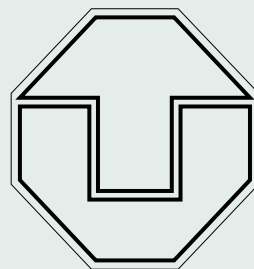


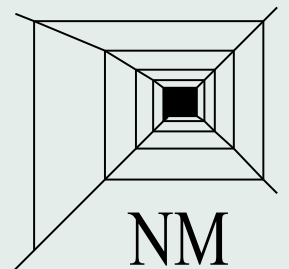
International Workshop on FLOW Control by Tailored MAGnetic Fields
Open and closed-loop Control of Weakly Conducting Fluids

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Technische Universität Dresden
Institut für Numerische Mathematik



**TECHNISCHE
UNIVERSITÄT
DRESDEN**



Rossendorf, April 2, 2004

Outline:

Part 1

2position p-controller

Part 2

Nonlinear feedback control recipes

Motivation

Given:

a validated mathematical model for a physical process.

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Goal:

Model based control of the physical process

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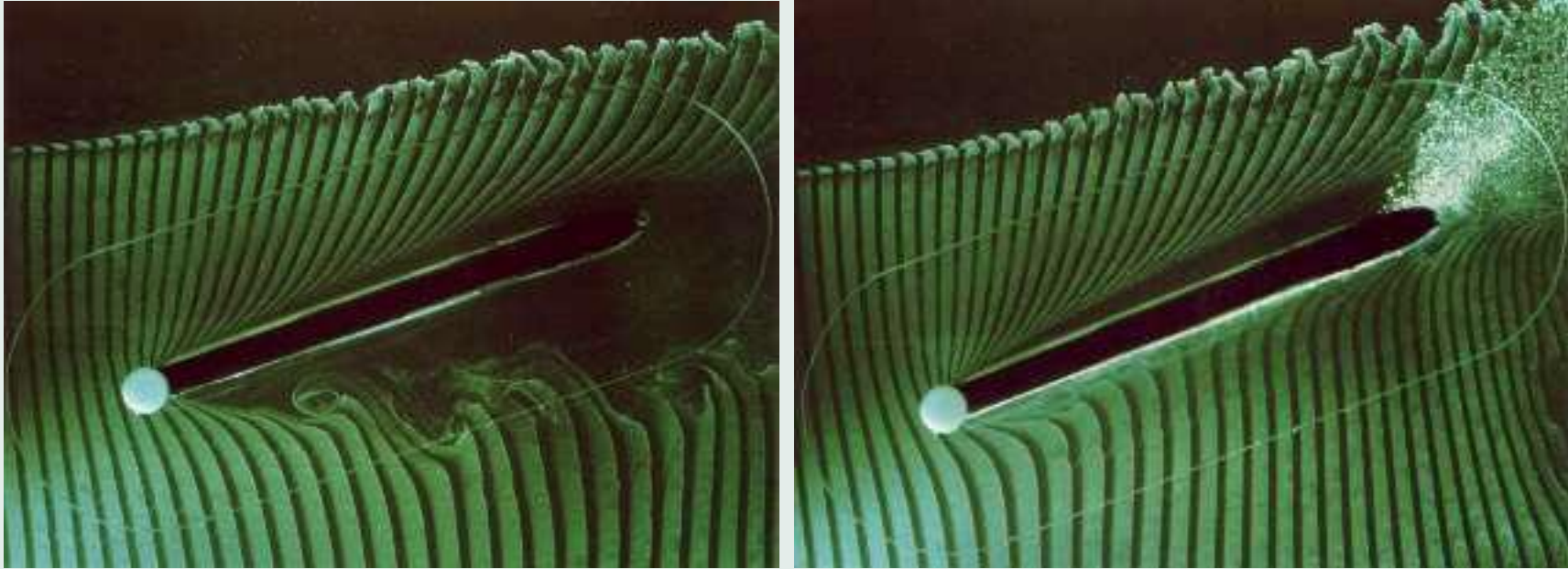
conducted by requirements of an experiment

Applications

Control of weakly conducting fluids(Sfb 609)

