

# Large Scale Shell Model Calculations with a New Algorithm

**D. Bianco**

**Naples**

**F. Andreozzi**

**N. Lo Iudice**

**A. Porrino**

**F. Knapp**

**Dresden 2010**

# Currently adopted methods

- Direct Diagonalization: **Lanczos:Antoine** (E. Caurier et al. Rev. Mod. Phys. 77, 427 (2005) for short review)
- Stochastic methods: **Monte Carlo** (C.W. Johnson et al. PRL 92), suitable for **ground state**. *Minus sign* problem.
- Truncation methods:
  - **Importance Sampling: Quantum MC** (M. Honma et al. PRL 95) selects the relevant basis states.
  - **Density Matrix Renormalization Group (DMRG)** (J. Dukelsky and S. Pittel, Rep. Prog. Phys. 67, 513 (2004))

# Iterative diagonalization algorithm

A. Andreozzi, A. Porrino, and N. Lo Iudice J. Phys. A 02

• Let

$$I = \sum |i\rangle\langle i|$$

$$A (= A^\dagger) \equiv \{ a_{ij} \} = \{ \langle i | \hat{A} | j \rangle \}$$

In our case  $A=H$

$$A \equiv \begin{pmatrix} a_{11} & a_{12} & a_{13} & a_{14} & \dots & a_{1N} \\ a_{21} & a_{22} & a_{23} & a_{24} & \dots & a_{2N} \\ a_{31} & a_{32} & a_{33} & a_{34} & \dots & a_{3N} \\ a_{41} & a_{42} & a_{43} & a_{44} & \dots & a_{4N} \\ \dots & \dots & \dots & \dots & \dots & \dots \\ a_{N1} & \dots & \dots & \dots & \dots & a_{NN} \end{pmatrix}$$

**Goal:**

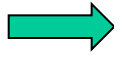
Determine the **lowest eigenvalue** and **eigenvector**

• *1° iteration loop*

$$A_0 = \begin{pmatrix} a_{11} & a_{12} & a_{13} & \dots & a_{1n_0} \\ a_{21} & a_{22} & a_{23} & \dots & a_{2n_0} \\ a_{31} & a_{32} & a_{33} & \dots & a_{3n_0} \\ \dots & \dots & \dots & \dots & \dots \\ a_{n_0 1} & \dots & \dots & \dots & a_{n_0 n_0} \end{pmatrix}$$

$$n_0 \ll N$$

$$v < n_0$$



$$\Lambda_0 = \begin{pmatrix} \lambda_1 & 0 & \dots & 0 \\ 0 & \lambda_2 & \dots & 0 \\ \dots & \dots & \dots & \dots \\ 0 & \dots & \dots & \lambda_v \end{pmatrix}$$

$$\begin{pmatrix} \Lambda_0 & B_{01} \\ B_{10} & A_1 \end{pmatrix}$$

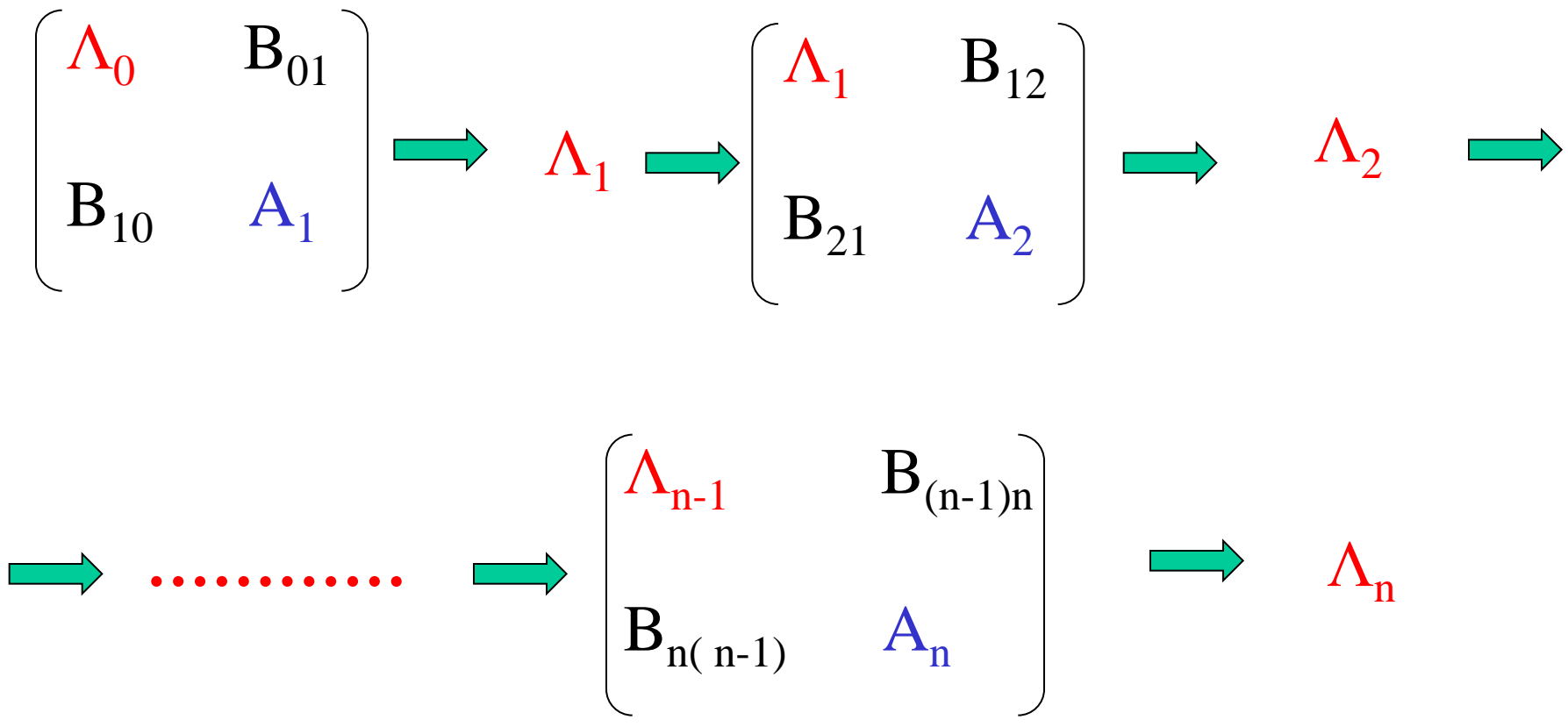
$$A_1 = \begin{pmatrix} a_{(n_0+1)(n_0+1)} & \dots & a_{(n_0+1)n_1} \\ \dots & \dots & \dots \\ \dots & \dots & \dots \\ a_{n_1(n_0+1)} & \dots & a_{n_1 n_1} \end{pmatrix}$$

