



Statistical Model Calculations for (n,γ) Reactions

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Nuclear astrophysics



Waste transmutation



uses include



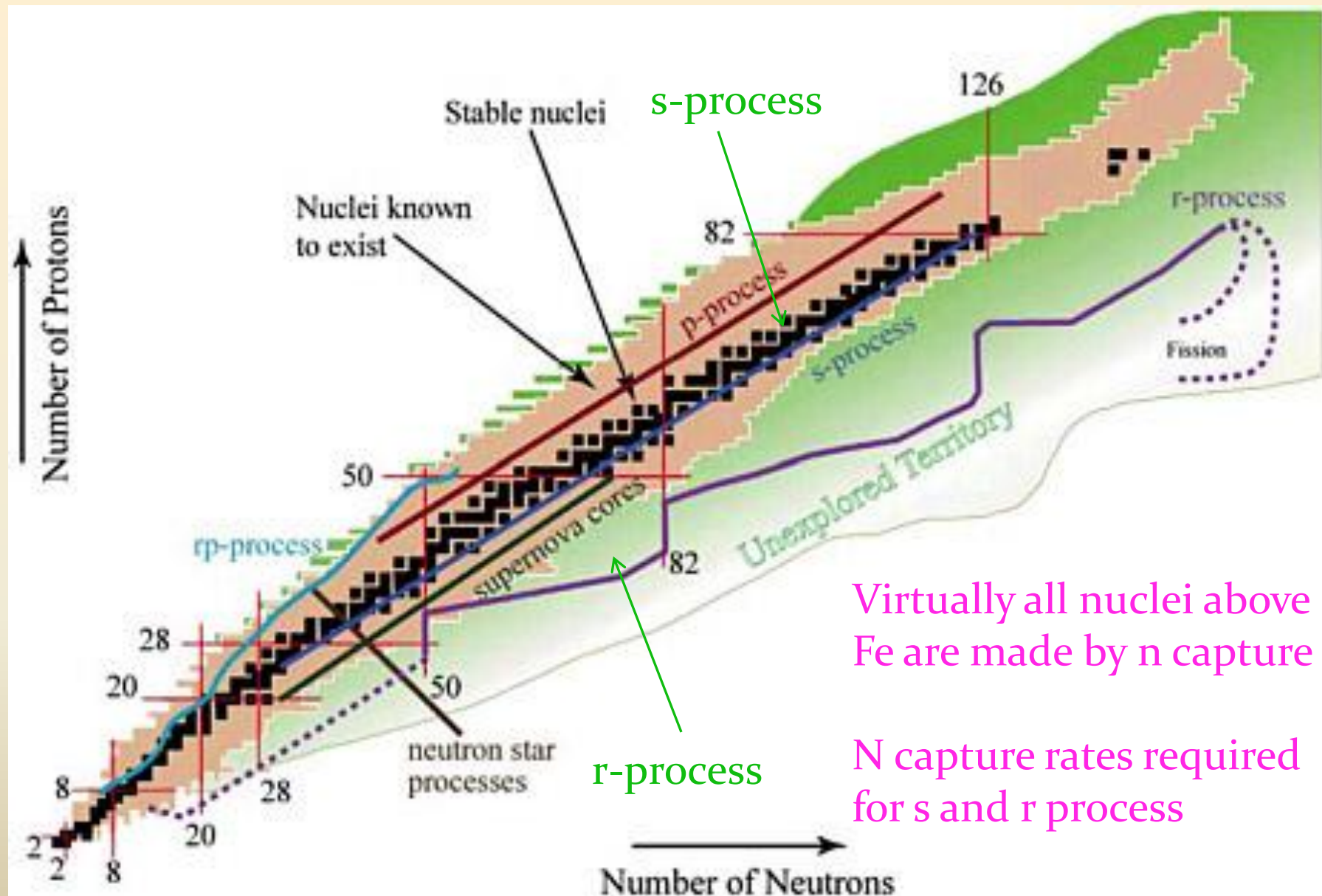
Reactor technologies

28 Aug 2014

Mary Beard CGS15, Dresden