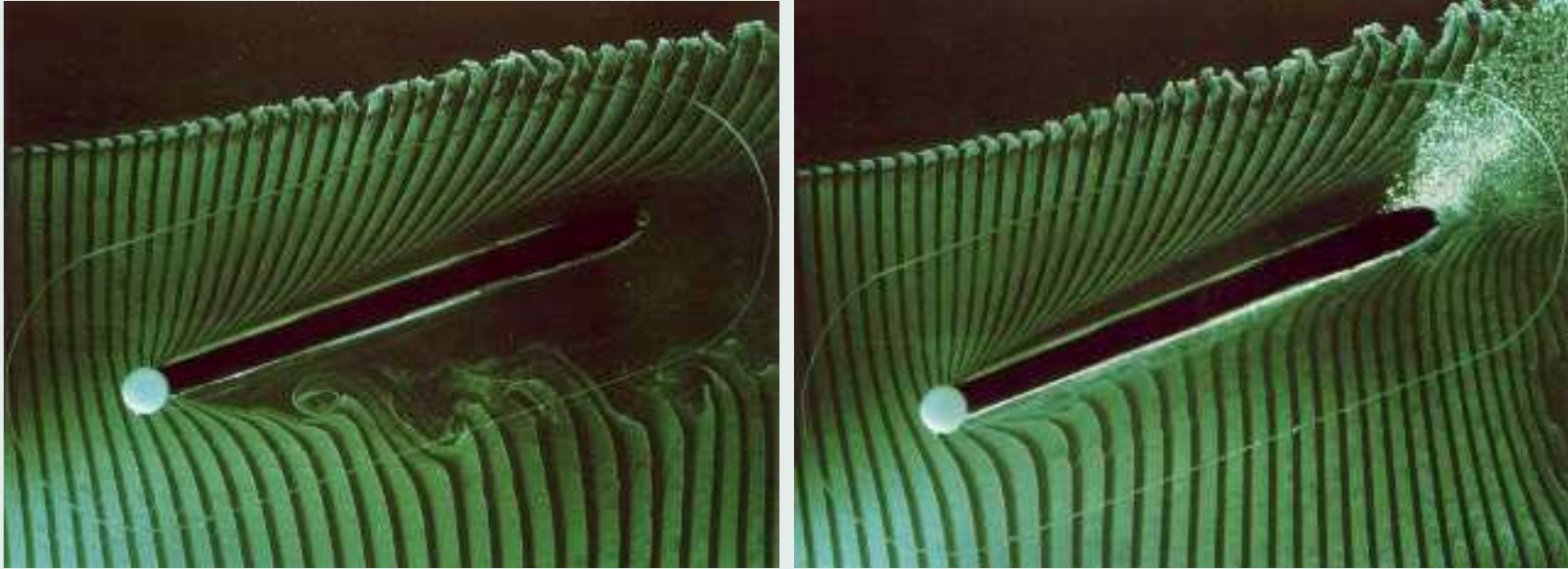
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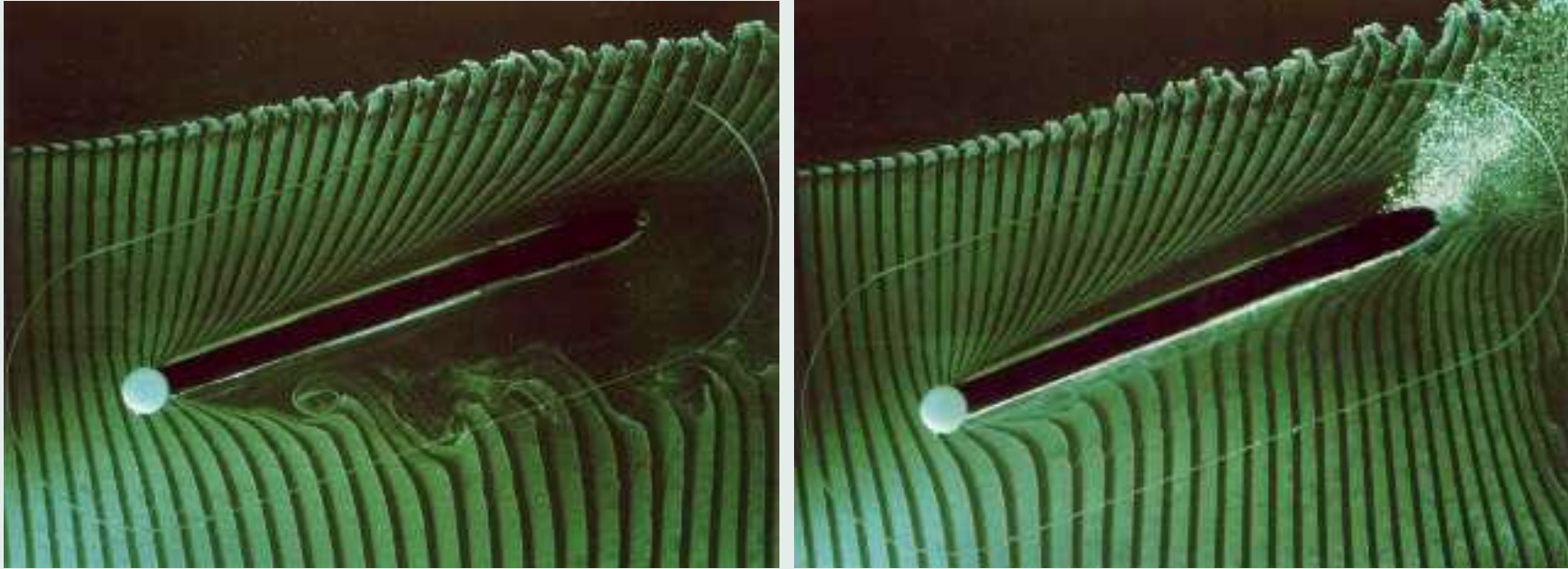