$$b_{ij} = \langle j | A | \phi_i \rangle$$

$$i = 1, v$$

$$j = n_0 + 1, n_1$$

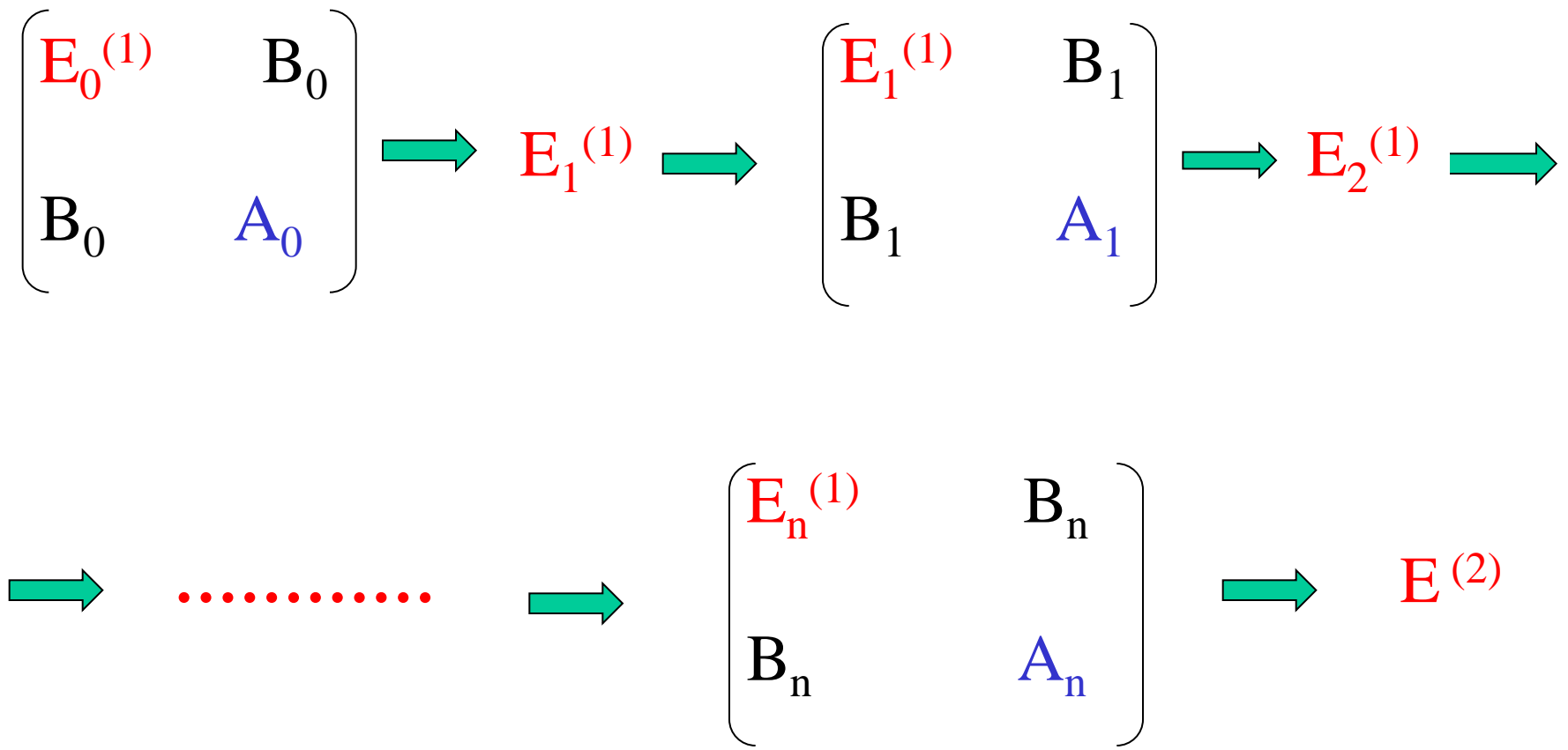
- 1° iteration loop*



$$\Lambda_n \equiv E^{(1)}, \quad |\Phi_N\rangle = |\Psi^{(1)}\rangle = \sum_{i=1, N} c_i^{(1)} |i\rangle$$

**End first iteration loop**

- *2° iteration loop*



$$E^{(2)}, \quad |\Psi^{(2)}\rangle = \sum_{i=1, N} c_i^{(2)} |i\rangle$$

**End 2° iteration loop**

# Iteration loops

$E^{(1)}, \Psi^{(1)}$

$E^{(2)}, \Psi^{(2)}$  .....

$E^{(i)}, \Psi^{(i)}$  .....

## THEOREM

If the sequence  $E^{(i)}$  converges, then

$$E^{(i)} \Psi^{(i)} \longrightarrow E \Psi = H \Psi$$

# Features of the algorithm

- Easy implementation
- Variational foundation
- Robust
  - Convergence to the extremal eigenvalues
    - Numerically stable and **ghost-free** solutions
    - Orthogonality of the computed eigenvectors
- Fast :  $O(N^2)$  operations
- But not enough!