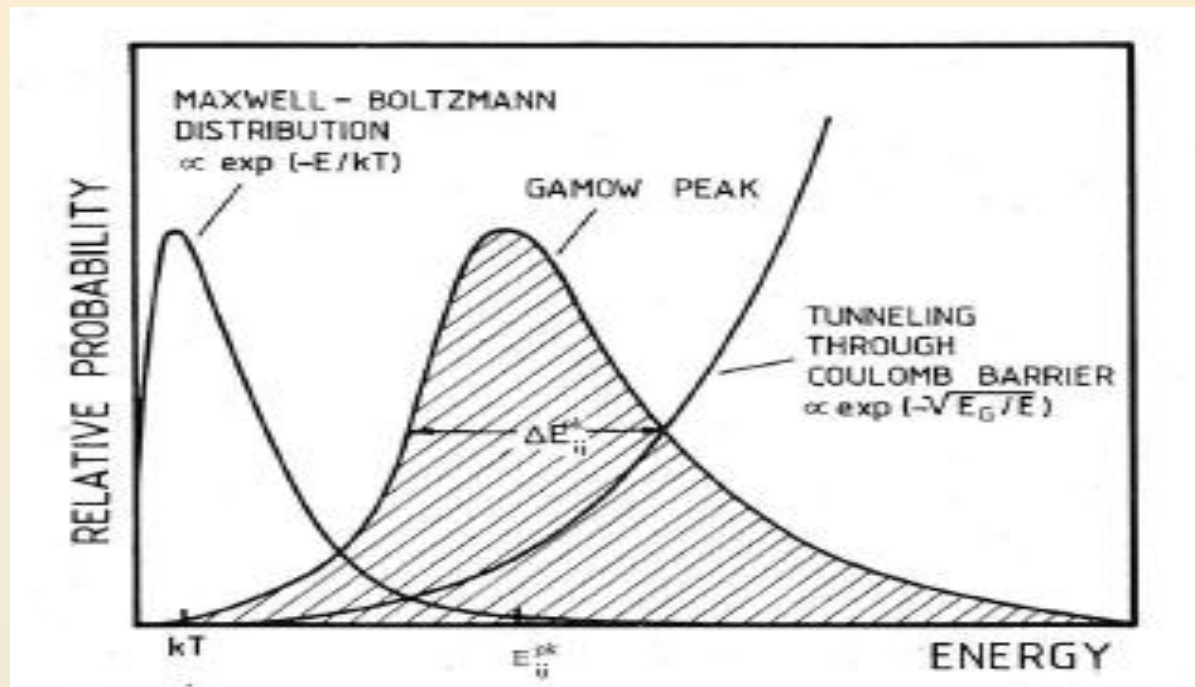
Medical applications



Virtually all nuclei above Fe are made by n capture

N capture rates required for s and r process

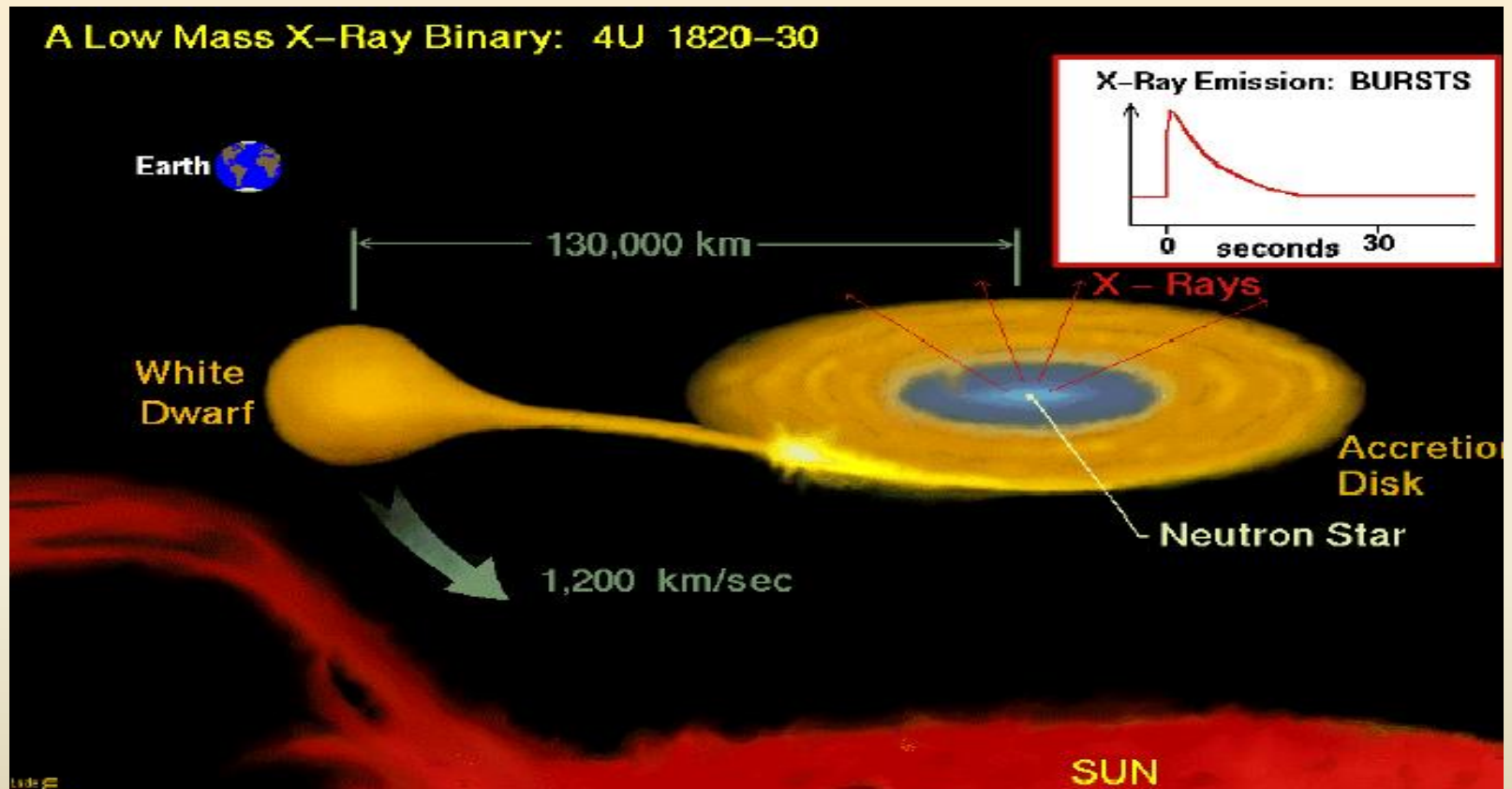
Most astrophysical sites need thermonuclear n capture rates