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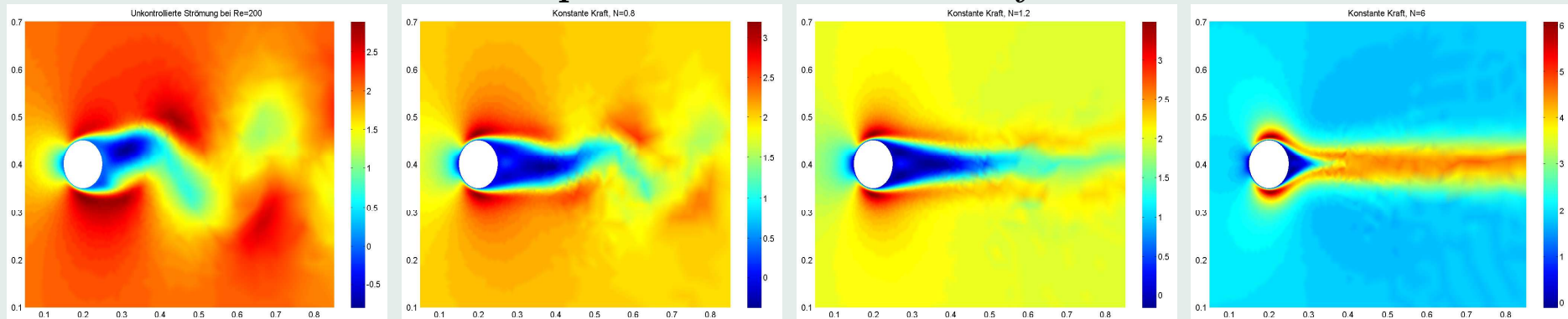
Develop methods for circular cylinder