# M-scheme

- **Virtue** : Construction of H matrix very easy.
- **Shortcoming**: the basis dimensions become huge.
- **Remedy**: H is sparse in the m scheme

# M-scheme Implementation

## i. Preliminaries

❖ Sort  $\{j_1^{n_1} \dots j_m^{n_m}\}$  ( $\equiv$  partitions)  $n_1 + \dots + n_i + \dots + n_m = v$   
(according to **increasing energies**)

❖ Choose  $D$

$$\mathbf{I}_d = \sum^d |i\rangle\langle i| = \sum^d |j_1^{n_1} \dots j_m^{n_m}\rangle \langle j_1^{n_1} \dots j_m^{n_m}|$$

### Property

$$J_k |j_1^{n_1} \dots j_m^{n_m}\rangle \in D$$

❖ New Hamiltonian

$$\mathbf{H}_J = \mathbf{H} + \mathbf{c} [\hat{\mathbf{J}}^2 - J(J+1)]$$

## ii. Space decomposition

$$\mathbf{I} = \mathbf{M}_0 \oplus \mathbf{M}_1 \oplus \dots \oplus \mathbf{M}_p$$

### Property

$$|i\rangle_i \in \mathbf{M}_i \longrightarrow \mathbf{H} |i\rangle_i \in \mathbf{M}_{i-1} \oplus \mathbf{M}_i \oplus \mathbf{M}_{i+1}$$

## iii. Starting iteration

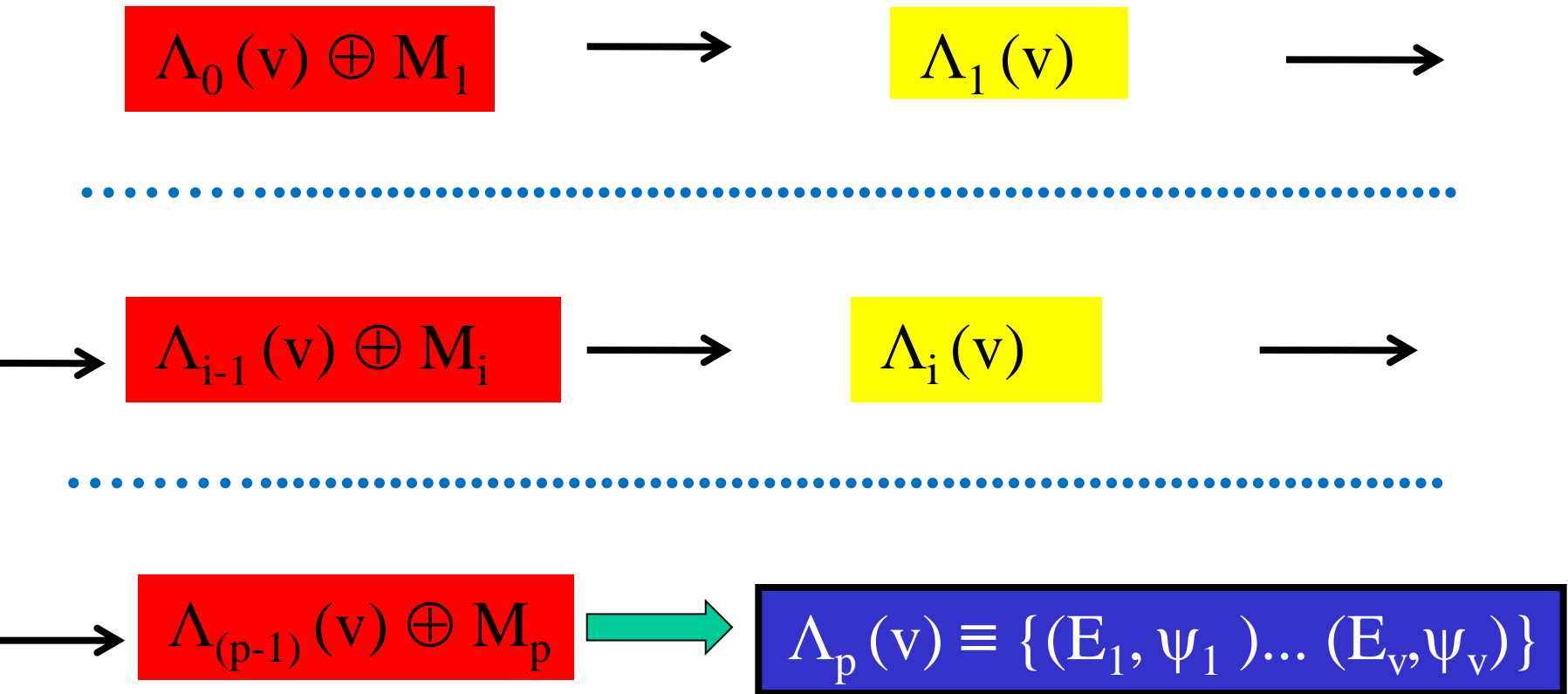
$$\mathbf{M}_0 \longrightarrow \Lambda_0(v) \equiv \{ (E_1^{(0)} \psi_1^{(0)}) \dots (E_v^{(0)}, \psi_v^{(0)}) \}$$

$$\mathbf{H} \psi_k^{(0)} \in \mathbf{M}_0 \oplus \mathbf{M}_1$$



$$\Lambda_0(v) \oplus \mathbf{M}_1 \longrightarrow \Lambda_1(v) \equiv \{ (E_1^{(1)} \psi_1^{(1)}) \dots (E_v^{(1)}, \psi_v^{(1)}) \}$$

### iii. Iterative process



**Analogy:** real space (Wilson) e density matrix (White) renormalization group

# Importance Sampling

$$\varepsilon_1 > \varepsilon_2 > \dots > \varepsilon_{p-1} > \varepsilon_p$$

$M_0$

$$\Lambda_0(v) \equiv \{(E_1^{(0)}, \psi_1^{(0)}) \dots (E_v^{(0)}, \psi_v^{(0)})\}$$

