffuse reflector.

In order to further test the sensitivity of the rudimentary liquid scintillation counter mentioned above, measurements were also performed with blank scintillation cocktail in FEP hosing with and without lead shielding and by having a radioactive source (plate of uranium glass) nearby.

The results are shown in Fig. 2. It can be seen that the uranium plate as well as the different environment are clearly recognized by the apparatus. Due to the small size of the SiPM, this could give rise to pen sized ionizing radiation monitoring devices.



Fig. 1: Mean histogram of ¹⁴C citric acid in Ultima Gold with lead shielding in FEP hosing and in PP microvial with reflector.



Fig. 2: Mean histogram of blank Ultima Gold in FEP hosing in different environments.

OUTLOOK. Currently, we are testing different ADC options, having 8 bit/2.5 GHz, 10 bit/2.5 GHz and 14 bit/125 MHz specifications. Especially for the latter case, a pulse shaper was also implemented in the setup. Additionally, two different multichannel analyzers are available, which allow hardware pulse analysis as well as automated histogram formation. By hardware pulse analysis, pulses can be captured at a much higher rate, which will shorten the measurement time. The automated histogram formation, on the other hand, will reduce the time for data analysis.

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SCIENTIFIC CONTRIBUTIONS (PART IV)

Nuclear Reactor SAFETY RESEARCH

The dominant mechanisms for the formation of solute clusters in steels exposed to irradiation

F. Bergner

Until recently, there was much debate about the dominant mechanisms of the formation of solute atom clusters in low-Cu reactor pressure vessel steels exposed to neutron irradiation. The crucial role of solute segregation to TEM-invisible irradiation-induced dislocation loops has now been uncovered by atomistic simulations supported by nm-scale experimental evidence.

Fast neutrons give rise to displacement cascades in reactor pressure vessel (RPV) steels. As a direct result, a high number of point defects is produced, namely essentially equal amounts of vacancies and self-interstitial atoms. At reactor-ambient temperatures, both vacancies and interstitials migrate in the body-centered cubic Fe lattice of low-alloy RPV steels. Interstitials migrate faster. Under continued irradiation they reach a quasi-steady state concentration at the μs time scale. For vacancies the quasi-steady state is reached at the ms to s time scale. The long-term evolution of the microstructure covers the following processes:

- recombination of a vacancy and an interstitial,
- clustering of vacancies to form three-dimensional sub-nm voids,
- clustering of interstitials to form two-dimensional dislocation loops.
- Loss of vacancies and interstitials in microstructural sinks such as dislocations and grain boundaries.
- Binding of point defects with solute atoms to form complexes often also capable of migrating in the lattice.

Given this background, a huge variety of possible mechanisms for the rearrangement of solute atoms as well as the formation of solute clusters and stable precipitates is theoretically possible. The question is, which mechanism(s) is (are) dominant in RPV steels exposed to reactor neutron irradiation.

For Cu-rich first- and second-generation RPV steels, the answer was already anticipated in the 1980s.^[1] The dominant mechanism was essentially confirmed and worked out in more detail in the decades that followed. It was found that the increased vacancy concentration gives rise to irradiationenhanced Cu diffusion resulting in the formation of coherent Cu-rich vacancy-containing precipitates. At longer time scales, alloying elements such as Mn and Ni do contribute as well.

For low-Cu second- and third-generation RPV steels, the problem turned out to be more puzzling. The same is true for Cu-rich RPV steels irradiated at such high fluences that the remaining matrix-Cu already went below the solubility limit. Collaborative international efforts including multiscale modelling and complementary nm-scale experimental approaches^[2] recently allowed the dominant mechanisms of cluster formation to be uncovered as described in brief below:

- (1) Interstitials cluster together to form dislocation loops. In pure Fe (but not in RPV steels), these loops grow to TEM-visible sizes.
- (2) The sub-nm sized loops are mobile in one dimension. Upon loop migration, they encounter solute atoms which bind to the loops and immobilize them. The loops are now essentially unable to grow further.
- (3) Solute atom migration via vacancy or interstitial drag gives rise to an enrichment of loops with solute atoms. Such segregated loops are often observed by means of atom probe tomography (APT), see Fig.1 for an example.^[3]



- Fig. 1: Segregated half-toroid shaped loop reconstructed from an APT measurement on a neutron-irradiated low-Cu A533B-type RPV model steel.^[3] The probable position of the original dislocation loop was superposed to the as-measured APT map.
- (4) Upon further segregation, the loops lose their toroidal shape to become three-dimensional solute clusters or even stable precipitates as evidenced by APT and smallangle neutron scattering (SANS).

The mechanisms described above in (1)-(4) and the mechanism outlined before for Cu-rich steels may operate in parallel, and, if this is the case, influence one another. It is important to note that the segregated-loop mechanism of cluster formation is essentially irradiation-induced as opposed to irradiation-enhanced, that means it would not operate under thermal conditions.

Evidence for the multistep process was derived from objectkinetic Monte-Carlo simulations built upon a carefully parameterized model. The simulation results were compared with experimental results obtained for the mean size, number density, volume fraction and composition of clusters that formed in numerous neutron-irradiated RPV steels.^[2]

Another interesting point is the time scale at which the vacancy-Cu-based mechanism and the segregated-loop mechanism operate. There is experimental evidence indicating that the former process is effectively faster. On the other hand, once the impurity Cu in excess of the solubility limit is consumed, no or relatively few new clusters will form via this mechanism. In contrast, the segregated-loop mechanism is mainly based on solute atoms of alloying elements such as Mn, Ni and Si, which are available in RPV steels practically without limit to form Mn-Ni-Si-enriched clusters. A consequence of these competing mechanisms is a relatively wide variability of the functional dependence of the total cluster volume fraction on neutron fluence.

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Use of the small punch test for the estimation of ductile-to-brittle transition temperature shift of irradiated steels

E. Altstadt, F. Bergner, M. Houska

The correlation between ductile-to-brittle transition temperatures (DBTTs) obtained from the small punch test technique and from the standard Charpy impact test is investigated. Tests at different temperatures were performed on various steels including materials from original reactor pressure vessels in the unirradiated, neutron irradiated and annealed condition. It is demonstrated that the small punch test is a reliable and effective supportive means for the estimation of the irradiationinduced shift of the ductile-to-brittle transition temperature.

The small punch (SP) test has been recognized as a possible replacement of the Charpy impact test in cases of insufficient amounts of material available for the fabrication of the necessary number of Charpy samples. As a screening procedure, the SP test is intended to provide estimations of the material properties with small amounts of material, allowing an evaluation of ageing mechanisms such as irradiation-induced hardening and embrittlement.^[1,2] The shift of the DBTT is a key element of reactor pressure vessel (RPV) embrittlement surveillance. It is therefore interesting to evaluate whether the DBTT shift can be reliably measured by means of the SP test. As for the SP-test-based estimation of the DBTT, T_{SP} , the following simple correlation is used:

$$T_{SP}[K] = \alpha \cdot T_{CVN}[K] \tag{1}$$

where T_{CVN} is the DBTT from the Charpy impact test using standard specimens $10 \times 10 \times 55$ mm with a 2 mm V-notch (socalled CVN specimens). The factor of proportionality α depends on SP geometry, the material class, the fitting procedure for T_{SP} , and the criterion for the definition of T_{CVN} (*e.g.*, T_{41J} , T_{47J} , T_{68J} , $T_{50\%\text{US}}$).

EXPERIMENTAL. A number of seven different steels were investigated. These materials include bainitic RPV steels, base metals as well as weld metals. The experimental program covers different irradiation conditions, namely unirradiated (u), neutron irradiated (i) and neutron irradiated and annealed (ia). CVN specimens were tested according to the standard ISO 148-1. The test temperature range was from -25 to 225 °C for the irradiated steels and from -150 to 150 °C for the unirradiated steels. The impact energy Wt was calculated as integral of the force-deflection curve F(s). SP specimens of $10 \times 10 \times 0.54$ mm were manufactured from tested or unused Charpy-sized SE(B) specimens. In total, more than 500SP tests were performed. The test temperatures were in the range from -172 to +100 °C for the irradiated materials and from -194 to +100 °C for the unirradiated materials. For each test, the force-displacement curve F(v) was recorded and the small punch energy Esp was calculated as integral of F(v) up to maximum force F_m . The normalized small punch energy was obtained by $E_n = E_{SP}/F_m$. The DBTTs of the Charpy impact and small punch test series were determined by means of a tanh-fit of the energies W_t and E_n as function of the test temperature. T_{471} is the temperature at which the tanh-fit of W_t takes the value of 47 J, while T_{SP} corresponds to the inflection point of the tanh-fit of E_n .

RESULTS. Figure 1 shows the correlation between T_{SP} and T_{CVN} (T_{471}). The coefficient of determination is $R^2 = 0.986$.



Fig. 1: Correlation between E_n based T_{SP} (tanh-fit) and T_{47J} .

Considering that seven different steels are involved, the correlation Eq.(1) works surprisingly well. The application of the $T_{SP}-T_{47I}$ correlation to the core weld of the RPV of NPP Greifswald 4 is shown in Fig. 2. For the as-irradiated condition, a characteristic through-thickness dependence of T_{471} is observed in the Charpy impact test, which exhibits a pronounced minimum close to the welding root. This is nicely reproduced with the SP-based T_{SP} for the as-irradiated condition. The recovery of mechanical properties due to the annealing treatment is clearly indicated by significantly lower DBTTs (Fig.2). Again, the Charpy- and SP-based values of T_{47J} give a coherent picture. In the annealed condition, there is no significant difference between the DBTT for the welding root and the filling layers. This is in contrast to the asirradiated condition. At a first glance, the particular throughthickness behavior seems to be unexpected, as the neutron fluence is decreasing from the inner surface to the outer surface. However, this finding is well in line with the throughthickness variation of the concentration of Cu, which is the most important promoter of irradiation-induced embrittlement. Indeed, the Cu concentration was reported to be lowest in the welding root.[3]



Fig. 2: DBTT T_{47J} obtained from Charpy impact tests and from SP test, recalculated with Eq. (1), as a function of position in the core welding seam of the RPV of Greifswald unit 4 (material 10KhMFT); as-irradiated (irr) and annealed (ia) conditions; the background picture shows the structure of the weld.

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Microstructural characterization of inhomogeneity in 9Cr ODS EUROFER steel

P. Chekhonin, A. Das, E. Altstadt, F. Bergner, C. Kaden

Inhomogeneous regions of an oxide dispersion strengthened (ODS) steel were investigated. It was demonstrated that optical microscopy alone is insufficient to identify these inhomogeneous regions as residual ferrite and advanced techniques based on electron microscopy as well as site-specific nanoindentation are required.

Given their enhanced resistance against creep and hightemperature strength, ODS steels are interesting candidate materials for cladding tubes in Gen-IV nuclear fission reactors and blankets for fusion reactors. The presence of residual α -ferrite, the phase which remains untransformed during austenitization, was often reported in 9 wt.-% Cr ODS steels. Typically, residual ferrite is identified as bright, smooth and somewhat elongated region (bright stripe) against a dark matrix under optical microscopy (OM). Recently, there was confusion regarding similarly looking bright regions that are attributed to the absence of nanoparticles.