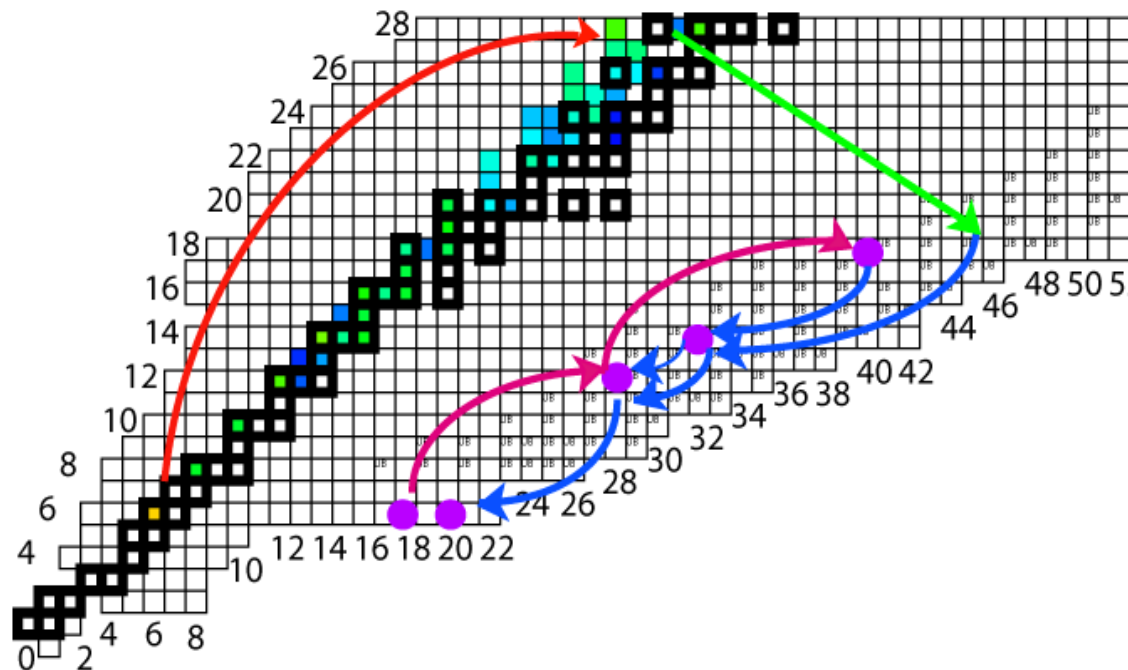
$$\langle \sigma v \rangle = \left(\frac{8}{\pi \mu_{ij}} \right)^{1/2} \frac{1}{(k_B T)^{2/3}} \int_0^{\infty} \sigma(E) E \exp\left(\frac{-E}{k_B T}\right) dE.$$

$$\text{MACS: } \langle \sigma v \rangle_T = \frac{\langle \sigma v \rangle}{\sqrt{2k_B T / \hat{A}}} \quad \rightarrow \text{KADoNiS}$$

Neutron Star Crust

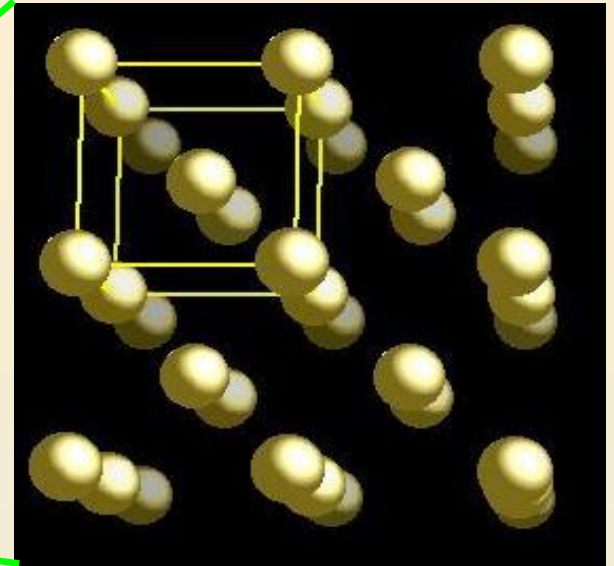
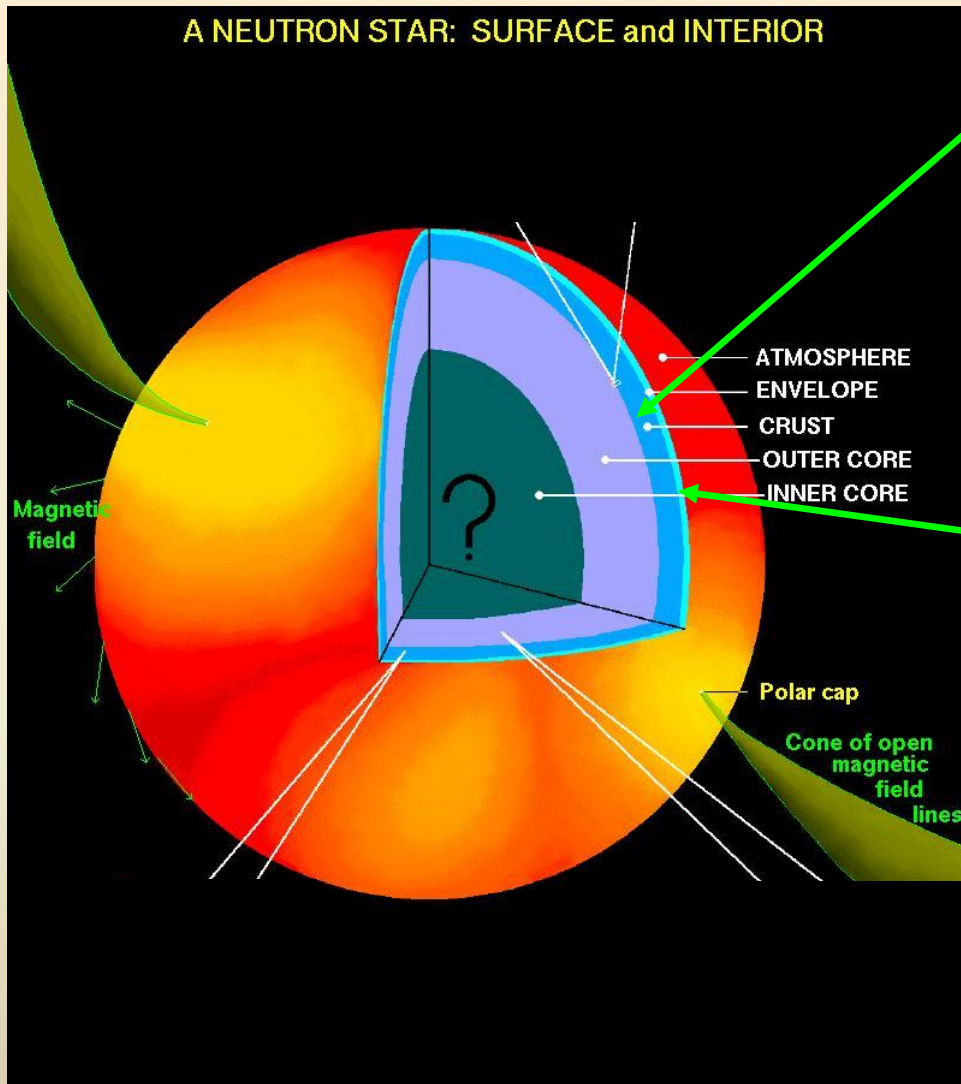