1.

$$|\langle \Psi_k^{(0)} | H | j \rangle|^2 / (a_{jj} - E_k^{(0)}) > \varepsilon_1$$

$$|j\rangle \subset M_1^{(\varepsilon)}$$

$$\Lambda_0(v) \oplus M_1^{(\varepsilon)}$$

$$\Lambda_1^{(\varepsilon)}(v)$$

2.

$$|\langle \Psi_k^{(1)} | H | j \rangle|^2 / (a_{jj} - E_k^{(1)}) > \varepsilon_2$$

$$|j\rangle \subset M_2^{(\varepsilon)} \subset M_1 \oplus M_2$$

# Importance Sampling: Iterative process

$$M_0 \longrightarrow \Lambda_0(v)$$

$$\Lambda_0(v) \oplus M_1^{(\varepsilon)} \longrightarrow \Lambda_1^{(\varepsilon)}(v) \longrightarrow$$

---

$$\longrightarrow \Lambda_{i-1}^{(\varepsilon)}(v) \oplus M_i^{(\varepsilon)} \longrightarrow \Lambda_i^{(\varepsilon)}(v)$$

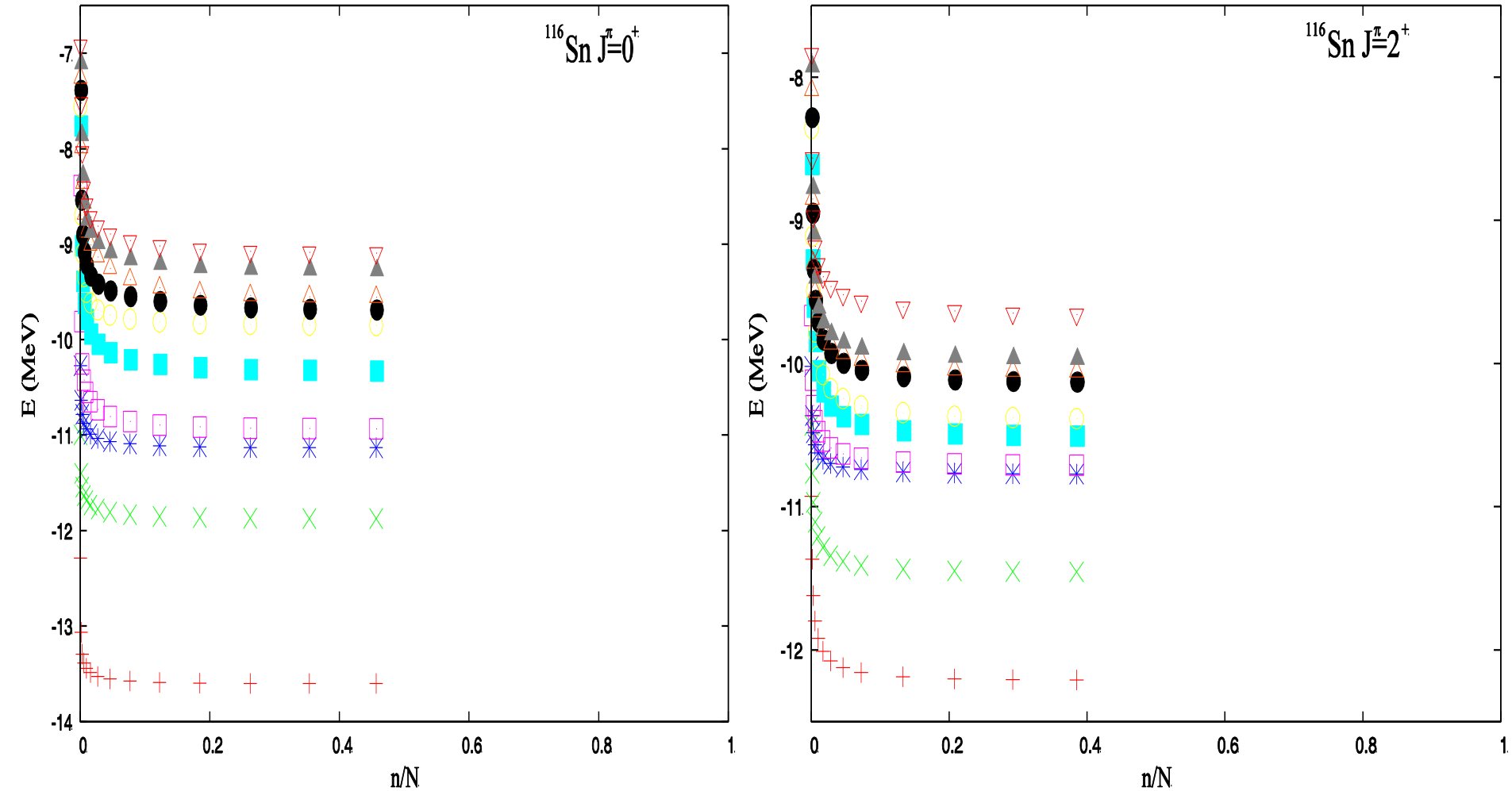
---

$$\longrightarrow \Lambda_{p-1}^{(\varepsilon)}(v) \oplus M_p^{(\varepsilon)}$$

↳

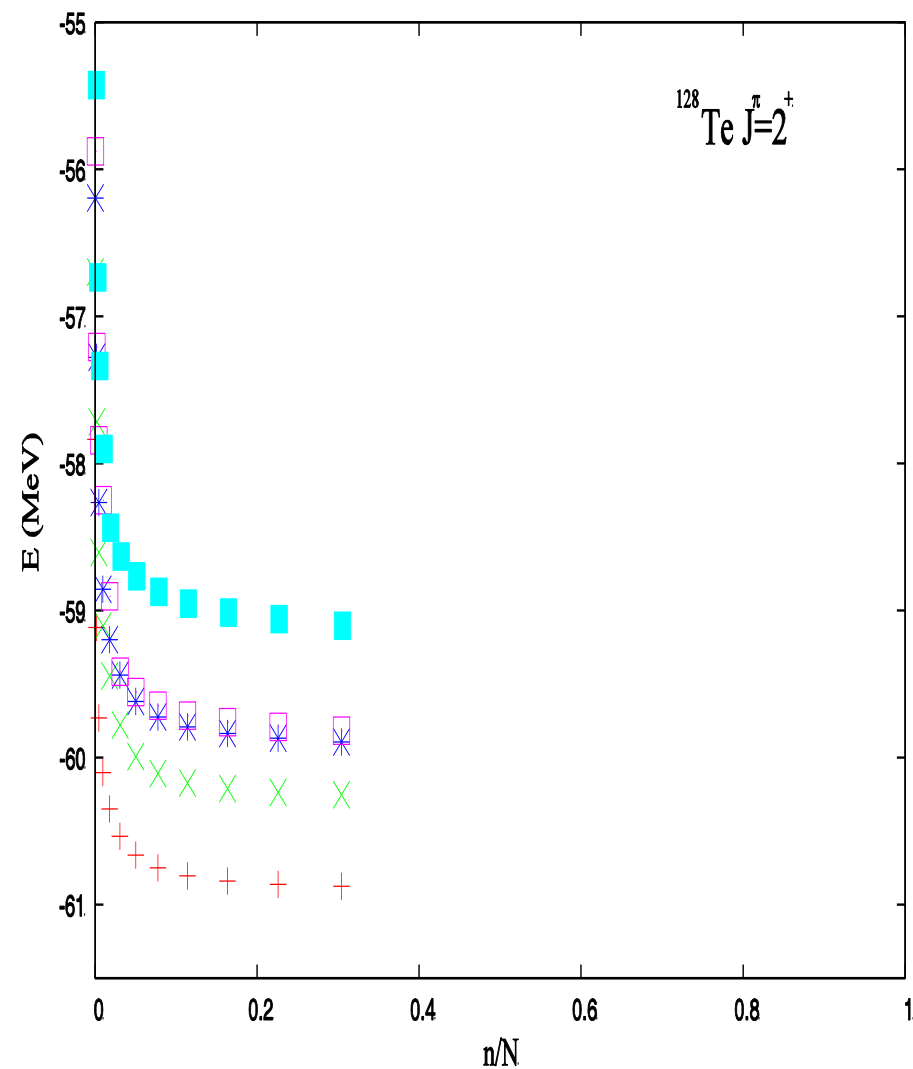
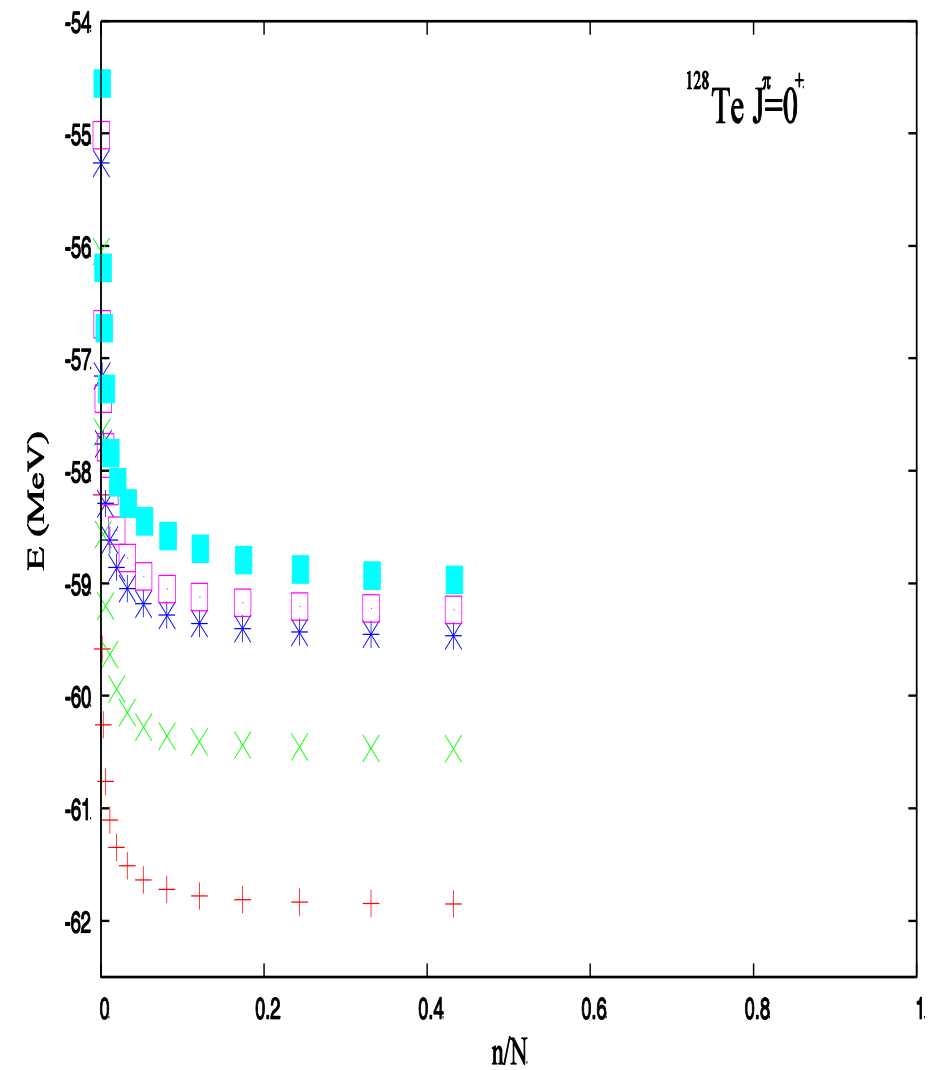
$$\Lambda_p^{(\varepsilon)}(v) \equiv \{(E_1^{(\varepsilon)}, \psi_1^{(\varepsilon)}) \dots (E_v^{(\varepsilon)}, \psi_v^{(\varepsilon)})\}$$