The present investigation focuses on electron microscopy to clarify this issue in the case of the air-cooled EU batch of ODS EUROFER.

EXPERIMENTAL. Samples of ODS EUROFER have been analyzed by means of OM, nanoindentation, scanning electron microscopy (SEM), electron backscatter diffraction (EBSD) and scanning transmission electron microscopy (STEM). Focused ion beam (FIB) preparation was used to select specific sample areas for subsequent TEM analysis.

RESULTS. Optical microscopy on etched samples indicates the occasional presence of bright stripes. Their volume fraction is estimated to be approximately 4%. SEM investigations demonstrate a lower carbide density (Fig. 1) inside bright stripes, while EBSD mappings confirm a body centered cubic steel phase and indicate a larger grain size. Nanoindentation experiments prove that the bright stripe region is softer in comparison to the matrix. Applying the FIB technique a thin lamella was extracted containing both, a bright stripe region as well as the adjoining matrix (Figs. 2a and 2b). STEM investigations reveal the absence of nanoparticles inside the bright stripe (Fig. 2c), while many nanoparticles are recognizable inside the matrix (Fig. 2d). Further-



Fig. 1: Backscatter electron contrast micrograph with a bright region and adjoining matrix regions. The carbides appear as bright spots.



Fig. 2: SEM micrograph showing the FIB lamella taken out from the EU batch of ODS EUROFER including areas from a bright stripe and the adjoining matrix region(a). STEM bright field image of the lamella including both the regions (b). STEM bright field image from the bright stripe region (c), and from the matrix region (d).

more, STEM indicates a significantly lower dislocation density of approximately an order of magnitude, thus, qualitatively as well as quantitatively explaining the lower yield strength in the bright stripes.^[1]

These results present clear evidence that in this ODS steel the bright stripes cannot be residual ferrite or retained austenite. Instead, these bright stripes are most probably formed due to improper mechanical alloying (MA) resulting in the absence of nanoparticles in certain regions. This gives rise to higher grain boundary mobility, larger grain size and lower number density of grain boundary carbides, which finally is believed to result in the distinct etching behavior as observed by optical microscopy. An extension of the MA period during manufacturing should lead to a homogenous microstructure, eliminating particle-free regions which are not expected to contribute to the creep resistance and high temperature strength of ODS steels.

Finally, it should be pointed out that these inhomogeneities do not only affect the yield strength, but additionally lead to a reduced fracture toughness, particularly at temperatures at and below room temperature,^[2] in comparison to differently produced ODS steels that exhibit more homogeneous MA. More details on the presented investigation including an advanced discussion can be found in reference [1].

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Microstructure and fracture toughness characterization of three 9Cr ODS EUROFER steels with different thermo-mechanical treatments

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Ferritic martensitic oxide dispersion strengthened (ODS) steels are promising structural materials for future Gen-IV nuclear fission and fusion reactors. The dependence of fracture toughness on microstructure was investigated. For this, three 9Cr ODS EUROFER steels manufactured through different thermo-mechanical processing routes were compared using fracture mechanics testing and microstructural characterization. It was found that variations in thermomechanical treatments can produce significant differences in microstructure and fracture toughness behavior of ferritic martensitic ODS steels.

Systematic development of reduced activation ferriticmartensitic (RAFM) ODS steels as structural material for nuclear applications in Europe resulted in ODS EUROFER, which can withstand operating temperatures up to 650 °C and exhibits high-temperature creep strength and resistance towards irradiation induced effects due to the oxide particles dispersed in the microstructure. ODS EUROFER batches were manufactured via different thermo-mechanical routes and hence developed slight variations in microstructure. The fracture toughness of a material is directly related to the microstructure, and is an extremely important parameter affecting the loading capacity, machinability and formability. Therefore, it becomes important to study the impact of variation in microstructure on the fracture toughness of 9Cr ODS EUROFER steels.

EXPERIMENTAL. The nominal composition of ODS EUROFER is (wt.-%) 8.9 Cr, 1.1W, 0.2V, 0.14Ta, 0.42Mn, 0.06Si, 0.11C, 0.3Y₂O₃ and Fe for the balance. The production route of all the three ODS EUROFER batches involved mechanical alloying followed by hot isostatic pressing, hotrolling at 1150 °C followed by austenitization at 1100 °C for 30 min. Hereafter, slightly different heat treatments were followed as indicated in Tab. 1. Tempering was then performed for all the materials at 750 °C for two hours.

Tab. 1: Heat treatments of the three ODS EUROFER batches.

Batch Name	Heat Treatments	Ref.
ODS_WQ_450	water quenching to 450 °C	[1, 2]
ODS_WQ_RT	water quenching to RT	[1, 3–5]
ODS_AC_RT	air cooling to RT	[6]

Microstructure characterization of the bulk was performed using OM, SEM, EBSD and TEM. Quasi-static fracture toughness testing was performed at -100 °C to 600 °C using sub-sized compact tension 0.25T C(T) specimens. The fracture surfaces were characterized using SEM.

RESULTS. After austenitization, ODS_WQ_450 was water quenched to 450 °C, which is above the martensite start temperature. Hence, its microstructure consists only of coarse globular ferrite grains. ODS_WQ_RT and ODS_AC_RT exhibited martensitic microstructure with ODS_WQ_RT containing residual α -ferrite phases which are harder than the matrix. ODS_WQ_450 contained larger sub-micron particles enriched with Cr, W, C and O while smaller particles were observed in the other two materials.

At low temperatures (-100 °C and RT), ODS_WQ_RT exhibits higher fracture toughness and more stable fracture behavior as compared to ODS_AC_RT (Fig.1) due to higher amount of secondary cracking which may be explained by the mismatch of mechanical properties between the elongated hard residual ferrite phases and the soft tempered martensite matrix. At these temperatures, ODS_WQ_450 exhibits the lowest fracture toughness because it contained the largest sub-micron particles. No information about the occurrence of secondary cracks on the fracture surfaces could be obtained from literature.^[2]



Fig. 1: Fracture toughness comparison over a wide temperature range for ODS_AC_RT and ODS_WQ_RT along with ODS_WQ_450.^[2] For ODS_WQ_450, the results for -100 °C and 22 °C are taken from Lucon et al. (grey symbols).^[1] The dashed line indicates the acceptable fracture toughness value of $J_Q = 45$ kJ/m² ($K_{JQ} = 100$ MPa \sqrt{m}).^[5]

At temperatures between 100 °C and 500 °C, ODS_WQ_RT and ODS_AC_RT exhibit more or less similar fracture toughness as the matrix ductility of these two materials is similar. On the other hand, ODS_WQ_450 exhibits the highest fracture toughness which can be explained by its higher matrix ductility.

At temperatures above 500 °C, reduced fracture toughness values for ODS_WQ_RT and ODS_AC_RT were observed due to a shift from transgranular ductile to intergranular fracture. On the other hand, ODS_WQ_450 still exhibited higher fracture toughness as its coarse-grained microstructure behaved better against intergranular fracture.

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Relationships between primary radiation damage, irradiation-induced microstructure and hardening of ion-irradiated Fe-Cr and ODS Fe-Cr alloys

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Ion irradiations are a powerful tool for investigating radiation effects like radiation hardening. However, the extraction of radiation hardening as a function of displacement damage from the nanoindentation (NI) response of ion-irradiated materials is a challenge. In this study, we analyzed the microstructure and the NI response for two ion-irradiated Fe-Cr alloys. Crosssectional Transmission Electron Microscopy (TEM) indicated the appearance of irradiation-induced dislocation loops but no other types of visible microstructural changes. NI indicated maxima of radiation hardening as a function of contact depth. Links between the depthresolved primary radiation damage, the observed depthdependent characteristics of the loops and the measured hardening are considered.

The irradiation with MeV ions gives rise to a depth-dependent displacement of atoms from their lattice sites, a deposition of injected ions at interstitial sites of the crystal lattice and a formation of irradiation-induced defect clusters such as dislocation loops. The consequence of these effects is a degradation of the mechanical properties in terms of embrittlement and hardening. One of the long-term goals of ionirradiation studies is the extraction of radiation hardening as a function of displacement damage. In our study, the effect of 5 MeV Fe-ion irradiation of Fe-Cr alloys on microstructure and hardening is addressed. Empirical models are applied to link primary radiation damage with the irradiated microstructure and the microstructure with irradiation hardening.^[1]

EXPERIMENTAL. The materials studied are an Fe-Cr model alloy (8.4 wt.-% Cr) and an oxide dispersion strengthened (ODS) Fe-Cr (9.1 wt.-% Cr) alloy. The ion irradiations were performed with the 3MV tandetron accelerator of the Ion Beam Center at HZDR. Fe2+ ions of 5 MeV energy were implanted at a temperature of 300 °C. The depth profiles of displacements per atom (dpa) and injected interstitials per atom (ipa) were calculated using the binary collision code SRIM-2013.00. The microstructure was studied by cross-sectional TEM using a FEI Talos F200X microscope. Irradiationinduced dislocation loops were imaged under various diffraction conditions. Depth profiles of loop number density and size were obtained by counting and sizing the loops in TEM bright-field images. The sample thickness was estimated by means of convergent beam electron diffraction. NI hardness was measured using a Universal Nanomechanical Tester (UNAT, ASMEC GmbH). Hardness vs. depth plots were calculated from about 30 indentation tests by averaging.

RESULTS. For both materials, the TEM images indicate the appearance of irradiation-induced dislocation loops as the one and only type of visible microstructural changes. Figure 1a displays a TEM bright-field overview image of the ion-irradiated ODS FeCr sample. A defect-rich band aligned parallel to the specimen surface is visible in the depth range from 1.2 to $1.6 \,\mu$ m, which is close to the peak positions in the dpa- and ipa-profiles (not displayed here). The band contains a large number of dislocation loops appearing as black dots under two-beam kinematical bright-field conditions and as white dots under weak-beam dark- field conditions as shown in Figs. 1b and 1c. The depth-dependent number



Fig. 1: TEM images of ion-irradiated ODS FeCr. Overview (a), two-beam kinematical bright-field (b) and weak-beam dark-field image (c) of the area marked in (a).



Fig. 2: Indentation hardness measured for the unirradiated and irradiated state of ODS FeCr.

density of dislocation loops determined from the TEM images was decomposed into fractions arising from displacement damage and injected interstitials (see ref. [1] for details). The latter fraction corresponds to 2 % of the total injected interstitials for ODS FeCr.

For both materials, a maximum of the irradiation-induced hardness increase as a function of contact depth was found as demonstrated in Fig. 2 for ODS FeCr. This observation is rationalized by a combination of the dispersed barrier hardening model and the half-sphere approximation of the indentation plastic zone, with the amount of the predicted hardening depending on the selected superposition rule of the contributions from pre-existing and irradiation-induced barriers for dislocation glide.