- r-p process path
- electron capture
- electron capture + neutron emission
- pycnonuclear fusion



- Matter continually compressed – ignites via triple α and hot cno breakout
- Follows rp-process path
- $\rho \sim 1.5 \times 10^9 \text{ g/cm}^3$ electron captures dominate
- At drip line neutron emissions process ashes into light material in the C to Mg range
- EC leads to high concentrations of free neutrons
- \rightarrow degenerate neutron capture

Neutron Star Environment



Tightly bound lattice structure

BCC, FCC, Imperfections, dislocations.....

Neutrons are confined in Coulomb crystal. → Not MB dist.

Degenerate N capture

$$N_A \langle \sigma v \rangle = K \int_0^{\infty} E \sigma(E) f(E, T) dE$$

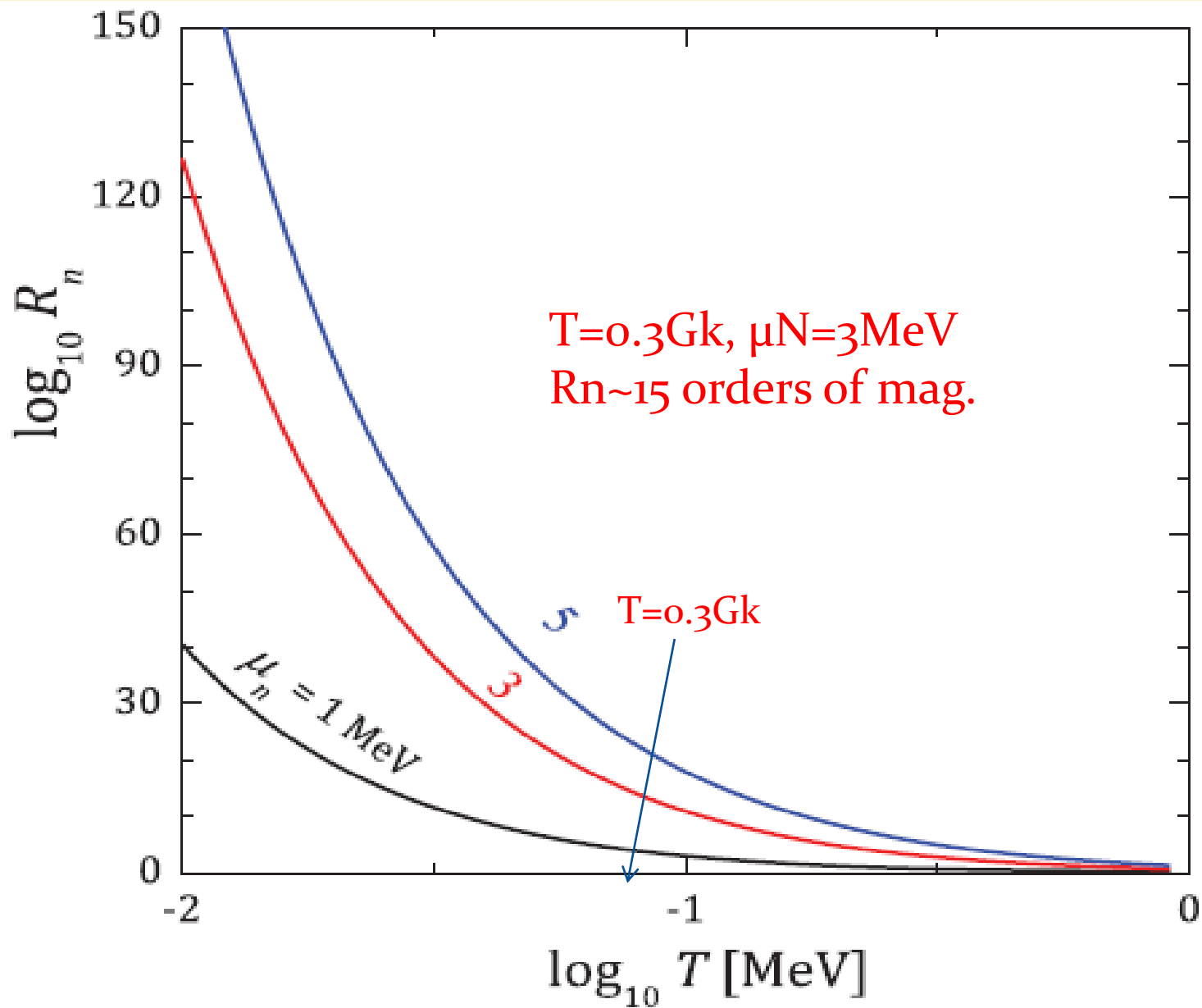
$$f_{FD} = \frac{1}{1 + e^{(E - \mu_N)/T}}$$

Introduce correction
parameter R_n :

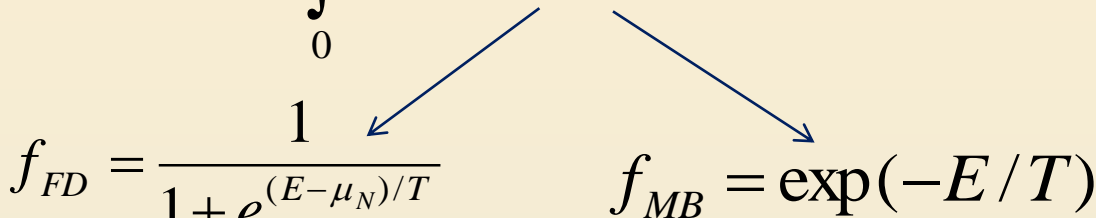
$$R_n = \frac{\langle \sigma v \rangle_{FD}}{\langle \sigma v \rangle_{MB}}$$

Measure of how
applicable MB rates
are

REACLIB, BRUSLIB....