# Convergence properties: Eigenvalues



Semimagic nucleus:  $^{116}\text{Sn}$

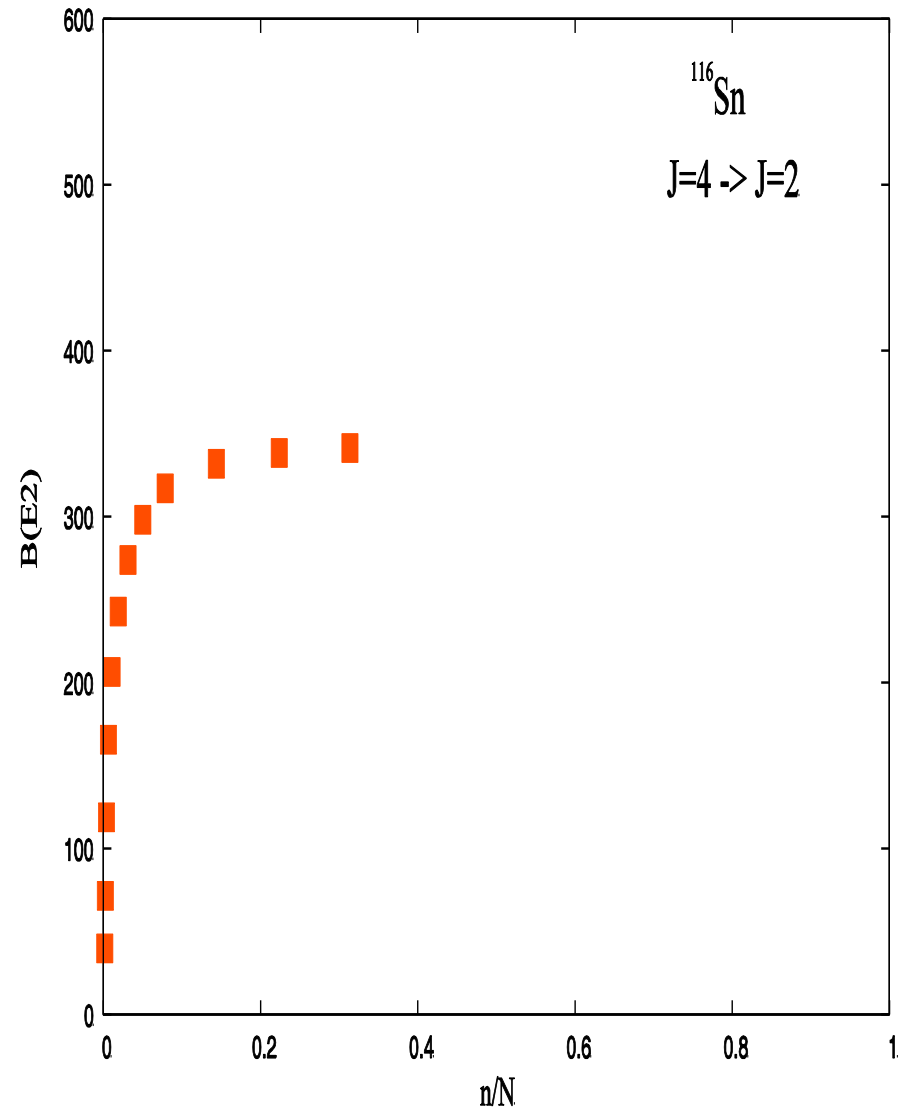
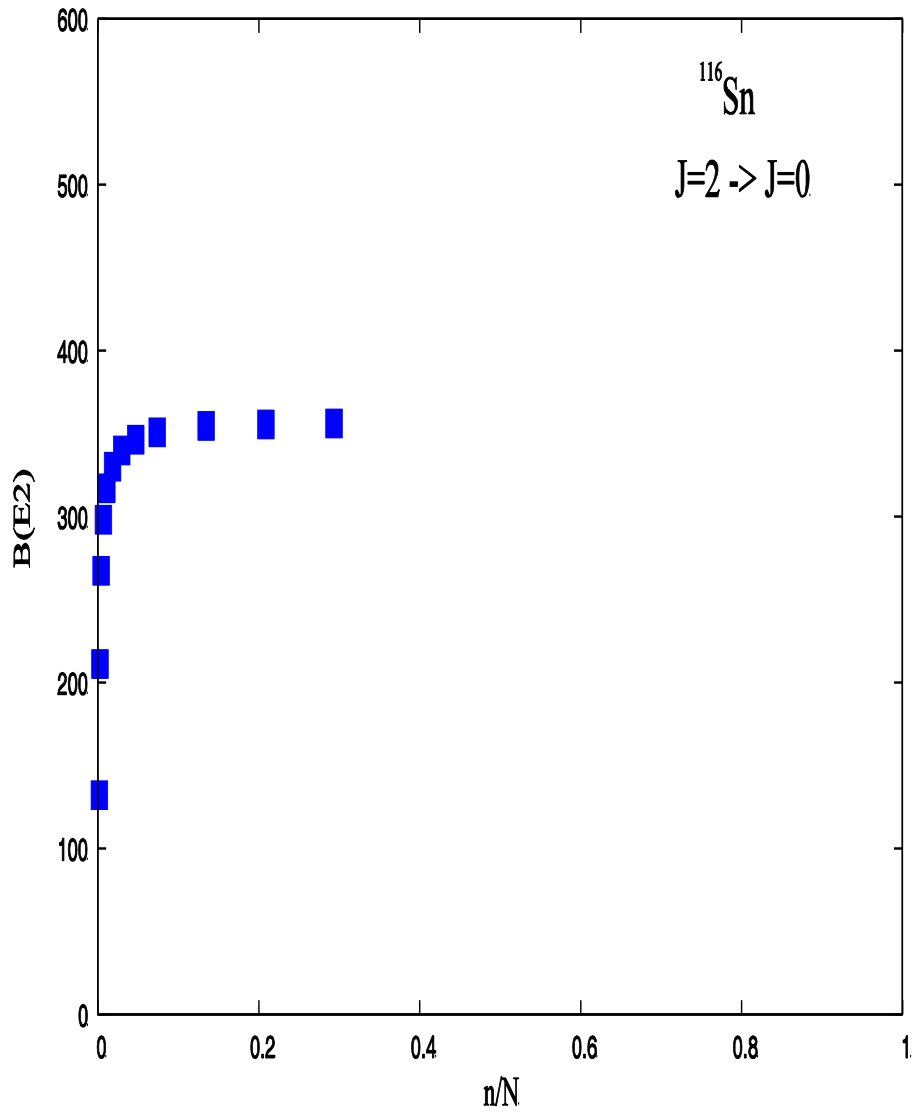
# Convergence of Eigenvalues