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Vacancy-solute clustering in Fe-Cr alloys after neutron irradiation

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Vacancy-solute clustering in neutron irradiated Fe-Cr alloys with various concentrations of Cr and minor solutes (Ni, Si and P) were studied by using positron annihilation spectroscopy (PAS) and small-angle neutron scattering (SANS). The results indicate the existence and formation of vacancy-CrNiSiP clusters that play a detrimental role in irradiation hardening. In Fe9Cr alloys with significantly reduced concentration of minor solute elements, the main defects are vacancy clusters with a cluster size up to 10 vacancies. A high concentration of α' -precipitates was observed in Fe14Cr(NiSiP). However, both vacancy clusters and α' -precipitates provide significantly less impact to hardening in comparison to vacancy-CrNiSiP clusters.

High Cr ferritic/martensitic (F/M) steels are candidate structural materials for the construction of key components in future reactors, since they offer better radiation resistance in comparison to austenitic steels. In order to provide a solid interpretation of the mechanical properties and to verify the association between solutes and point-defects, a thorough microstructure examination is needed. The defect characterization by PAS and SANS shows that vacancy-CrNiSiP clusters are formed and are dominant in irradiation hardening.

EXPERIMENTAL. The investigated materials are listed in Tab. 1. The composition of materials, the performed neutron irradiation up to 0.11 displacements per atom (dpa) and detailed descriptions of the PAS and SANS experiments are given in [1] and references therein. The A-ratio is defined as the ratio of total and nuclear irradiation-induced scattering. It is a one-parameter signature of the cluster composition.

Tab. 1: FeCr model alloys, volume fraction *c*, and *A*-ratio of irradiationinduced nano-features derived by SANS and changes of Vickers hardness ΔHV10 after irradiation (0.11 dpa).

Material	c (vol%)	A-ratio	$\Delta HV10$
Fe9Cr F/M	0.13 ± 0.01	2.20 ± 0.2	27 ± 3
Fe9Cr	0.15 ± 0.01	1.45 ± 0.1	22 ± 3
Fe9CrNiSiP	0.29 ± 0.01	1.90 ± 0.15	72 ± 3
Fe5CrNiSiP	0.13 ± 0.01	3.50 ± 0.5	68 ± 4
Fe14CrNiSiP	1.75 ± 0.05	2.15 ± 0.15	94 ± 7

RESULTS. All irradiated materials exhibit an increase of the S-parameter and a decrease of the W-parameter in the PAS experiment, indicating that positrons are trapped and annihilate at vacancy-rich complexes induced by neutron irradiation. Fe9Cr ferrite has the strongest enhancement in the low momentum region of all materials. In the high momentum region, the smallest difference with respect to the reference Fe sample was observed for Fe9Cr F/M alloy. The coincidence Doppler broadening spectra of all other alloys fall in between these two extreme cases (Fig. 1).

The size distribution of irradiation-induced clusters determined by SANS are shown in Fig 2. An average radius of about 1 nm was estimated for all alloys containing minor solutes. The *A*-ratio (Tab.1) observed for ferritic pure Fe9Cr indicates that the majority of irradiation-induced scatterers are vacancy clusters (theoretical value 1.4). The higher *A*-ratios of the other alloys indicate a dominant contribution



Fig. 1: PAS results: S vs. W parameter. Non-irradiated condition (a), neutron irradiated FeCr alloys (b).



Fig. 2: SANS results: size distribution of irradiation-induced nano-features in different Fe9Cr alloys after the same neutron irradiation up to 0.11 dpa.

of solute clusters. It is worth to note that an *A*-ratio of 2.05 is expected for Cr-rich α '-phase particles, as observed in Fe14CrNiSiP. Moreover, the high-volume fraction in Fe14CrNiSiP provides additional argument for the presence of Cr-rich α '-phase particles in this alloy.

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Characterization of the porous network structure of concrete

Q. I. Roode-Gutzmer, J. Kulenkampff

The porosity and specific surface area of a small intact core of concrete (~0.2 g) was characterized nondestructively by micro-computed tomography (μ -CT), N₂/BET and destructively by Mercury Intrusion Porosimetry (MIP).

Biological shielding concrete that has been subjected to neutron radiation is prone to degradation *via* the alkali silica reaction.^[1] This has two implications: loss of mechanical strength (which impacts service life), and a change in the porous structure and reactive surface in terms of the sorption characteristics and transport of radionuclides (either occurring *in situ via* neutron activation or *ex situ* due to contamination of leaked reactor cooling water). The objective of this work is to investigate the characterization of pore structure on intact specimens of concrete (as small as possible) so that reactive transport processes can be directly correlated.

EXPERIMENTAL. Non-radioactive concrete (ca. 48 years old) from the dome section of a reactor pressure vessel forming the base of the reactor was acquired during the dismantling of a decommissioned nuclear power plant, Stade, Niedersachsen, Germany. A diamond-coated core driller (dry; 200 rpm) was used to remove a small core ($\emptyset = -5 \text{ mm}$; L = 2 mm; 0.17 g) from a larger drill core (\emptyset = 73.6 mm; L = 150.7 mm) that remained intact and without any visible cracks after being subjected to a standard test (EN 13412:2006-11) for modulus of elasticity under compression (21.7 GPa). This small core was analyzed nondestructively by µ-CT and Scanning Electron Microscopy (SEM) (sample surface: rough, uncoated), then dried at 100 °C overnight prior to an N₂/BET measurement in triplicate. MIP was subsequently conducted on the small intact core in accordance with ISO 1509-1. µ-CT scans were acquired on the large and small drill cores using a Nikon XT H225 employing a tungsten target. A SEM image was obtained using a JEOL JSM-F100. The Pascal 140 and 440 Series from Thermo Electron Cooperation were used to conduct MIP. Specific surface area determined by N₂/BET was performed using a Thermo Scientific Surfer Analyzer.

RESULTS. The results are summarized in Figs. 1 and 2. It is clear that the methods of MIP and µ-CT yield pore-size information on different scales with a small overlap in the region 4–75 µm. The specific surface area calculated using the cylinder and plate model from the MIP results is 9.94 m²/g, whereas N₂/BET yields a value of $6.07 \text{ m}^2/\text{g}$ (s. d. = 0.54). A higher value for MIP is due both to the ink-bottle effect (large pores with narrow throats)-smaller pores are overestimated at the expense of larger pores-and to fracking. The hysteresis of the mercury intrusion and extrusion curves suggests needle-like pores.^[2] The initial constant volume with decreasing pressure and the fact that 51.5% of the mercury remains in the sample after extrusion is indicative of ink-bottle pores, as well as non-equilibrium compression and decompression conditions. Larger cores (i.e. ~1g) revealed similar pore features in their differential pore-size distribution curves, but cores < 0.1 g displayed erroneous results. While MIP yields information on the connectivity of the transport pores, the specimen is unfortunately unusable for further reactive transport experiments.



Fig. 1: Differential volume pore-size distribution curve (a) of a small intact core of concrete determined by mercury intrusion porosimetry illustrating on the logarithmic scale gel pore features (5–17 nm) and capillary pores of ~1 μm, which can clearly be seen in the SEM image of the cement matrix (b) and a thickness map (c) of a sub-volume of the small concrete core obtained from μ-CT portraying air voids (yellow) in millimeter range, and micropores of ~40 μm (red) and ~20 μm (blue).



Fig. 2: Accumulated porosity vs. pore-size on the logarithmic scale determined by MIP and μ -CT. True air voids and pores are calculated with background greyscale values (inset: pink) for a sub-volume of the real small core (with voxel size 4 μ m) (a) and a digitally extracted volume from the big drill core including the small core as sub-volume (with voxel-size 75 μ m) (b). Between the cement matrix and immediately adjacent to the aggregate exists a greyscale that is neither air nor grain (inset: red).^[3] This greyscale is not well resolved and could suggest obscure pores, possibly in a network of unresolved needles. This volume fraction is illustrated for the real small core (c).

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Volume expansion of quartz aggregate in ion-irradiated concrete

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Quartz aggregate in concrete irradiated by Si-ions with a fluence of 5×10^{14} ions/cm² at 300 keV exhibited an out-of-plane expansion of ~80 nm.

Quartz is the primary mineral component of most concretes and exhibits the highest neutron radiation-induced volume expansion (*i.e.* 17.8%) of all minerals yet examined.^[1] The volume expansion is due to structural relaxation that occurs when a critical concentration of defects (E' centers) has been produced by the interaction of the radiation imparted to the crystal lattice.^[2] The mechanism of this amorphization process appears to be indifferent to the type of radiation (*i.e.* neutrons, electrons, ions, and photons, e.g. gamma). The chemical durability of concrete is largely dependent on the reactivity of the silicate aggregates to alkaline pore water, which ultimately results in the alkali-silica-reaction. The rate-determining step for the dissolution of quartz is associated with $Q^4 \rightarrow Q^3$, which is considerably slower relative to the subsequent dissolution steps $Q^n \rightarrow Q^{n-1}$, where the Q^n nomenclature represents the Si-tetrahedron with *n* being the number of Si-atoms the tetrahedron is bound to.^[3] As quartz amorphizes due to radiation, the concentration of Q³ species increases. The ultimate goal of this work is to examine how radiation-induced changes in quartz effects dissolution behavior in alkaline media and to resolve (spatially and temporally) the influences of grain boundaries, defects, crystal strain, and Q³ content. In this paper we examine the volume change of quartz.

EXPERIMENTAL. Non-radioactive concrete (ca. 48 years old) from the dome section of a reactor pressure vessel forming the base of the reactor was acquired during the dismantling of a decommissioned nuclear power plant, Stade, Niedersachsen, Germany. Two slices (ca. 5 mm thick) of concrete from a drill core were obtained using a diamond-coated and water-cooled circular saw. Four polished sections with a diameter of 3 cm were prepared by vacuum impregnation using a two-component epoxy-based resin (Araldite 2020). The samples were ground smooth using SiC paste with a final diamond polish (3 µm). Mineralogical composition of the aggregates was identified in situ by µ-Raman spectroscopy (HORIBA Jobin Yvon LabRAM Aramis Vis) using an excitation wavelength of 532 nm. Aluminium foil was placed on each polished section so as to respectively protect and expose parts of the aggregate to radiation. Si-ion irradiation with a fluence of 5×1014 ions/cm2 at 300 keV was conducted under vacuum at room temperature without cooling. Some sputtering could be observed on the samples after irradiation. Depth profiles of the irradiated samples were examined by Vertical Scanning Interferometry (VSI) and Confocal Microscopy using an S Neox 3D Optical Profiler (Sensofar Metrology).