Astrophysical networks require 1000's of rates!

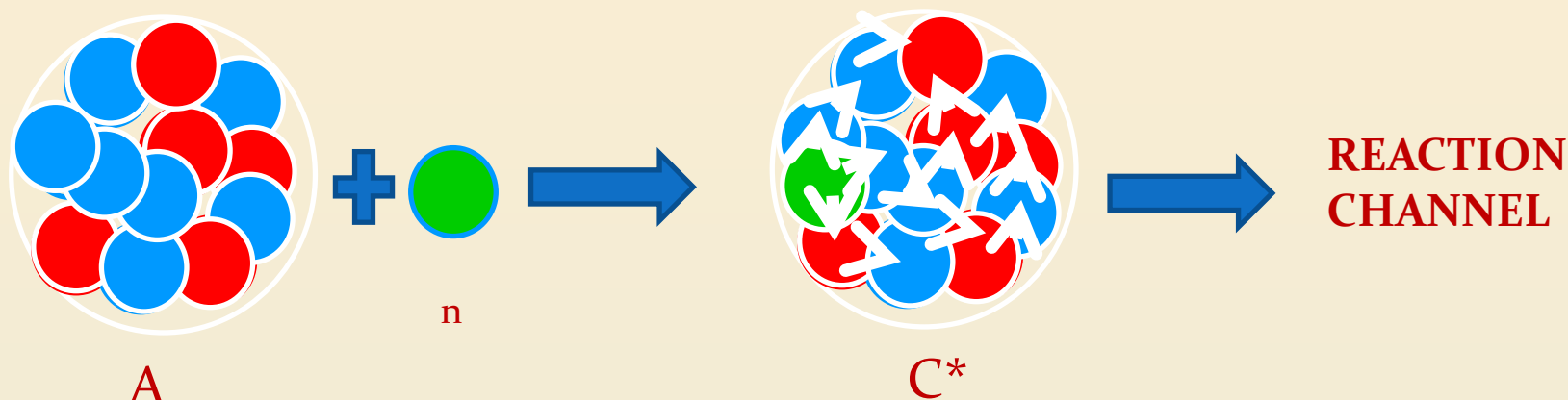
$$N_A \langle \sigma v \rangle = K \int_0^{\infty} E \sigma(E) f(E, T) dE$$


The diagram shows two blue arrows originating from the integrand $E \sigma(E) f(E, T)$ of the equation above. One arrow points to the Fermi-Dirac distribution function $f_{FD} = \frac{1}{1 + e^{(E - \mu_N)/T}}$, and the other points to the Maxwell-Boltzmann distribution function $f_{MB} = \exp(-E/T)$.

$$f_{FD} = \frac{1}{1 + e^{(E - \mu_N)/T}} \quad f_{MB} = \exp(-E/T)$$

$\sigma(E)$ typically obtained from HF calculations

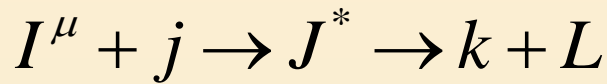
The Statistical Model: $N_A \langle \sigma v \rangle = K \int_0^\infty E \sigma(E) f(E, T) dE$



Reaction takes place via intermediary C^*
 C^* must reach a state of thermodynamic eq.

Statistical model normally valid if LD in C^* is high, so that resonances overlap

HF Reaction σ



Reaction Cross Section needs

- transmission co-efficients \rightarrow n, gamma

$$\sigma_{jk}^\mu(E) = \pi \hat{\lambda}_j^2 \frac{1}{(2J_I^\mu + 1)(2J_j + 1)} \sum_{J,\pi} (2J + 1) \frac{T_j^\mu(J^\pi) T_k(J^\pi)}{T_{tot}(J^\pi)}$$

Spins :
tar., proj., & comp.

Transmission coefficient
(make/destroy J^π
from $I^\mu + j$)

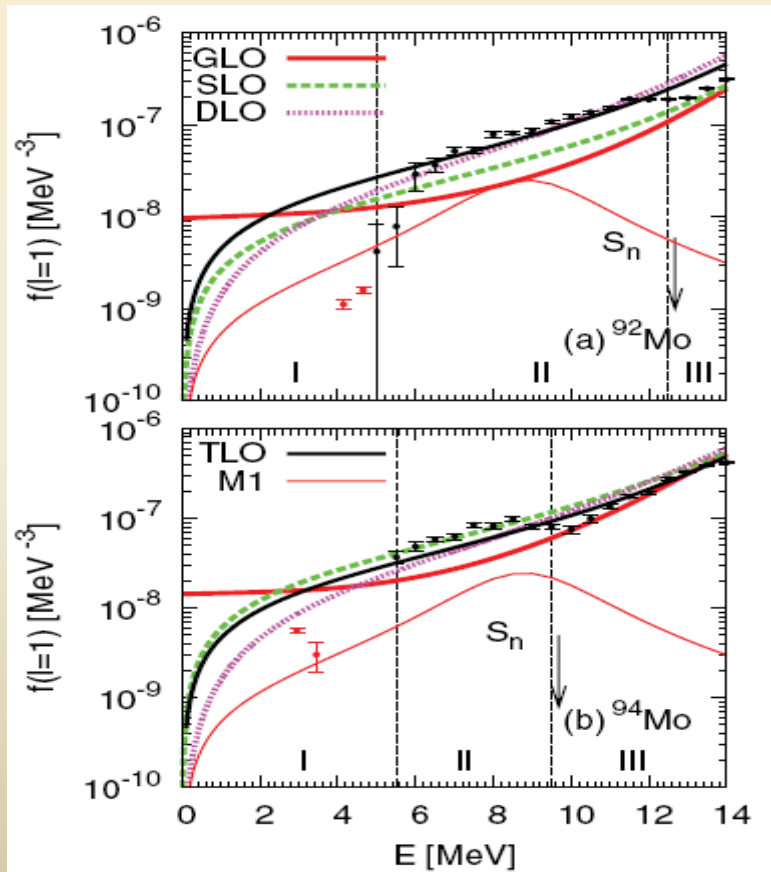
Transmission coefficient
(make/destroy J^π
from $k + L$)