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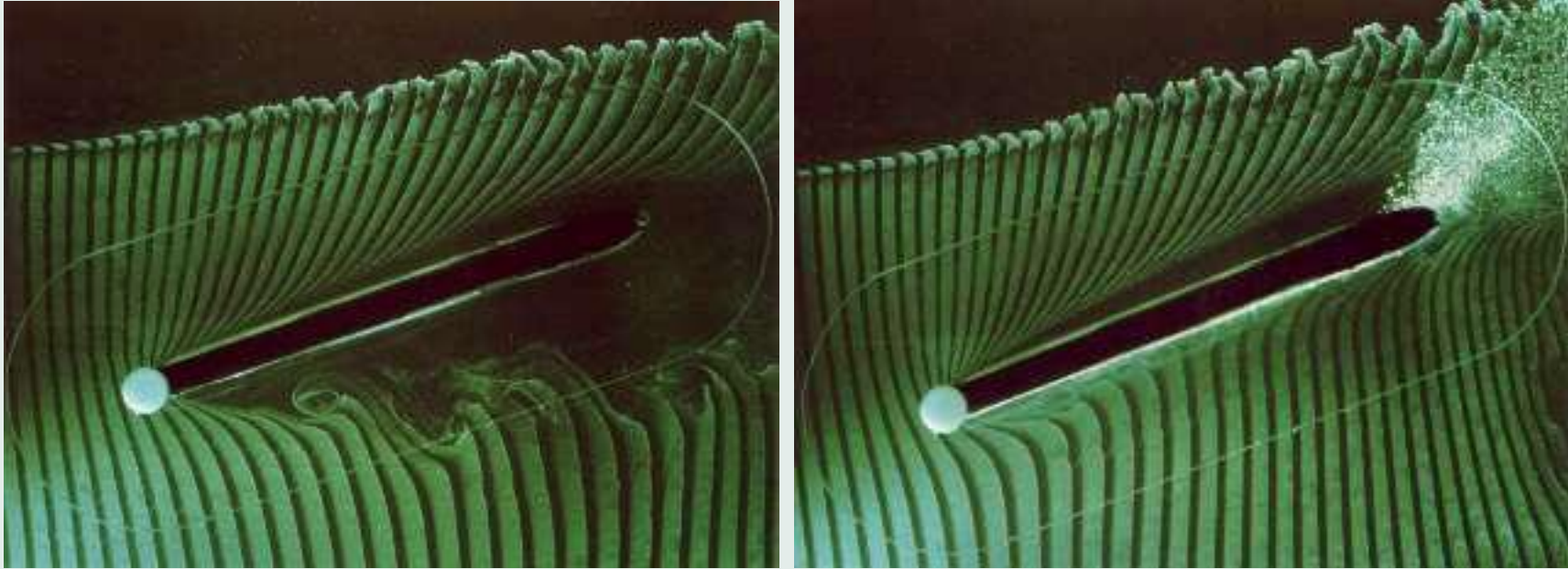
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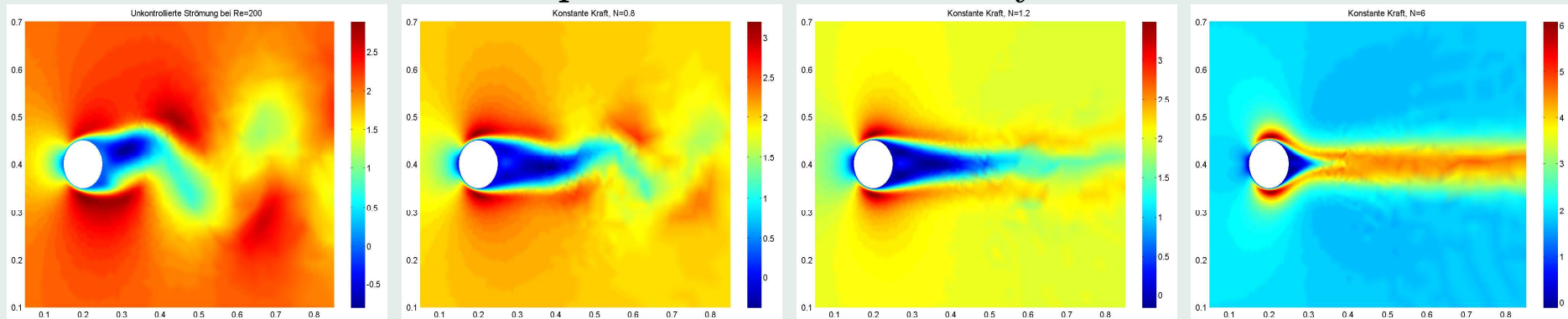
Control target: Re-attach flow utilizing near wall Lorentz forces