RESULTS. The minerals making up the aggregate were clearly identified by μ -Raman spectroscopy to be α -quartz and potassium feldspar (microcline: KAlSi₃O₈), the latter always intergrown with quartz. The results are illustrated in Fig. 1. The most significant depth profile change observed in the sample is the ablation of the epoxy resin (-600 nm). The volume change in quartz is observed as an increase in height of 80 nm. Taking into account an ion-penetration depth of 429 nm (estimated in the Kinchin-Pease calculation by SRIM)^[4] for pure quartz, the volume expansion amounts to



Fig. 1: Image showing radiated area (epoxy resin tarnished brown) left of the red dashed line and non-radiated area (epoxy resin clear) right of the dashed red line. Area of interest (AOI) in the white square depicting the boundary between radiated and non-radiated regions on large quartz aggregate. A = quartz (SiO₂), B = feldspar (microcline: KAlSi₃O₈), C = sputtering from Si-ion irradiation. VSI of AOI indicating depth profile, exhibiting +80 nm on quartz, -600 nm on epoxy resin, and unspecified on feldspar.

15.7%. No observable difference in average height between radiated and non-radiated areas on the black aggregate containing feldspar could be ascertained. Aggregates containing feldspar were generally found to display a rougher surface than aggregates containing only quartz. The maximum volume expansion observed for potassium feldspar has been reported to be 7.7 %.[1] Considering the same extent of radiation damage as observed for the quartz aggregate in this work, and a calculated ion-penetration depth for KAlSi₃O₈ (estimated by SRIM) of 604 nm, a height increase in microcline of ~41nm is expected. The troughs that the rougher surface of microcline exhibits could accommodate the expansion laterally in the plane, which decreases the expected out-of-plane expansion to ~3.4 nm, which is close to the resolution for VSI (1–2nm). In addition, curvature of granular surfaces deflects light, which inhibits VSI measurements. Finer surface polishing of the specimen is therefore required in order to accurately observe radiation-induced volume expansion of microcline.

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Microscopic characterization of un-irradiated and irradiated steel shielding material of a nuclear power plant

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Nuclear power plants (NPPs) have been considered an important energy source in Europe over a long time. However, in Germany the government decided that NPPs must terminate their operation end of 2022 at the latest, due to their possible threats on humans and the environment. The question arises, how much the long term irradiation of an NPP poses safety risks, in particular, in the reactor pressure vessel (RPV), which is the main shielding barrier of the radioactive fuel. For this purpose, the microstructural features, the composition, and the radioactive inventory of the steel shielding material in the RPV are investigated.

Figure 1 shows a photograph of a dismantled RPV of the decommissioned VVER prototype reactor from Greifswald NPP.^[1] The region next to the reactor core in the RPV is normally the part exposed to the highest neutron fluencies.^[2] Its steel mainly consists of both base metal and weld metal, which needs to be investigated in understanding the irradiation impact on the RPVs. Here, blocks of base metal and weld metal were taken from Greifswald NPP unit 8 RPV. It was produced by ŠKODA steel works in former Czechoslovakia by the end of 1970s, however it never went into normal operation mode.^[3]

EXPERIMENTAL. Specimens were extracted from the inner core of the RPV and cut into segments. These were neutron irradiated in JRC-Petten laboratories for a period of time 271 days, which mimics the long-term irradiation process. Analysis followed of both un-irradiated and irradiated samples, which were sliced into 10×10×1 mm sections and underwent a polishing process. Optical emission spectroscopy (OES) measurements on the un-irradiated samples provided the chemical composition of the alloying elements, which are sensitive to the neutron irradiation. Furthermore, surface analysis based on scanning electron microscopy/energy dispersive X-ray spectroscopy (SEM/EDX) of both un-irradiated and irradiated samples revealed microstructural features related to their composition and irradiation conditions.

RESULTS. Table 1 contains accurate mass measurements of the alloying material within the RPV shielding steel samples based on OES experiments. Results from the SEM/EDX measurements are shown in Fig. 2, they document an increase of the dark cluster size after the irradiation process and a distribution of white precipitates. Based on the EDX analysis, our results revealed the presence of manganese(II) sulfide (MnS) inclusions within the dark clusters, and the presence of molybdenum in the white distributed precipitates. Such a phenomenon can be explained as the long term irradiation on the steel shielding material may give rise to enhanced diffusion and rearrangement of the solute atoms in the Fe lattice. This consequently leads to the formation of clusters and segregation of the atom particles.

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Fig. 1: RPV of the decommissioned Greifswald NPP.^[1]



Fig. 2: SEM measurements of the base metals from the RPV showing their bainitic ferrite microstructural features of un-irradiated (a) and irradiated (b) steel shielding samples. The latter reveals the presence of high surface area of dark clusters containing MnS inclusions, and the distribution of white precipitates of Molybdenum.

Tab. 1: Mass concentration (%) of the chemical composition identified in the base metal of shielding steel samples in the RPV by OES.

	С	Si	Mn	Cr	Ni
%	0.15	0.26	0.47	2.87	0.07
SD*	0.006	0.003	0.003	0.01	0.0006
	V	Мо	Cu	Fe	
%	0.29	0.72	0.048	95	
SD*	0.002	0.008	0.0002	0.007	

*: denotes the standard deviation.

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Neutron fluence calculations and validation based on metal foils placed inside a German PWR

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In 2011, the German Atomic Energy Act was amended to phase out nuclear energy for the commercial production of electricity. As a result, all German nuclear power plants (NPPs) will have been shut down by the end of 2022. In order to dismantle the NPPs in a safe, economical, and timely manner, decommissioning studies have to be performed. The planning and implementation of a decommissioning strategy for an NPP require knowledge of the neutron activation and contamination levels, which have emerged during its operation. The extent and levels of activation in an NPP depend on the neutron fluence distribution within the NPP components. The neutron fluence is, therefore, one of the most fundamental and important parameter in the decommissioning studies. Hence, it should be as accurately estimated as possible. In this study, a 3D detailed Monte-Carlo (MC) model of a German PWR was developed to calculate the neutron fluence characteristics within the reactor components. The neutron fluence prediction capability of the developed model was validated based on activation measurement of several monitors composed of thin metal foils placed along the ex-core instrumentation channel. A very good agreement was obtained between the calculations and the measurements.

REACTOR MODEL AND MC CALCULATIONS. The German PWR was simulated using the Monte-Carlo code MCNP (version 6.2).^[1] The reactor was modeled in detail based on original technical drawings. The model extends from the reactor core through the reactor pressure vessel (RPV) and its internals into the ex-vessel components including the thermal insulation and the interior and exterior biological shields. A schematic view of the model is shown in Fig. 1. The model was used to calculate the neutron fluence characteristics (spectrum and distribution) of the reactor. In addition, experiments were carried out in two reactors for the validation purpose of the model. For this, the activation of several metal isotopes were calculated at 10 points inside the ex-core instrumentation channel located in the interior biological shield (Fig.1). The results were then compared with the activation measurements of monitors composed of thin metal foils placed in these locations.

EXPERIMENTAL MEASUREMENTS. Several activation monitors connected with a steel cable were placed inside the ex-core instrumentation channel of an active reactor (parallel to the reactor core). Each monitor was composed of an aluminum box filled with thin metal foils: Ti, Fe, Ni, Cu, Zn, In, and Sn. The metal foils have a thickness of 0.1 mm and a size of 5×5 mm to 10×10 mm or 10 mm diameter. The monitors were irradiated in the reactor for about 6 months. After recovery, the neutron activation products were analyzed by gamma spectroscopy using coaxial high purity germanium (HPGe) detector (GMX, 30% efficiency, Ortec-Ametek, Inc.).

CALCULATIONS *VS.* **MEASUREMENTS.** The activation was calculated for each metal foil, and the obtained result was compared with the corresponding measured value. Because the exact position of the monitors inside the channel is not known, the calculations were performed for 3 possible locations of the monitors: close to the outer wall of the channel (C1 - Out), in the middle of the channel (C2 - Mid.), and close to the inner wall of the channel (C3 - In). The activation



Fig. 1: 3D MCNP model of a German PWR.



Fig. 2: Activation calculations vs. measurements for ^{113}In (n,y) $^{114m}In.$

calculations vs. measurements for ¹¹³In (n,χ) ^{114m}In are shown in Fig.2. As can be seen, the results of the calculations agreed very well with the measurement values. The same observation was obtained for all of the examined metal foils activation reaction.

ACKNOWLEDGEMENTS. This work is funded by the Federal Ministry of Education and Research (BMBF) under contract number 15S9409A and supported by PreussenElektra GmbH.

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Coupled thermal hydraulic and neutronic simulations of the ESFR-SMART core

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The in-house reactor dynamics code DYN3D^[1] was used to perform coupled core thermal hydraulic (TH) and neutronic simulations of the European Sodium Fast Reactor core in the framework of the Horizon 2020 ESFR-SMART project, hereinafter referred to as the ESFR-SMART core. The consideration of burn-up dependence of the fuel stoichiometry and porosity was found to be critical for fuel temperature calculations. As compared to stand-alone mode calculations, the effect of the coupling was evaluated by comparing radial distributions of power, average fuel temperature and coolant heat-up. Lastly, the safety relevant global reactivity coefficients were evaluated in coupled simulations.

REFERENCE CORE DESCRIPTION. The ESFR-SMART core is a large 3,600 MWth oxide core containing 216 inner core and 288 outer core fuel subassemblies (S/As), 31 corium discharge tubes, and 24 control and 12 dedicated shutdown devices. The active core is surrounded by three types of reflector layers with a total of 268 S/As. At 293 K, the S/A pitch size is 20,985 cm, and the height of the inner and outer active cores are 75 cm and 95 cm, respectively. In this study, the core was modeled at End of Equilibrium Cycle (EOEC) with nominal power, 668 K core inlet temperature and a total coolant flowrate of 18,705 kg/s. The burn-up distribution was obtained with the Serpent code in six batches per zone,^[2] separately for fissile and fertile axial regions. The full specification is available in the project deliverables.^[3]

NEUTRONICS. In stand-alone neutronic calculations, a uniform fuel temperature of 1,500K and coolant density of 0.837 g/cm³ was assumed. In coupled calculations, the spatial distribution of these two parameters was provided by the TH module of DYN3D. The applied homogenized few-group cross sections, generated with Serpent, were parametrized by fuel temperature and coolant density.

THERMAL HYDRAULICS. In stand-alone TH calculations, the power distribution was provided from the stand-alone neutronic results of DYN3D. In coupled calculations, the power distribution was calculated by the code. According to the recommendation of the project's fuel performance and gap conductance task group,^[4] the following new correlations were added to the TH module for proper simulation of the fuel rod thermal conductance in mixed oxide sodium cooled fast reactor fuels:

- Philipponneau correlation for thermal conductivity,^[5]
- burn-up dependent fuel stoichiometry,
- burn-up dependent fuel porosity,
- burn-up and linear heat rate dependent gap conductance.

FUEL ROD TEMPERATURES. Typically in system analyses, fixed values of fuel stoichiometry (x) and porosity (p) are used, *e.g.*, x = 0.02 and p = 0.045 for fresh fuel. Such simplification of disregarding burn-up (bu) dependence leads to an opposite tendency of fuel conductivity with increasing burn-up (Fig. 1). It was found that applying fresh fuel stoichiometry and porosity in the ESFR-SMART core will result a significantly lower average fuel temperature (by about 80 K) within S/As with a low-burn-up of 23.21 MWd/kg. Furthermore, in the core center, between S/As with low-burn-up and high-burn-up such as 112.42 MWd/kg, the difference in



Fig.1: Difference in fuel thermal conductivity at 1,500 K. Fixed vs. parametrized fuel stoichiometry and porosity.

fuel temperature changes drastically (swings from -10 K to +120 K) when considering burn-up dependence in *x* and *p*.