Total decay transmission coefficient
(any channel)

Transmission Functions

Particle – OMP, LD (often BSFG , CT)

Photon – GDR. Normally E1 dominates. See talk by S. Frauendorf



$$T_{E1}(E_\gamma) = 2\pi E_\gamma^3 f_{E1}(E_\gamma)$$

Parameterized
Lorentzian(s)

Ingredients

data

- masses
- deformations
- $J\pi$, E_x

models

- level density
- γ -strength function
- optical model

Calculating Trans. requires nuclear inputs

→ HF calculations are sensitive to:
choice of models
quality/quantity of data

XS uncertainties

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graph TD; A([XS uncertainties]) --> B[Is CN model applicable?<br/>•Pre-equilibrium<br/>•direct reactions]; A --> C[Model/data uncertainties<br/>•Masses<br/>•Deformation<br/>•Jπ<br/>•Level densities<br/>•GSF<br/>•OMP]; D[3rd source:] --> E[•Code approximations<br/>•Implementation details (matching E, GDR param, a, δ....)];
```

Is CN model applicable?

- Pre-equilibrium
- direct reactions

Model/data uncertainties

- Masses
- Deformation
- $J\pi$
- Level densities
- GSF
- OMP

3rd source:

- Code approximations
- Implementation details (matching E, GDR param, a, δ)

To test the 3rd point, xs calculations have been performed with:

TALYS

SAPPHIRE

(Statistical Analysis for Particle and Photon capture and decay of High energy REsonances)

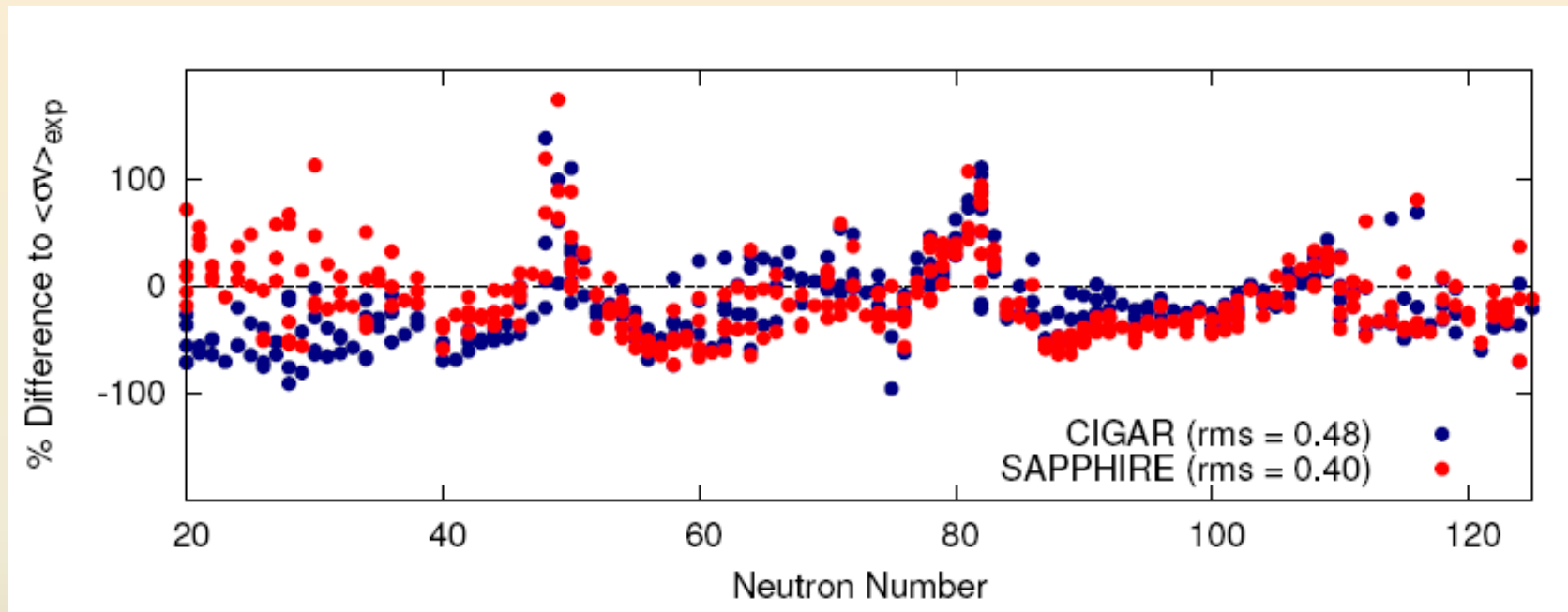
CIGAR (Capture Induced Gamma-ray Reactions)

Contain identically implemented nuclear models

Calculations performed for ~340 nuclei in KADoNiS database

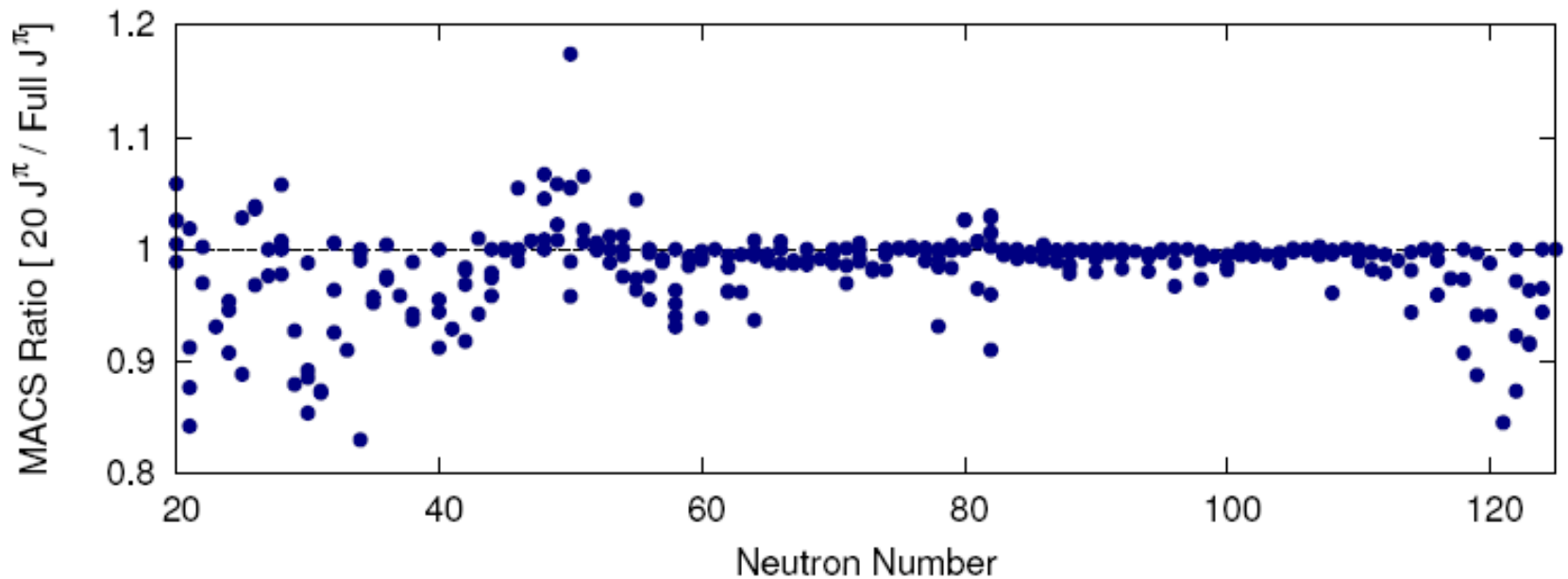
→calculated MACS @ 30keV compared to exp. MACS

CIGAR vs SAPHIRE