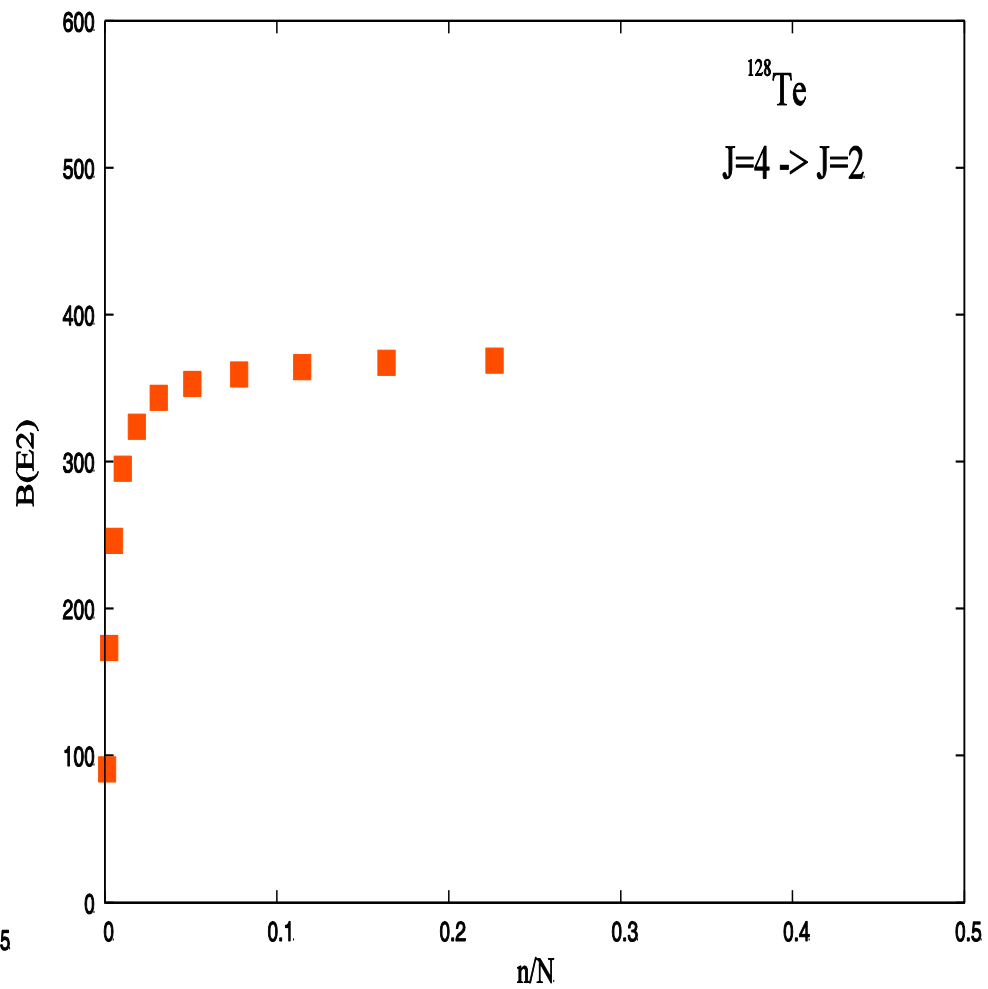
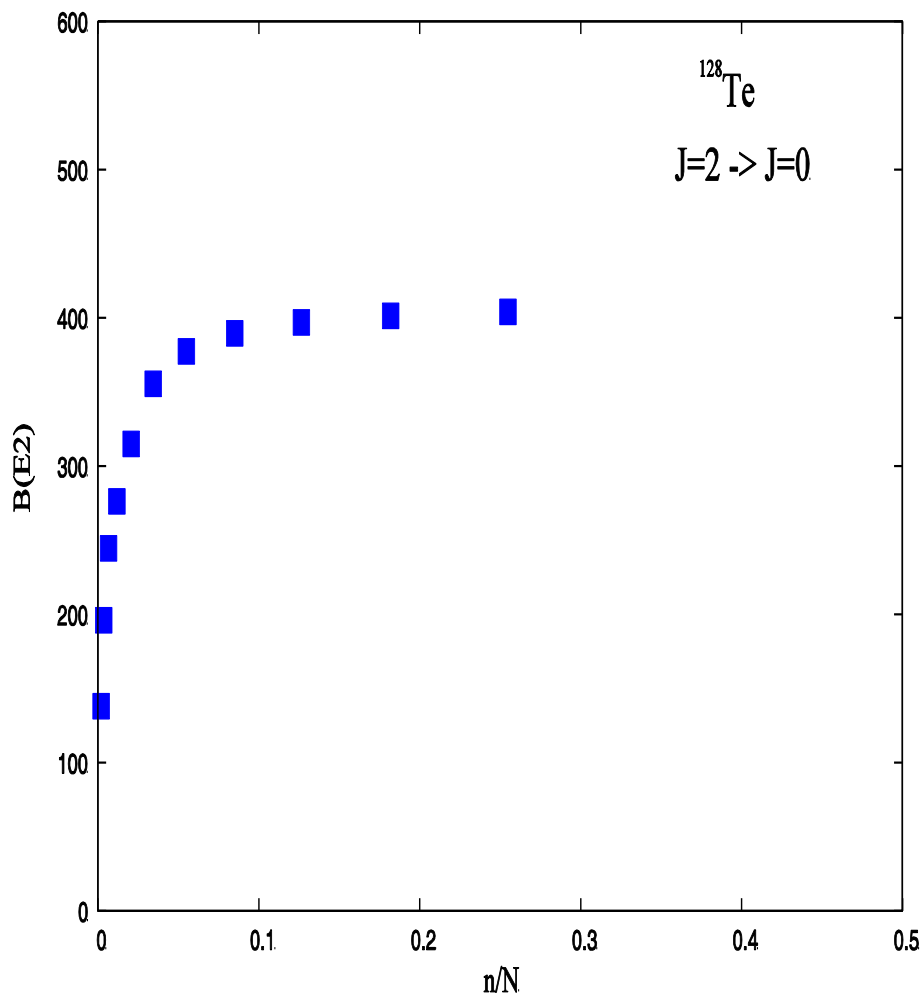
Open shell Nuclei :  $^{128}\text{Te}$