COUPLED RESULTS. The spatial distributions of power, fuel temperature and coolant heat-up obtained in coupled simulation were compared against the stand-alone mode calculation, as seen in Fig. 2. It can be concluded that taking into account the TH feedbacks with a non-uniform temperature field can deviate the power distribution up to 0.8% (Fig. 2a). Consequently, this will affect the average fuel temperature and coolant heat-up in S/As up to about 7 K and 2 K (Fig. 2b-c), respectively.



Fig. 2: Differences in radial distributions of selected quantities between coupled vs. stand-alone modes. Power (A), fuel temperature (B), and coolant heat-up (C).

Finally, by using the coupled model, the following global reactivity coefficients were obtained at nominal core state:

- k = -0.59 pcm per 1 K change of the inlet temperature,
- -g = -0.30 pcm per 1 K change of the core heat-up,
- h = -3.21 pcm per 1 % change of reactor power.

These negative values indicate a high level of ESFR-SMART core self-protection and safety.

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Validation of a CFD simulation against a ROCOM mixing experiment

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The CFD simulation of a boron dilution scenario is compared with experimental results obtained at the ROCOM test facility. The main features of the scenario are asymmetric, transient mass flow conditions in the affected loops 1 and 2 of a KONVOI-type reactor vessel and reduced density of the underborated coolant slugs fed into the reactor trough the cold legs of both loops.

Boron is a neutron absorber. In pressurized water reactors, it is added as boric acid to the primary loop coolant at varying concentrations over the fuel cycle in order to remove excess reactivity from the reactor core. Under normal conditions with forced circulation, boric acid is homogeneously dissolved in the coolant inventory. In the aftermath of a small break loss of coolant accident (SB LOCA), decay heat is removed from the core by coolant evaporation. Subsequent condensation in the steam generator leads to the formation of underborated liquid coolant that accumulates in the loop seals of the primary circuit. After refilling and restart of natural circulation (NC), slugs of underborated coolant reach the reactor core and give rise to re-criticality and power insertion. Boron dilution was identified as a key issue of PWR safety. Its analysis is part of the design and of the licensing and oversight procedures. Nevertheless, the underlying physical process - single phase natural circulation flow with liquid mixing - is a challenging task for numerical analysis.

EXPERIMENTAL. The ROCOM (ROssendorf COolant Mixing) test facility, formerly located at Helmholtz-Zentrum Dresden-Rossendorf (HZDR), was a 1:5 acrylic glass mockup of the German four-loop pressurized water reactor KONVOI (Siemens KWU) including the primary coolant circuit. The E2.3 experiment serves as the validation base for the CFD study, its measurement data are publicly available.^[1,2] It models asymmetric flow conditions with an ethanol-water mixture representing underborated slugs which arrive at different and time-dependent mass flow rates through vessel inlets of loops 1 and 2; loops 3 and 4 remain unaffected from underboration and mass flow changes. Mixing is measured by means of wire-mesh sensors based on local electrical conductivity.

NUMERICAL STUDY. The mixing process was simulated with the CFD software ANSYS CFX. The computational geometry is shown in Fig. 1. The highlighted sections of loops 1 and 2 hold the entire volume of underborated coolant at the start of transient calculation. Measured and simulated boron distributions in the inlet cross sections indicate that slug and regular coolant are layered one on top of the other (Fig.2). After opening the valves, the initially vertical concentration front is tipping over because of the different densities of the fluids. In the top half of the pipe, the lighter, underborated slug is moving towards the pressure vessel. Coolant mixing in the downcomer region is shown in Fig. 3. Simulated and experimental distributions resemble each other, presenting similar vortex-like flow structures. In the core inlet plane, underboration dominates the peripheral coolant channels of the upper and lower left quadrants, where loops 1 and 2 are located (Fig. 4). A band of underboration across the central part can be observed, creating two zones of channels with little underboration. The shown distributions suggest that



Fig. 2: Experimental and simulated boron distributions at inlet sensor positions at the time of minimum area-averaged concentration.



Fig. 3: Experimental (top) and simulated boron distributions in the unwrapped downcomer.

$t = 65 \mathrm{s}, C_{\mathrm{B,min}} = 1218 \mathrm{ppm}$	$t = 65 \mathrm{s}, C_{\mathrm{B,min}} = 1290 \mathrm{ppm}$	$C_{\rm B}/{\rm ppm}$ - 2200
		-2100
		-2000
	••••••••••••••••••	1 900
		- 1800
	•••••••	1 700
••••••••••••••	•••••••••••••••••	1 600
		- 1 500
		1 400
		1 300
		1 000

Fig. 4: Boron distributions in the core inlet plane.

three-dimensional CFD is able to properly simulate buoyancy-driven fluid mixing in the different parts of a pressure vessel.

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Validation of McSAFE simulation tools

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One of main goals of the EU Horizon 2020 McSAFE project was the development of a high-fidelity multiphysics tool combining Monte Carlo neutron transport, subchannel thermalhydraulics and fuel-performance analysis. The final stage of this work was validations *versus* nuclear power plant measurement data. A full-core pinby-pin depletion of the first operating cycle of a Pre-Konvoi PWR plant was simulated. The critical boron concentration and a set of pin-level neutron flux profiles were compared against measurements with very good agreement.

In the framework of the EU Horizon 2020 McSAFE project, the continuous-energy Monte Carlo code Serpent 2 has been coupled to SUBCHANFLOW (SCF), a subchannel thermal-hydraulics code, and TRANSURANUS (TU), a fuel-performance code (Fig. 1).^[1-4] The purpose of this three-code coupling is to perform highly detailed depletion simulations combining high-fidelity neutronics with pin-by-pin thermal-hydraulic and thermomechanic analysis. This methodology allows for the direct calculation of burnup-dependent local safety parameters at pin level for Light Water Reactor (LWR) designs, which is one of the main objectives of the project.



Fig. 1: Software design.

Each code is implemented as a C++ solver class derived from a common base class, defining a standard coupling interface and masking the internal calculation methods, data structures and programming languages. Feedback data is handled through fields defined in unstructured meshes using the MEDCoupling open-source library, which facilitates the data exchange and the code-to-code interpolations. The actual coupling scheme is implemented in a C++ supervisor program using object-oriented features. This methodology is particularly useful for this three-code coupling, since it increases the maintainability and reusability of the code with respect to traditional approaches such as master-slave coupling.

VALIDATION. The validation of Serpent-SCF-TU pin-level burnup capabilities were validated on data from a German Pre-Konvoi PWR Nuclear Power Plant (NPP). For the beginning of the first cycle of operation, numerical results are compared with measured critical boron concentration and neutron flux profiles. To assess the impact of the fuel solution method, the results obtained performing the thermomechanic analysis in TU, *i. e.* with the three-code coupling, are compared with the ones using Serpent-SCF without TU, solving the fuel behavior in SCF with a simplified approach.



Fig. 2: Critical boron concentration.

RESULTS. The calculated and measured critical boron concentration is shown in Fig. 2. The RMS difference between the results and the experimental data is 3.5 and 4.8 ppm for Serpent-SCF and Serpent-SCF-TU respectively. The deviations with the measurements are well below the acceptability criterion for this type of calculation, which is usually of the order of 50 ppm, and within the typical uncertainties due to specifications and nuclear data. Moreover, the differences between results with and without TU, which are lower than 3 ppm for the whole range, are not significant. These results serve to verify the calculation of global parameters like reactivity, which involve the burnup simulation with equilibrium xenon and the critical boron iteration.

Pin-level neutron flux profiles obtained with Serpent are compared with aeroball measurements and the results are consistent with the experimental data within their statistical uncertainty. These results demonstrate the accuracy of the simulation tool, both in terms of global behavior and of local pin-level parameters.

The fuel solution method has shown minor effect on the neutronics for the amount of burnup considered in this work. The advantage of using TU should be clearer at higher burnup due to phenomena such as the corrosion layer and the behavior of the high-burnup structure, and therefore the current work should be extended to higher depletion.

ACKNOWLEDGEMENTS. This work has been funded from the McSAFE project which is receiving funding from the Euratom research and training program 2014-2018 under grant agreement No 755097.

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Assessment of ATHLET heat transfer models by comparison to steam condensation experiments performed at the COSMEA test rig

M. Jobst

Steam condensation experiments performed at the HZDR COSMEA test rig are modelled with the system code ATHLET and numerical results are compared to experimental data. With the default Dittus-Boelter correlation applied in ATHLET, the code under-predicts the experimental data significantly. Alternative correlations were applied considering the increased heat transfer on secondary side due to swirl elements.

Future nuclear power plant designs, such as the boiling water reactor KERENA, will be based on passive safety system for decay heat removal which are independent from electrical power. For safety assessment of nuclear power plants, system codes such as AC² ATHLET or RELAP are applied. These codes have been qualified and validated for active thermal-hydraulic circuits, but there is still further work necessary to validate them for passive decay heat removal systems. In order to obtain experimental data for such code validation, the COSMEA (COndenSation test rig for flow Morphology and hEAt transfer studies) has been implemented to the TOPFLOW test facility at HZDR.^[1]

EXPERIMENTAL SET-UP. The COSMEA test rig is a slightly inclined tube-in-tube condenser. Its inner condenser tube is made of stainless steel with 43.1 mm inner diameter, 2.6 mm wall thickness and 2,992 mm length.^[1] Saturated steam at specified inlet quality is injected at different flow rates and the condensation mass flow rate is measured. The inner tube is cooled by the coaxially arranged outer tube. Experiments have been performed for different steam pressures from 5 to 65 bar.^[1]

NUMERICAL MODEL. An ATHLET model of COSMEA has been set up, which includes primary and secondary side of the condenser tube, and thus the major heat transfer processes of the experiment are included: steam condensation on the inside of the tube, heat conduction and convective heat transfer on the outside of the tube. Figure 1 shows the nodalization scheme of the model with axial division of the condensation tube into 100 nodes (control volumes). The Dobson-Chato condensation model is applied on the primary side, the convective heat transfer on the secondary side is calculated according to the Dittus-Boelter correlation.^[2]



Fig. 1: Nodalization scheme of the COSMEA ATHLET model.