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Control target: Re-attach flow utilizing near wall Lorentz forces

Desired: Closed-loop control strategy

Control strategies (problem depending)

- Optimal (robust) open-loop control
- (robust) feedback control

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3 concepts

- Design of a two-position controller on the basis of the mathematical model
- model-predictive control, instantaneous control
- model-reduction and control

Part 1: Control of the flow past a cylinder,
conducted by the experiment.

On the surface of the cylinder a Lorentz force is generated with actuators. It decays exponentially in the flow (T. Berger et al., Phys. Fluids 2000).

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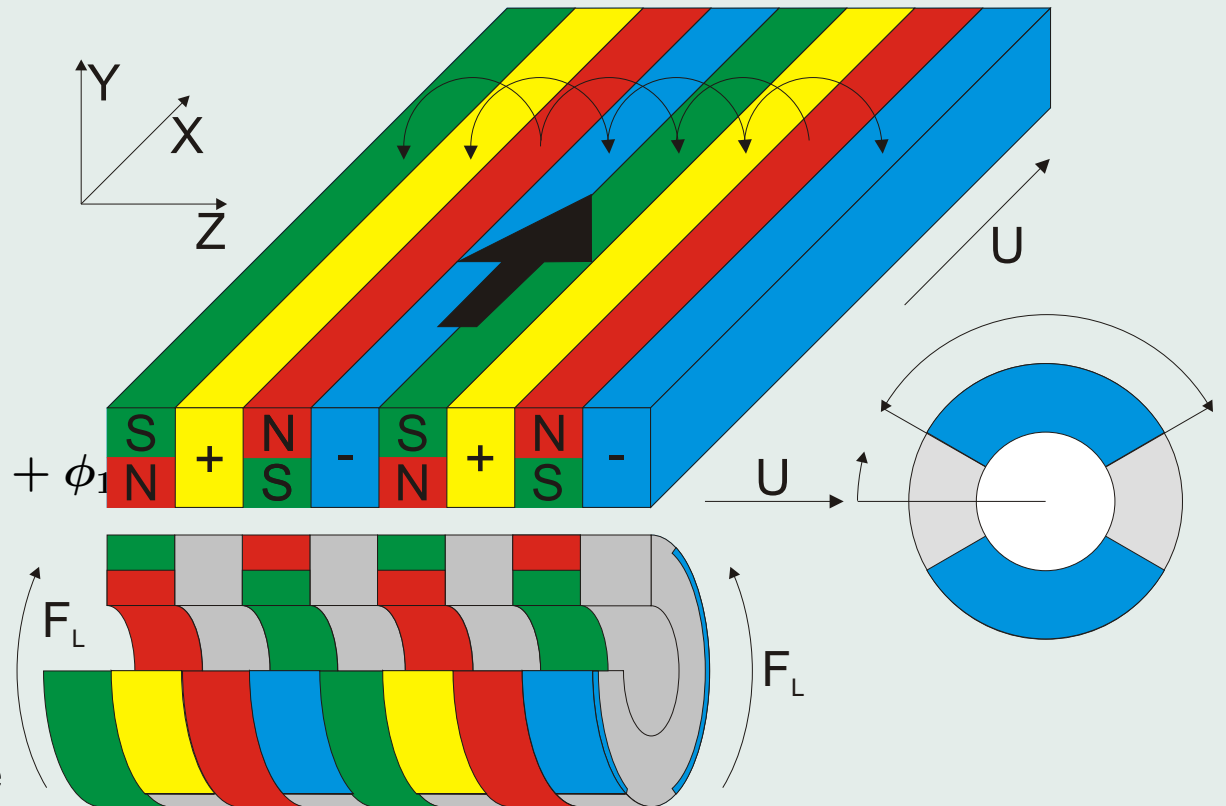
The Lorentz force has the form

$$F_L(x, y) := e_\phi g(\phi) e^{-\frac{\pi}{a} \cdot \text{dist}[(x,y), \text{cylinder}]},$$

$$g(\phi) = \begin{cases} 1, & \phi_0 \leq \phi \leq \phi_1 \\ -1, & 180^\circ + \phi_0 \leq \phi \leq 180^\circ + \phi_1 \\ 0, & \text{else} \end{cases}$$

a electrode-/magnet-spacing.

The direction of the Lorentz Force is defined by the arrangement of the magnets/electrodes.



Navier-Stokes-Equations

$$\begin{aligned}u_t + (u \cdot \nabla)u + \nabla p - \frac{1}{\text{Re}} \nabla^2 u &= N F_L \\ \nabla \cdot u &= 0\end{aligned}$$

plus initial and boundary conditions.

u denotes the flow velocity, p the pressure, and $\text{Re} = \frac{U_\infty D}{\nu}$ the Reynolds number, with ν the constant kinematic viscosity, D the cylinder diameter and U_∞ the characteristic velocity. The Interaction parameter N describes the ratio of the electromagnetic and the inertial forces of the flow,

$$N = \frac{J_0 B_0 D}{\rho U_\infty^2} ,$$

with J_0 the current density and B_0 magnetic induction.