# Convergence of $B(E2)$ : $^{116}\text{Sn}$



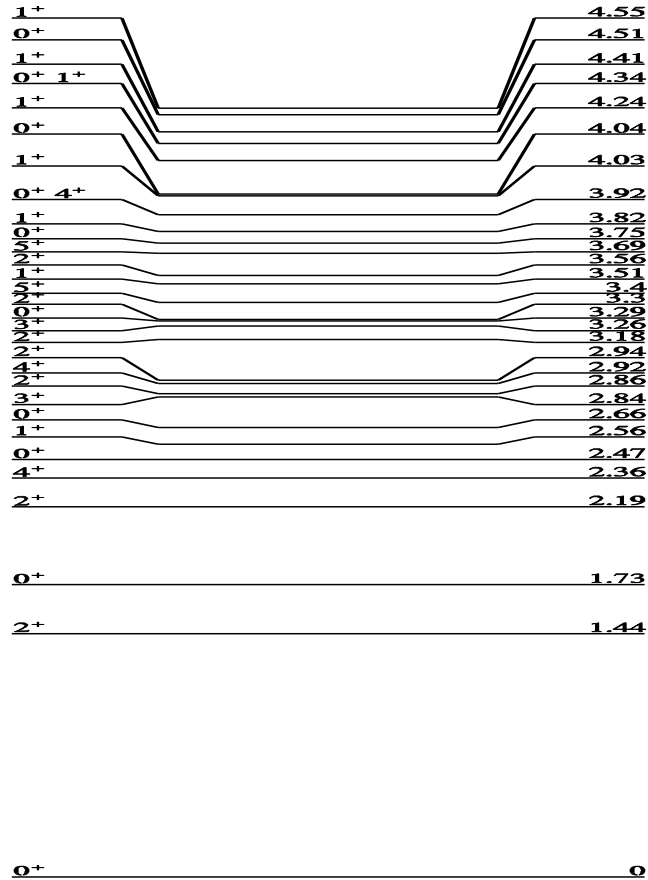
# Convergence of $B(E2)$ : $^{128}\text{Te}$



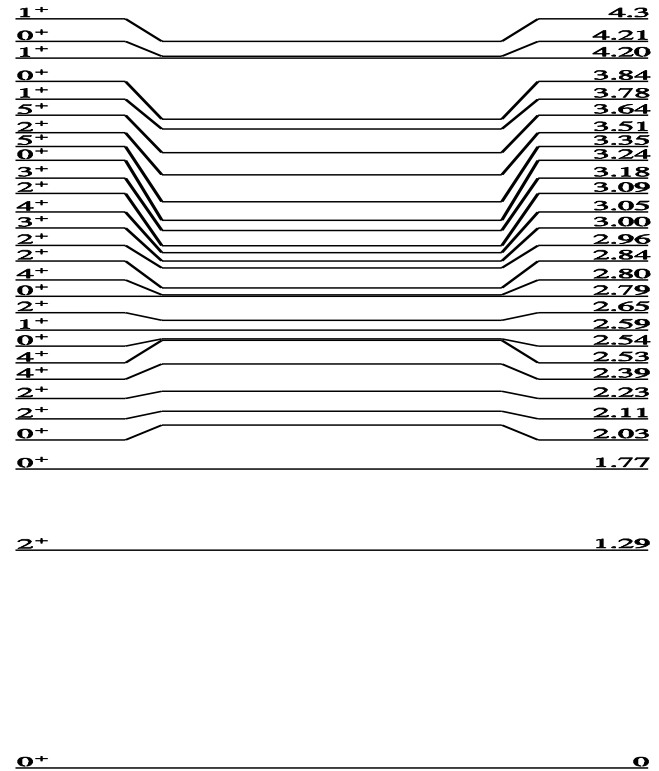
# Spectra

$^{116}\text{Sn}$

Th.



Exp.



# Conclusions

- The algorithm is **simple**, **robust** and has a **variational** foundation
- Once endowed with the **importance** sampling,
  - a) it keeps the extent of space **truncation** under strict **control**
  - b) It reaches saturation very early
  - c) it allows for **extrapolation** to exact eigensolutions

THANK YOU