RESULTS. A total of 28 COSMEA tests from the 2018 measurement campaign^[1] with different boundary conditions were calculated with ATHLET, from which 17 have been selected for discussion in this paper. Figure 2 shows the calculated condensation mass flow rates (solid lines) for different primary pressures and averaged steam volume flows. The comparison to experimental data (dashed lines) reveals, that ATHLET under-predicts the condensation rates (up to 15%). In the experiments carried out in 2018, five swirl elements were attached to the secondary side in order to achieve the



Fig. 2: Condensation mass flow rate depending on steam pressure and averaged steam volume flow in the test section (TS). Heat transfer on secondary side is calculated by Dittus-Boelter correlation.

best possible mixing and a high and uniform heat transfer. By definition, the Dittus-Boelter correlation cannot describe this increase in the heat transfer coefficient. Therefore, the literature was searched for suitable correlations. The swirl elements deflect the flow by about 30°, and this angle decreases with increasing distance from the swirl element. The correlations after Cavazzuti and Corticelli describe both the decrease of the swirl angle and the dependence of the Nusselt number on the swirl angle.^[3] Thus, an average increase of the heat transfer coefficient compared to the Dittus-Boelter correlation of 40 % was determined. This factor can be accounted for the ATHLET input data set and a much better agreement with the measurement results was obtained after recalculation (Fig.3).





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Investigations on the multiscale and multi-physics dynamic behavior of German PWRs during reactivity accidents

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Despite the high degree of maturity reached by system codes in nuclear safety research, the accurate simulation of complex accident sequences, such as reactivity accidents still presents several gaps. The main uncertainties arise from the correct modelling of asymmetries in the reactor coolant system (RCS) together with local thermohydraulic phenomena in the vessel and reactivity feedback effects in the core. The accurate representation of such phenomena requires a combined multi-physics and multiscale approach. This is investigated with ATHLET-DYN3D, exploiting the 3-D Neutron Kinetics of DYN3D together with the 3-D features of ATHLET 3D-Module.

The present project is enclosed in the KEK (*Kompetenzerhalt in der Kerntechnik*) initiative, which fosters the preservation of the competence on safety-related problems in nuclear technology in Germany. The main goal is the development and application of new and improved reactor models for the investigation of the early in-vessel phase of hypothetical core melt scenarios with an eventual asymmetric cold-water injection into the RCS. The development of advanced reactor configurations in ATHLET to enhance the simulation of large 3-D effects in vessel and piping, together with multiphysics calculations by ATHLET-DYN3D fosters a more realistic simulation of complex accident sequences. Investigation on the effectiveness of the implementation of bleed and feed procedures is also foreseen in the project.^[1] The investigation is applied to generic German PWRs of Konvoi type.

VALIDATION OF ATHLET 3D-MODULE. During the project, the 3D-Module of ATHLET has been validated for the simulation of coolant mixing under buoyancy-driven flow conditions based on the experiments E2.3,^[2] T1.1 and T2.1^[3] conducted at the ROCOM test facility. With this aim, an ATHLET model of ROCOM has been developed and additional sensitivity runs have been performed based on the spatial resolution as well as verification of boron transport models. The results from the validation exercises have been successful and have demonstrated the ability of ATHLET to simulate coolant mixing during boron dilution and overcooling conditions.^[4]

ACCIDENT ANALYSES. Two postulated reactivity accident scenarios have been selected in the project, one consisting of a main steam line break (MSLB) with doubleguillotine break in steam generator 1 and a small-break loss of coolant accident (SBLOCA) with inherent boron dilution. The break in the SBLOCA is located in cold leg 1 and accounts for 44 cm². The analysis of the SBLOCA sequence is divided into two different parts, one focusing on the study of inherent boron dilution, whereas second part models a simultaneous loss of AC power meant to investigate the plant dynamic response during the implementation of bleed and feed procedures to prevent core meltdown. An ATHLET model of a Generic PWR of Konvoi type has been developed for the analyses with 1-D and 3-D vessel topologies, meant for comparison of the results between both. **RESULTS AND DISCUSSION.** The results from the MSLB study show that even with conservative boundary conditions, in none of the runs a power surge takes place in the core. The negative reactivity from the control rods exceeds the positive reactivity by the Doppler and moderator density feedback effects. The results with the traditional 1-D vessel model show a temperature difference of about 40 K with respect to the validated 3-D approach, which in turn leads to higher reactivity insertion by the Doppler. Regarding the study of SBLOCA with inherent boron dilution, some important insights are:

- Large accumulation of condensate in U-tubes and loop seals (up to 30,000 kg) under failure of safety systems together with high thermal loads in steam generators.
- Counter-current flow limitation (CCFL) and flooding limit exceeded in hot legs and U-Tubes.
- Very effective mixing of condensate with water from the safety systems in the loop seals and cold legs.
- Increase in the boron concentration in the core along with evaporation. Acts as a safety inherent mechanism.

On the other hand, implementation of bleed and feed procedures show grace times of 1 h and 2 h during primary and secondary side depressurization, respectively. Feeding the secondary and primary sides with an external pump has proven to be a very efficient mechanism to ensure long-term core cooling and prevent core meltdown, being secondary feeding the preferred option.



Fig. 1: Vessel water level during secondary bleed and feed.

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Extension of ATHLET/DYN3D to SFRs: Verification against DYN3D on the Phénix reactor core

V.A.Di Nora

The ATHLET/DYN3D code is currently under extension for being applied to the analysis of entire sodium-cooled fast reactors (SFRs) systems. In this regard, modifications have been implemented on the code-system and then tested. This work provides an overview of the verification tests executed on the Phénix reactor core and the related outcomes obtained.

INTRODUCTION AND MOTIVATION. The DYN3D code is a three-dimensional neutron kinetic (NK) solver equipped with built-in thermal-hydraulic (TH) and fuel rod behavior modules.^[1] The code has been currently extended to the safety analysis of SFRs cores. The extension includes a newly implemented thermal-mechanical (TM) module that accounts for radial diagrid and axial fuel thermal expansions.^[2] The upgraded DYN3D was validated against the IAEA benchmark devoted to the Phénix natural circulation (NC) test.^[3,4] To further broaden the applicability of the code to safety analyses of whole SFR systems, DYN3D should be coupled with a TH code able to model the liquid sodium as a coolant. The TH code ATHLET was chosen for the purpose. A previous coupled ATHLET/DYN3D version was already developed and validated for LWRs.^[5] However, the application of the code system to SFRs has required the update and extension of the existing code coupling routines.

Preliminarily to the application onto whole SFR plants, the modified code version has to be tested. The prediction capabilities of the new ATHLET/DYN3D and DYN3D standalone are expected to be equivalent or comparable. The current research aims to verify such an expectation through the benchmark of the codes conducted at the core level of the Phénix SFR.

TEST DESCRIPTION AND METHODOLOGY. The Phénix reactor was a small-size SFR prototype with a MOX-loaded kernel subdivided into inner and outer core fissile regions, radial blanket, and peripheral reflector. On the occasion of the reactor shut-down, in 2009, several experiments were conducted, including the NC test. For such an experiment, the reactor was stabilized at a power level of 120 MWth. The test was then initiated by manually inducing the steam generators' dry out, originating thus a loss of heat sink scenario. Details on core geometry, material compositions, and system perturbations are available from the benchmark specifications. The verification is conducted by:

- Considering as references the static and transient DYN3D solutions of the NC test, respectively found at time zero and along the first 458 seconds of the test.
- Developing an ATHLET/DYN3D model of the Phénix NC test, equivalent to the existing DYN3D model, and comparing the related calculations results against the references. For the static analysis, the study compares the region-wise temperature distributions of fuel, cladding, and sodium. Evolutions of power, reactivity components, and the average temperatures of fissile and fertile channels are then benchmarked for the transient analysis.

RESULTS. Table 1 summarizes the results of the verification under static conditions. The fuel, clad and sodium temperature distributions are in good agreement with each other with errors below 1.5%. The best agreement is found for the sodium (Na) distribution which maximum error was 0.2%. Figure 1 shows the comparison of transient solutions against

Tab. 1: Phénix static comparison ATHLET/DYN3D against DYN3D: RMS error on temperature distributions.

Material	Inner core	Outer core	Blanket
Fuel	0.59	0.65	0.63
Clad	1.31	1.20	0.61
Na	0.16	0.14	0.08



Fig. 1: Phénix transient comparison evolutions of power, reactivity, and temperatures of core materials.

the reference. Figures 1a) and b) show the evolutions of power and reactivity components, namely, Doppler, Sodium density (Na-Q), clad-driven fuel axial expansion, and radial diagrid expansion. The curves are in good agreement and basically overlap. Figure 1c) shows the temperature evolutions of fuel, clad, and sodium averaged within the inner and outer core regions. Similarly, Fig. 1d) presents the average temperature trends of the blanket materials. Even in this case, the curves are very close to each other.

CONCLUSIONS. The current benchmark activity performed at the Phénix core level demonstrates that either using ATHLET/DYN3D or DYN3D-standalone, practically equivalent results are obtained. Consistent and close temperature distributions were found under static conditions, and excellent agreements were observed on the transient evolutions of power, reactivity, and temperatures of core materials. Such findings lay the foundations for applying ATHLET/DYN3D to the analyses of the entire Phénix primary system.

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New developments of the Monte Carlo code TRAMO and their validations on experiments of Russian VVER-1000 reactor equipment

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The Russian Partner SEC NRS (Scientific and Engineering Centre for Nuclear Radiation Safety) is very much interested in using TRAMO as an independent computer code for validation and verification of activity and shielding studies at Russian nuclear power plants. For this purpose, TRAMO and its auxiliary programs have been further developed in such a way that the Russian partner is now in principle able to perform independent calculations. This paper gives a brief overview of these development steps.

TRAMO^[1] is a Monte Carlo code for three-dimensional geometries developed at the Helmholtz-Zentrum Dresden-Rossendorf and was optimized especially for structures of a nuclear power plant. It can be used to calculate neutron and gamma fluences in the reactor pressure vessel and its internals. In principle, however, other geometries are also conceivable. The possibility of obtaining results with sufficient statistical accuracy in a reasonable time is ensured by various optimization methods. These include variance reducing methods such as weight window, Russian roulette and implicit capture as well as source biasing. The heart of the optimization, however, is the auxiliary program TRAWEI,[1] which generates spatial and energy-dependent weight factors precisely matched to the geometry. For a calculation with TRAMO, additional input variables are necessary. In detail, these are the source distribution, which is supplemented by a pin distribution if necessary, and various crosssection data based on common nuclear databases such as ENDF-B/VII or JENDL-4.0.[2,3]

Within the framework of the scientific-technical cooperation between Germany and Russia in the field of the peaceful use of nuclear energy on the basis of an agreement of 1987, there is a long-standing and intensive cooperation with the Russian partner, which is reflected in many joint projects. In the current project, TRAMO has to be further developed and improved in such a way that it is possible for the Russian partner to carry out fluence calculations on its own, since Russian nuclear safety regulations require calculations with an independent code.^[4,5] For this purpose, a number of necessary steps have been agreed upon.