The numerics are validated against the computations of R. Grundmann/O. Posdziech (TUD). Comparison with experiments by G. Gerbeth/T. Weier et al. (FZR) to be addressed.

Implementation in FEMLAB.

Spatial discretization: Taylor-Hood elements.

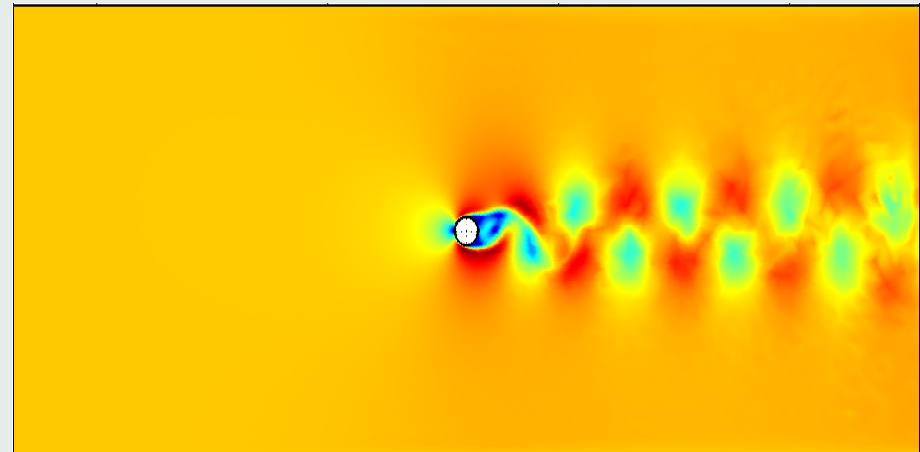
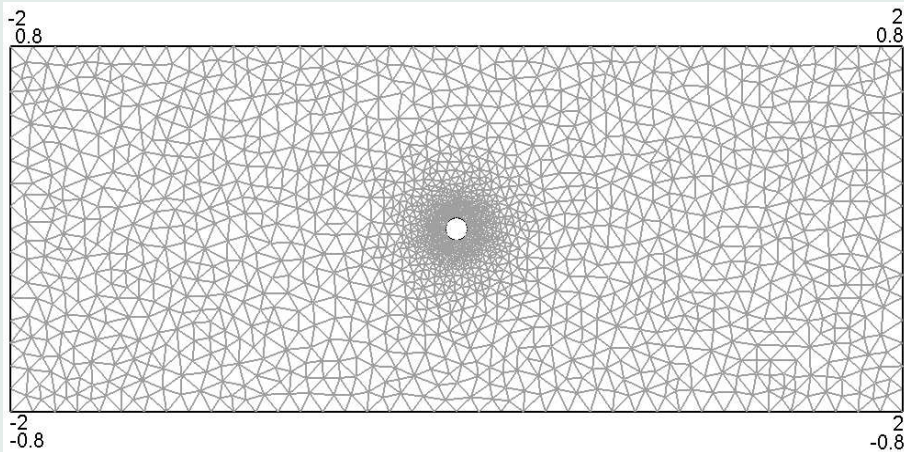
Integration schema: BDF (or other).

Inflow condition: Block/parabolic profile. Cylinder surface: no-slip.

Upper and lower boundary: slip/no-slip.

Outflow: do-nothing.

Computational domain:



The large (in relation to the cylinder) domain is needed in order to compute the correct drag coefficients.

The drag is defined as:

$$F_D = F_{D_f} + F_{D_p} + F_{D_{em}} = \underbrace{\int_{\partial cyl.} \rho \nu \frac{\partial u_t}{\partial n_y} dS}_{=F_{D_f}} - \underbrace{\int_{\partial cyl.} p n_x dS}_{=-F_{D_p}} - \underbrace{J_0 B_0 \frac{a}{\pi} D (\cos \phi_0 - \cos \phi_1)}_{=-F_{D_{em}}},$$

in which u_t is the tangential velocity and n_i denote the components of the normal vector. With F_D the drag coefficient C_D is given by:

$$C_D = \frac{2F_D}{\rho U_\infty^2 D}, \quad C_{D_f} = \frac{2F_{D_f}}{\rho U_\infty^2 D}, \quad C_{D_p} = \frac{2F_{D_p}}{\rho U_\infty^2 D}, \quad C_{D_{em}} = \frac{2F_{D_{em}}}{\rho U_\infty^2 D} .$$