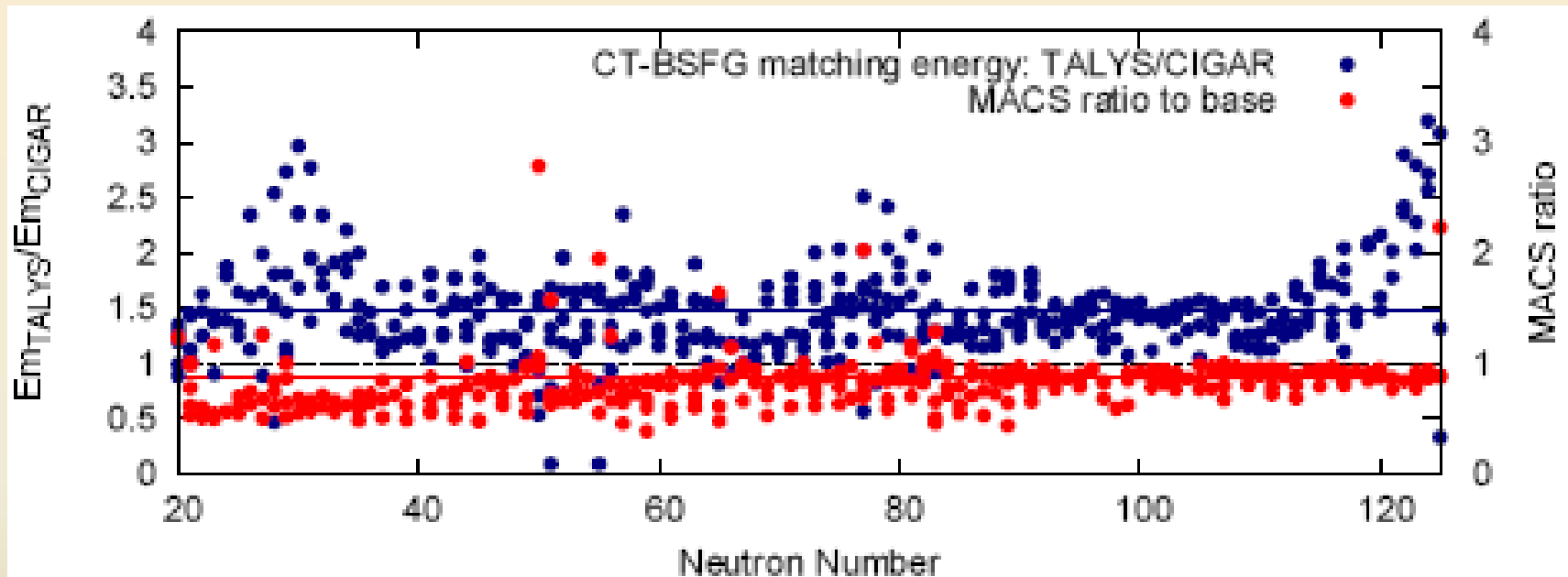
Identical nuclear models in both codes. Differences due to $J\pi$ data truncation

Code approximation:
 $J\pi$ truncation



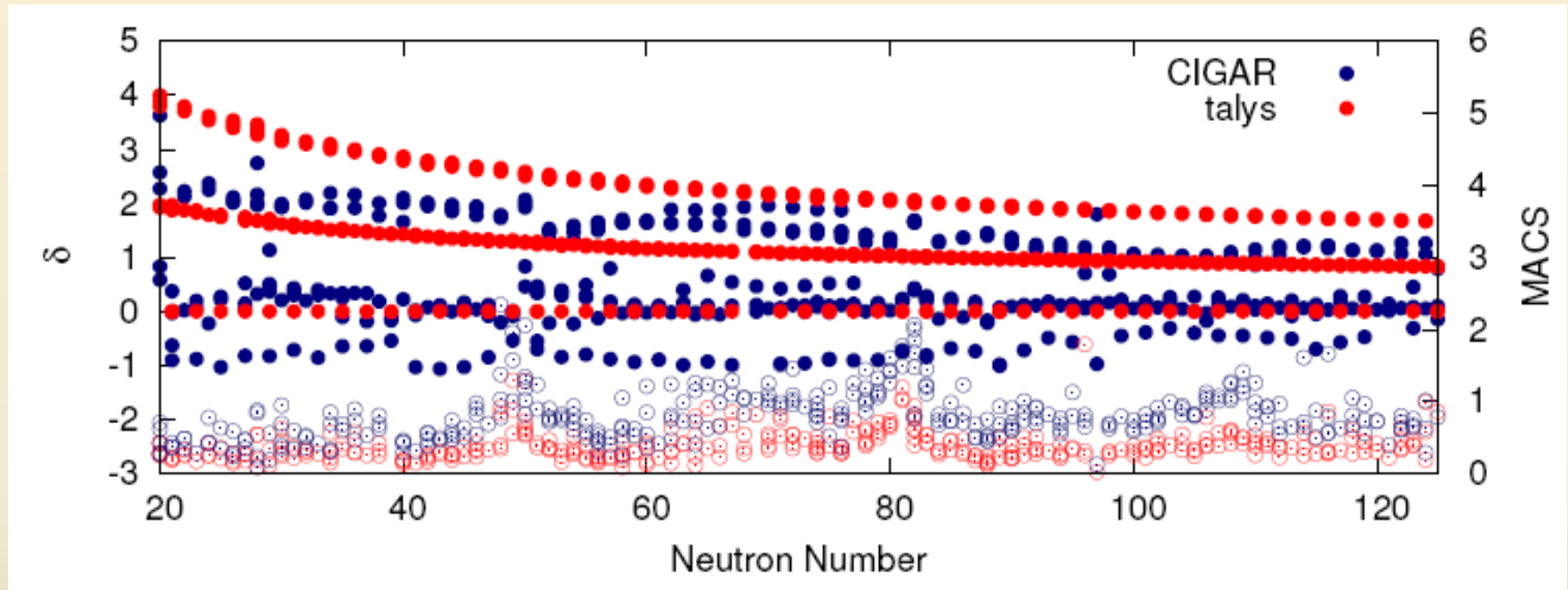
→ Truncating $J\pi$ data use can impact calculations

Model Implementation: CT – BSFG matching E



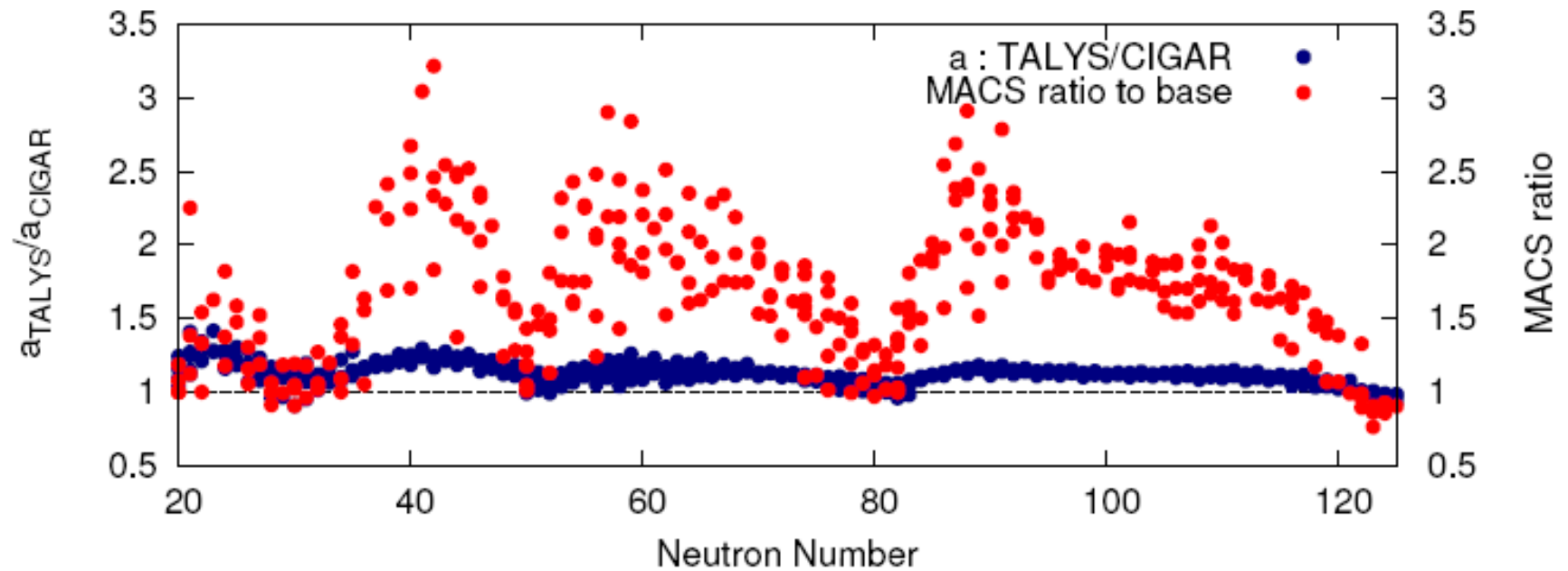
Larger Em leads to modest reduction of MACS

Model Implementation: LD model BSGF δ



Both codes use BSGF, but different def. of back-shift
→ Reflected in MACS

Model Implementation: LD parameter a



Conclusions

HF rates are integral to huge number of applications

Large number of different types of uncertainties,
good to \sim factor 3....but that is not all models

- But also from code assumptions and implementation

Collaborators

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