First of all, the TRAMO code itself was transferred into a modern Fortran and contemporary programming structures, so that on the one hand the readability was increased and on the other hand the possibility of further model development (also by externals) was given at all. Furthermore, the Russian partner sent us all material compositions of relevant VVER-440 and VVER-1000 reactors, so that an extensive database with cross section data and scattering matrices in two energy group structures was generated. A detailed geometry model was provided for one reference case each of a VVER-440 (Kola-3) and VVER-1000 (Rovno-3). Various material samples (monitors) are located at the outside surface of the reactor pressure vessel, which can be used to experimentally determine the neutron fluence within a cycle. On their basis, we developed a generalized TRAMO input model with special focus on the regions of the monitors. For each horizontal and vertical set of monitors, TRAWEI was used to generate corresponding weight data sets, and TRAMO was

used to generate source biasing data by activating and deactivating specific source regions.

To make it easier for the Russian partner to handle the code and the numerous helper programs, a shell script was developed. It regulates the whole process of a TRAMO calculation including input and output steps and includes in detail the selection of the reactor type, the orientation of the monitors, the positions of them (height and angle), the energy group structure and a selection of a 60° sector of the VVER reactor. In addition, global settings such as reactor power can still be specified in a simplified input file. The script reads and generates an appropriate base input file, selects the correct weight data sets, source biasing data and effective cross sections, and converts the source and pin data distribution to TRAMO format. Then, the actual calculation is performed and a detailed and a reduced output file are written. With the help of a data evaluation program, the fluences at the monitor positions are interpolated if necessary and converted into measurable quantities such as reaction rates. The entire software package, including the necessary data and a manual for the use and explanation of the input variables, was handed over to the Russian partner for independent use.

To verify the correct implementation of the new input models, the existing models of the reference cases mentioned above were adjusted and recalculated. The fluence values were found to be largely in agreement. Thereupon, the VVER-1000 model was adapted to the Kalinin-4 nuclear power plant. The active zone of the core is slightly larger than at Rovno, otherwise the design is almost identical. One horizontal and one vertical arrangement of monitors each was entered into the model according to the specifications of the Russian partner, and corresponding calculations were performed. The fluences agree well with those of Rovno-3. The deviations amount to a maximum of 10 %. The comparison with the experimentally obtained data will be performed by the Russian partner and is expected soon.

An outline for a follow-up project has been submitted and has successfully passed the preliminary review. Within this, the TRAMO package is to be turned into a holistic tool for reactor dosimetry and activity calculations. Besides the model development by introducing elliptical bodies, TRAWEI shall be integrated directly into TRAMO, so that the user can determine whether weight data sets shall be calculated or whether they are already available, the possibility for parallelization (and thus for a considerable increase in computing speed) shall be created and a graphical input tool shall facilitate the handling by users.

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Development of a generally usable weight window generator using the TRAWEI code

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Nuclear decommissioning is currently one of the most important activities of the nuclear industry. The optimal planning of the decommissioning requires good estimations of the neutron activation distribution in the structural materials of the nuclear power plant. These estimations are associated with deep penetration problems, which can be solved using Monte-Carlo (MC) codes such as the MCNP. In Monte-Carlo codes, a large number of simulated particles with corresponding long computational times are usually needed to obtain statistically reliable results in volumes far away from the source. Therefore, variance reduction techniques are required to achieve highly accurate results in a reasonable computation time. Weight window is one of the most widely used variance reduction techniques for MC deep penetration problem calculations. However, in complicated shielding calculations, the estimation of optimal weight window parameters is quite challenging. The present work is aimed to develop an optimal weight window generator program based on Recursive Monte-Carlo (RMC) method, which is to be tested in neutron fluence calculations in a German Pressurized Water Reactor (PWR) decommissioning study.

The MCNP code has a build-in weight-window parameters generator.^[1] In complicated shielding calculations the generator outputs zero weight windows. Several approaches have been suggested and applied to solve this problem and achieve more meaningful (non-zero) weight window parameters, such as the MAGIC method,^[2] Density method, and Distributed method. However, those approaches are not straightforward and require iterative calculations. Therefore, the present work aims to develop an optimal and direct zone-weights parameters generator based on Recursive Monte-Carlo (RMC) method. The development is based on the TRAWEI code.^[3] The developed generator will output space- and energy-dependent mesh-based zone weights, which can directly be used with the MCNP code.

THE RECURSIVE MONTE CARLO METHOD. The RMC method is specifically applied for deep-penetration problems, such as in the case of decommissioning of nuclear power plants.^[4] The method is used to calculate the importance functions of a focus region, *i.e.*, the probabilities that a neutron from one cell will reach the result region, using already generated results from the neighborhood cells and lower energy groups. At this point, the influence on the target zone is known, and the value calculated up to this point is weighted accordingly. Based on the results already obtained, the influence functions of the next cells are successively determined. The Monte-Carlo play of a particle can be interrupted when it crosses a cell with known influence. From this point on, the influence on the target zone is known and can be used for the calculation. By skillfully sequencing the cell calculations, the calculation can be accelerated. The same procedure is used for the down scattering of neutrons, starting with the lowest energy group and continuing with the next higher one. The already existing results can then be used for down scattering. The weights are estimated by taking the reciprocal values of the importance function. This procedure ensures that for each volume element (cell), a weight value is determined and no iteration is necessary.

THE TRAWEI CODE. The TRAWEI code was developed at Helmholtz-Zentrum Dresden-Rossendorf (HZDR) to calculate cell weights, which can be used in the TRAMO code. The TRAMO and TRAWEI codes are 3D multi-group Monte-Carlo codes applicable for typical reactor geometries.^[3] The codes were specially developed for solving particle transport problems in the field of reactor dosimetry and shielding calculations. They use various nuclear data files (ENDF/B-VII, JEFF3.3, own data supply). TRAWEI is an RMC code that is responsible for generating the weights. The weights can be used directly in TRAMO. Currently, an external weight transfer program is developed to transfer the TRAWEI weights also for MCNP calculations. A schematic view of the development structure is shown in Fig. 1.



Fig. 1: Utilization of TRAWEI weights for Monte Carlo calculations.

ONGOING DEVELOPMENT. The developed weight window generator program will be used together with the MCNP code for German PWRs decommissioning studies. Several realistic geometries have been prepared to validate the accuracy of the developed program. The preparation of the TRAWEI code has been completed and the development of the transfer programs has been initiated. First test calculations have been carried out. They will be evaluated by comparison with existing weight window generators, *e.g.*, MCNP's weight-window generators.

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PUBLICATIONS

ARTICLES (PEER-REVIEWED) EDITORIAL WORK FURTHER CONTRIBUTIONS PATENT ORAL PRESENTATIONS THESES

• ARTICLES (PEER-REVIEWED)



Statistics on the publication activity of the Institute of Resource Ecology (IRE) within the last 10 years. The diagram considers peer-reviewed articles that were published under primary responsibility of the institute. Note that the given impact factors based on data from 2019. The annual update might change the statistics retrospectively.

The following compilation provides 113 peer-reviewed articles published in 2020 representing any contribution of members of the Institute of Resource Ecology.

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(Technische Universität Dresden, 2020, except where noted)

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Eibl, M.

Actinide Interaction with Zr-bearing Phases: Spectroscopic Investigations of An³⁺ Sorption and Incorporation Reactions with Zirconia

Fichter, S.

Synthesis and Characterization of Tri- and Tetravalent Actinide Amidinates

Iric, K.

Assembly of DNA-encapsulated lipid bilayers and their application to studies of GPCRs

Lehmann, S.

Uran unter reduzierenden Bedingungen – Spektroskopische und thermodynamische Charakterisierung der wässrigen Speziation

MASTER THESES

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Philipp, T.

U(VI) retention by Ca-bentonite and clay minerals at (hyper)alkaline conditions

Wolter, J.-M.

Stability of U(VI) and Cm(III) doped calcium (alumino)silicate hydrate phases at high ionic strengths

Europium Adsorption on the (001) Muscovite Face: A Kinetic Monte Carlo Study

DIPLOMA THESIS

Füssel, T.

Die Wirkung von sorbierten redoxaktiven Spezies auf Aluminiumoxid-Nanopartikeln bei der Immobilisierung von Technetium

BACHELOR THESIS

Gerlach, L.-M.

Mikrokalorimetrische Messungen zur kinetischen Analyse einer statischen Kultur von Lactococcus lactis

Lückel, B.

Untersuchung der Endozytose als Aufnahmeprozess von Uran(VI) in Tabak BY-2-Zellen

Neubert, T.

Untersuchung des mikrobiellen Einflusses auf Kupfer- und GGG40-Grenzflächen im Bentonit unter halophilen Bedingungen

SCIENTIFIC ACTIVITIES

O WORKSHOP

O SEMINARS

O TEACHING ACTIVITIES

INTERNATIONAL WORKSHOP ON "HOW TO INTEGRATE GEOCHEMISTRY AT AFFORDABLE COSTS INTO REACTIVE TRANSPORT FOR LARGE-SCALE SYSTEMS"

HZDR, Germany, February 05-07, 2020

Experts from fields of reactive transport modelling (RTM), statistics, computational numerics, thermodynamics, radiochemistry or geosciences were brought together to discuss RTM on large scales, namely time (up to one million of years) and space (dozens of kilometers). Full-fledged incorporation of basic, geochemical processes, *e. g.* sorption, precipitation, or redox reactions, will thus demand inacceptable long computing times.

As an effective way to integrate geochemistry at affordable costs into RTM, different numerical concepts (*e. g.* multidimensional look-up tables, surrogate functions, machine learning, utilization of uncertainty and sensitivity analysis etc.) already exist and were extensively discussed throughout the workshop.

In total, 40 scientists from 28 research institutes and 8 countries followed the invitation of the organizers, Dr. Madlen Stockmann and Prof. Vinzenz Brendler, HZDR, and presented their research in keynote and regular lectures and during poster session. A tour on the radiochemical laboratories of the IRE had been offered.

Based on the intensive discussions and very positive feedback on the workshop, a continuation is intended to bundle and strengthen the respective research activities and stipulate the international network that started to form during the conference days.

The main topics of the workshop were:

- Introduction and discussion of existing geochemical concepts in reactive transport modelling to describe sorption and related retardation processes of contaminants on a variety of sediments and rocks.
- Focusing on large-scale natural systems as experienced, e.g. in nuclear waste disposal, carbon capture & storage, environmental remediation, or geothermal applications.
- Integration of the discussed approaches into current paradigms in Thermo-Hydraulic-Mechanic-Chemical (THMC) models and long-term safety assessments at affordable costs.
- Promotion of the exchange of scientific knowledge and practical experience between the workshop participants.