To determine the Strouhal Number the **frequency of vortex separation** f is used,

$$St = \frac{fD}{U_\infty}$$

St is specified by the Lift force:

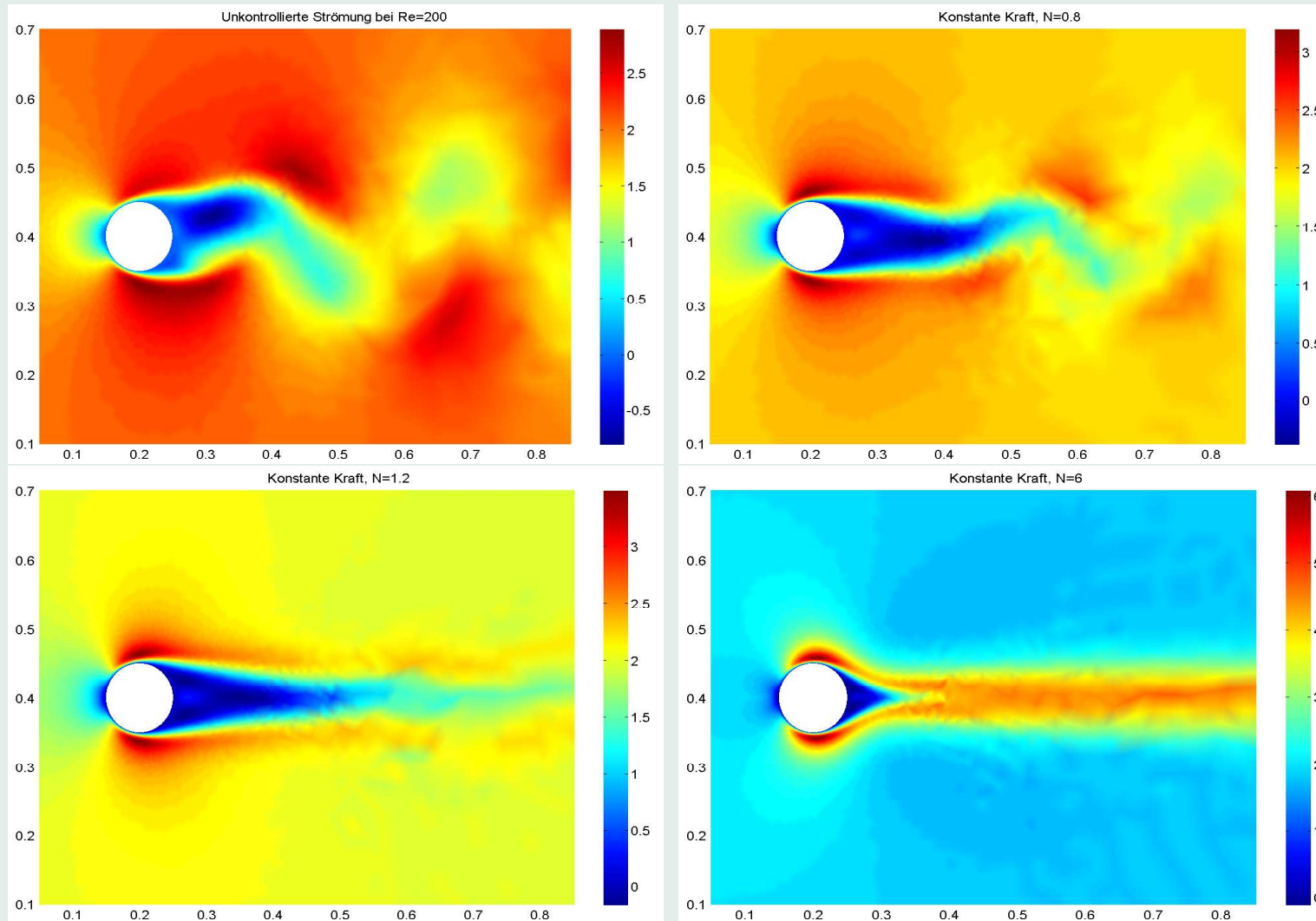
$$F_l = - \int_{\partial cyl.} \rho \nu \frac{\partial u_t}{\partial n_x} + p n_y dS$$