List of keynote speakers:

Allan Leal

ETH Zurich, Switzerland

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For more information please visit: <u>www.hzdr.de/smartkd</u>

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O SEMINARS (TALKS OF VISITORS)

Majzlan, Jurai

Friedrich-Schiller-Universität Jena, Germany **What controls the solubility of arsenic and antimony at old mining sites?** *February* 4, 2020

Qin, Guam

Department of Petroleum Engineering, University of Houston, U.S.A. Multi-scale and multi-physics numerial simulation of fluid transport processes in porous media March 23, 2020 Albrecht-Schmitt, Thomas

Department of Chemistry and Biochemistry, Florida State University, Tallahassee, U.S.A. Chemistry Beyond Plutonium: How Relativity Alters Electronic Structure in Heavy Elements June 30, 2020

Fränzle, Stefan

Professur für Umweltbiotechnologie, TU Dresden, Germany

Chritin: Going a long way from Nuclear Fuel Recycling, towards Mining, Applications in Fuel Cells and environmental Trace Element Monitoring – and maybe back December 3, 2020



(Winter term: WT; summer term: ST)

Lectures

Bok, F. Friedrich-Schiller-Universität, Jena ST 2020 Geochemische Modellierung mit Geochemist's Workbench[®]

Brendler, V. Dresden University of Applied Sciences, ST 2020 Radiochemistry

Fahmy, K. Technische Universität Dresden, WT 2019/2020 Molecular Biophysic: Concepts and Methods ST 2020 Biological thermodynamics

Fischer, C. Universität Bremen, ST 2020 Diagenese

> Universität Leipzig WT 2020/2021 Moderne Methoden in der Materialwissenschaft

Heine, T.

Technische Universität Dresden, WT 2019/2020 Quantenchemie WT 2019/2020, WT 2020/2021 PC II (Theorie der Chemischen Bindung) PC III (Einführung in die Computerchemie) WT 2020/2021 Theoretische Chemie Computational Spectroscopy Quantum Chemistry

Yonsei University Seoul, Korea ST 2020 Artificial Intelligence and Machine Learning

Huittinen, N. Technische Universität Dresden, ST 2020 Radioecology

Kuc, A. B. Universität Leipzig, WT 2019/2020, WT 2020/2021 Computational Chemistry of Solids

Lippold, H. *Universität Leipzig, ST* 2020 **Radioanalytik Entstehung und Eigenschaften ionisierender Strahlung** *WT* 2019/2020, *WT* 2020/2021 **Radiochemische Analysenmethoden** Raff, J.

Dresden University of Applied Sciences, WT2019/2020, WT 2020/2021 Mikrobiologie

Technische Universität Dresden, WT 2019/2020, WT 2020/2021 **Mikrobielle Laugung**

Schmidt, M.

Technische Universität Dresden, ST 2020 Chemistry of the f-elements

Courses

The laboratory course "Radiochemistry" was provided from August 24–September 04, 2020, as a part of a module of the chemistry master degree program at the Technische Universität Dresden.

Advisers:

Acker, M.	Kretzschmar, J.
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Kloditz, R.	Weiss, S.
Köhler, L.	Yassin, G.

The IRE provided the experiment "Alpha spectrometric isotope dilution analysis of uranium" of the laboratory course "Instrumental Analysis" held by the Institute for Analytical Chemistry, Technische Universität Dresden.

Advisers:

WT 2019/2020

Hilpmann, S Köhler, L Neumann, J. Roode-Gutzmer, Q. I. Weiss, S. Schymura, S. Universität Leipzig WT 2019/2020, WT 2020/2021 Radioanalytik

Stumpf, T.

Technische Universität Dresden, WT 2019/2020, WT 2020/2021 Radiochemistry

Biophysics course of the Dresden-International-Graduate School.

Advisers:

WT 2019/2020 & WT 2020/2021				
Fahmy, K.	Oertel, J.	Philipp, J.		

Courses and seminars were held referring to the following lectures:

Fahmy, K.

Technische Universität Dresden, WT 2019/2020, WT 2020/2021 Biophysics course of the Dresden-International-Graduate School

Faria Oliveira, A.

Technische Universität Dresden, WT 2020/2021 **Spektroskopie mit dem Computer**

Kuc, A.B.

Universität Leipzig, WT 2020/2021 Spektroskopie (PC) WT 2020/2021 Computational Chemistry for Solids

Springer, M.

Technische Universität Dresden, WT 2020/2021 **Theorie der chemischen Bindung**
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FUNDING ORGANIZATION / COMPANY	PROJECT TITLE	CONTRACT NO. (if applicable)
Commission of the European Communities	CONCERT – European Joint Programme for the Integration of Radiation Protection Research	H2020-662287
(EU)	ENTENTE – European Database for Multiscale Modelling of Radiation Damage	H2020-900018
	ESFR-SMART – European SFR – Safety Measures Assessment and Research Tools	H2020-754501
	EURAD – European Joint Programme on Radioactive Waste Management CORI, DONUT, FUTURE T2 + T3, SFC, UMAN, KMSoK	H2020-847593
	FRACTESUS – Fracture mechanics testing of irradiated RPV steels by means of sub-sized specimens	H2020-900014
	INSIDER – Improved Nuclear Slte characterization for waste minimization in DD operations under constrained EnviRonment	H2020-755554
	M4F – Multiscale Modeling for Fusion and Fission Materials	H2020-755039
	McSAFE – High-Performance Monte Carlo Methods for Safety Demonstration– From Proof of Concept to realistic Safety Analysis and Industry-like Applications	H2020-755097
	McSAFER – High-Performance Advanced Methods and Experimental Investigations for the Safety Evaluation of Generic Small Modular Reactors	H2020-945063
	R2CA: Reduction of Radiological Consequences of Design Basis and Design Extendion Accidents	H2020-847656
	RadoNorm – Towards effective radiation protection based on improved scientific evidence and social considerations – focus on radon and NORM + RadoNorm Open Call Grant	H2020-900009
	STRUMAT-LTO – STRUctural MATerials research for safe Long Term Operation of LWR NPPs	H2020-945272
	SurfBio – Innovation hub for surface and colloid biology research	H2020-952379
	TOP – ERC Starting Grant - Towards the Bottom of the Periodic Table	H2020-759696
Bundesministerium für Wirtschaft und Energie (BMWi) & Bundesmi- nisterium für Bildung und Forschung (BMBF)	EMPRADO – Entwicklung einer Methode zur Pre-Aktivitäts- und Dosisleistungsberechnung von reaktornahen Bauteilen auf Basis von Neutronenfluenzverteilungen Berechnung der Neutronenfluenzverteilung in reaktornahen Bauteilen und deren Validierung an Experimenten als Basis der Aktivitätsrechnungen	15 <i>S</i> 9409A
	f-Char – Spektroskopische Charakterisierung von f-Element-Komplexen mit soft donor-Liganden	02NUK059B
	FENABIUM – Struktur-Wirkungsbeziehungen zwischen f-Elementen und organischen Ligandsystemen mit Naturstoff-basierten Bindungsfunktionen in Hinblick auf eine mögliche Mobilisierung in der Umwelt	02NUK046B
	GRaZ – Geochemische Radionuklidrückhaltung an Zementalterationsphasen	02E11415B
	GRaZ II – Geochemische Radionuklidrückhaltung an Zementalterationsphasen	02E11860B
	iCross – Integrität von Endlagersystemen für radioaktive Abfälle – Skalenübergreifendes Systemverständnis	02NUK053B

FUNDING ORGANIZATION / COMPANY	PROJECT TITLE	CONTRACT NO. (if applicable)
	KRIMI – Kinetik der Radionuklidimmobilisierung	02NUK056C
	MgO-S3 – Spritzbeton für Streckenverschlüsse für HAW-Endlager im Steinsalz	02E11769B
	PANAS – Untersuchungen zu passiven Nachzerfallswärme- Abfuhrsystemen; Teilprojekt B: Untersuchungen zu Kondensationsprozessen im Notkondensator und numerische Simulation einer passiven Wärmeabfuhrkette	02NUK041B
	RADEKOR – Speziation und Transfer von Radionukliden im Menschen unter besonderer Berücksichtigung von Dekorporationsmitteln	02NUK057A
	ResKin_Move – Geschwindigkeitsfeld-Analyse mit Positronen- Emmissions-Tomografie (PET): Parametrisierung und Validierung von Transportmodellen (Plug-Skala)	03G0900A
	SMILE – Smart-K $_{\rm d}$ in der Langzeitsicherheitsanalyse – Anwendungen	02E11668B
	STROEFUN III – Strömungstechnischer Funktionsnachweis für Verschlussbauwerke und flüssigkeitsgestützte Abdichtung des Kontaktbereiches Phase III: Vertiefung Kenntnisstand Kontaktbereich & Injektionsmittel, in situ-Versuche	02E11748B
	TRANS-LARA – Transport- und Transferverhalten langlebiger Radionuklide entlang der kausalen Kette Grundwasser-Boden- Oberfläche-Pflanze unter Berücksichtigung langfristiger klimatischer Veränderungen	02NUK051B
	TransForm – Transnationale Forschungsinfrastruktur und multidisziplinäre akadamische Nachwuchsförderung	01DT20004
	VESPA II – Verhalten langlebiger Spalt- und Aktivierungsprodukte im Nahfeld eines Endlagers und Möglichkeiten ihrer Rückhaltung	02E11607B
	WERREBA – Wege zum effizienten Rückbau von Reaktorkomponenten und Betonabschirmung: Berechnung des Aktivitätsinventars und deren Validierung an Bohrkernen sowie Mobilitätsuntersuchungen von Radionukliden	1559412
DAAD	DAAD PPP Frankreich 2020–2021 – Programm des projektbezogenen Personenaustausches mit Frankreich 2020–2021	57512305
DFG	DFG-Programmpauschalen ab 2014 – Verwaltung von 25 % der ab 2014	
	zu den DFG-Projekten bewilligten Programpauschalen	
	DNA-Struktur Molekulare Mechanismen der Interaktion chaotroper Salze mit natürlichen und künstlichen DNA-Strukturen	FA 248/8-1
	SFB 1415 – Sonderforschungsbereich 1415 "Chemie der synthetischen zweidimensionalen Materialien"	SFB 1415/1 2020
EnBW Kernkraft GmbH	NanoExperte – Gutachten Nachweis von Nanomaterialien	P97417
Helmholtz-Gemeinschaft Deutscher Forschungs-	CROSSING – Crossing borders and scales – an interdisciplinary approach	PIE-0007
zentren e.V. (HGF)	Extract	
	Exzellenznetzwerk-Phase 3 "Physics of Life – The Dynamic Organization of Living Matter"	ExNet-0029-Phase2-3
	NUSAFE / iCross – Integrität von Endlagersystemen für radioaktive Abfälle – Skalenübergreifendes Systemverständnis	SO-093
HTW	Altstahl – Bruchmechanische Prüfung im Rahmen eines Forschungsprojektes der HTW	

FUNDING ORGANIZATION / COMPANY	PROJECT TITLE	CONTRACT NO. (if applicable)
PT Jülich	Ultrasens (WIPANO) – Ultrasensiver spektraler Sensor zur indirekten Messung ionisierender Strahlung	03THWSN004
TÜV NORD	Prüfung geologischer Fragen	M.ASS.06.015.01.F70
TÜV SÜD	Gemischte Reaktorkerne: Sicherheitstechnische Analysen für gemischte Reaktorkerne sowie zur Rekritikalität bei einem schweren Kernschmelzunfall für WWER-1000-Reaktoren (AP12)	3617R01520
	TÜV VVER 1200 code 1 – Support and consulting services	
	WWER AP03	500603389
UJV Rez, a.s.	DYN3D für UJV Wartung 2019/2020	

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