The lift coefficient is given by:

$$C_L = \frac{2F_L}{\rho U_\infty^2 D} .$$

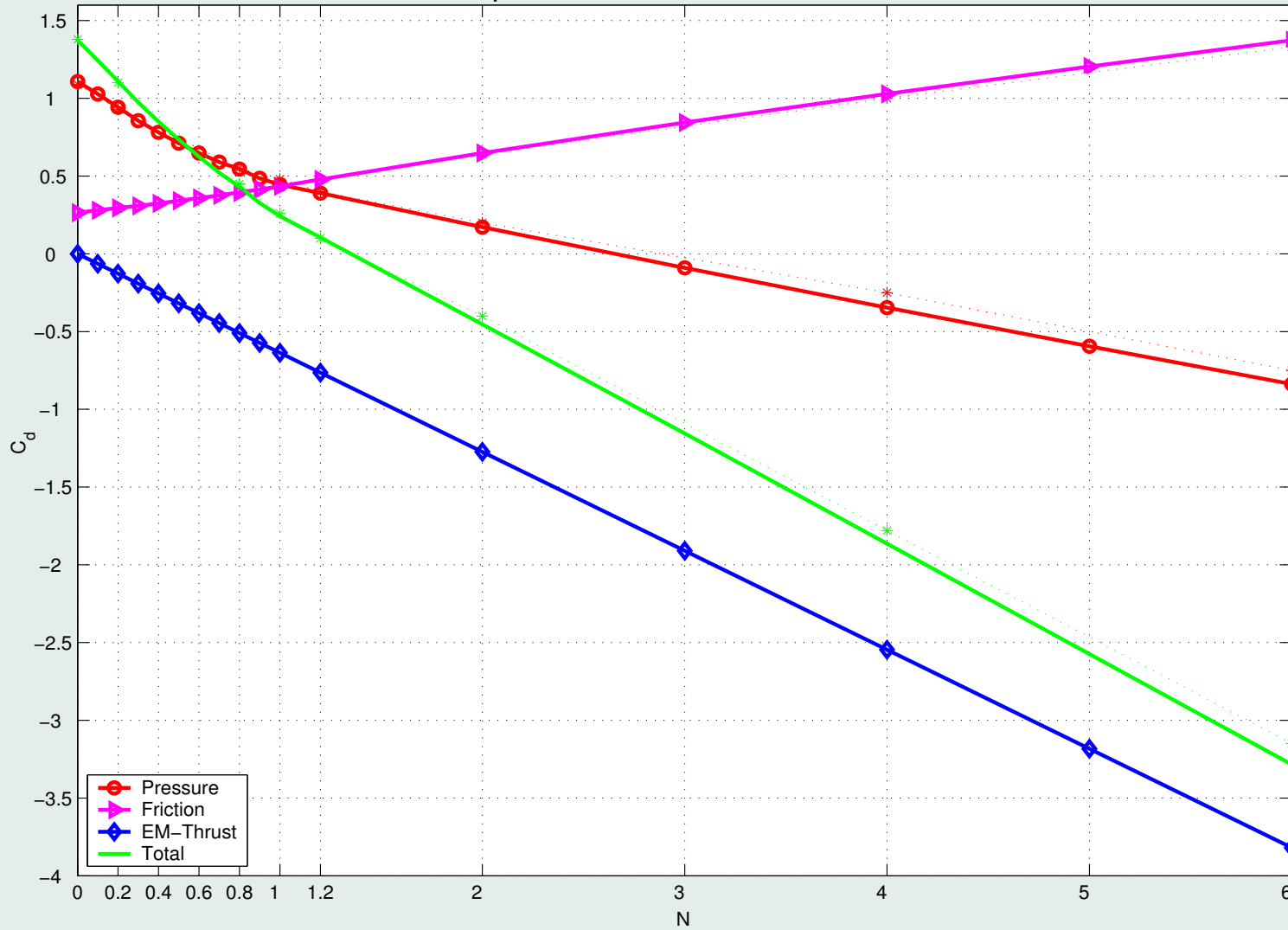
$Re = 200$, stepsize $\Delta t = 0.01$

The formation of the vortex street is suppressed by application of a tangential Lorentz force:

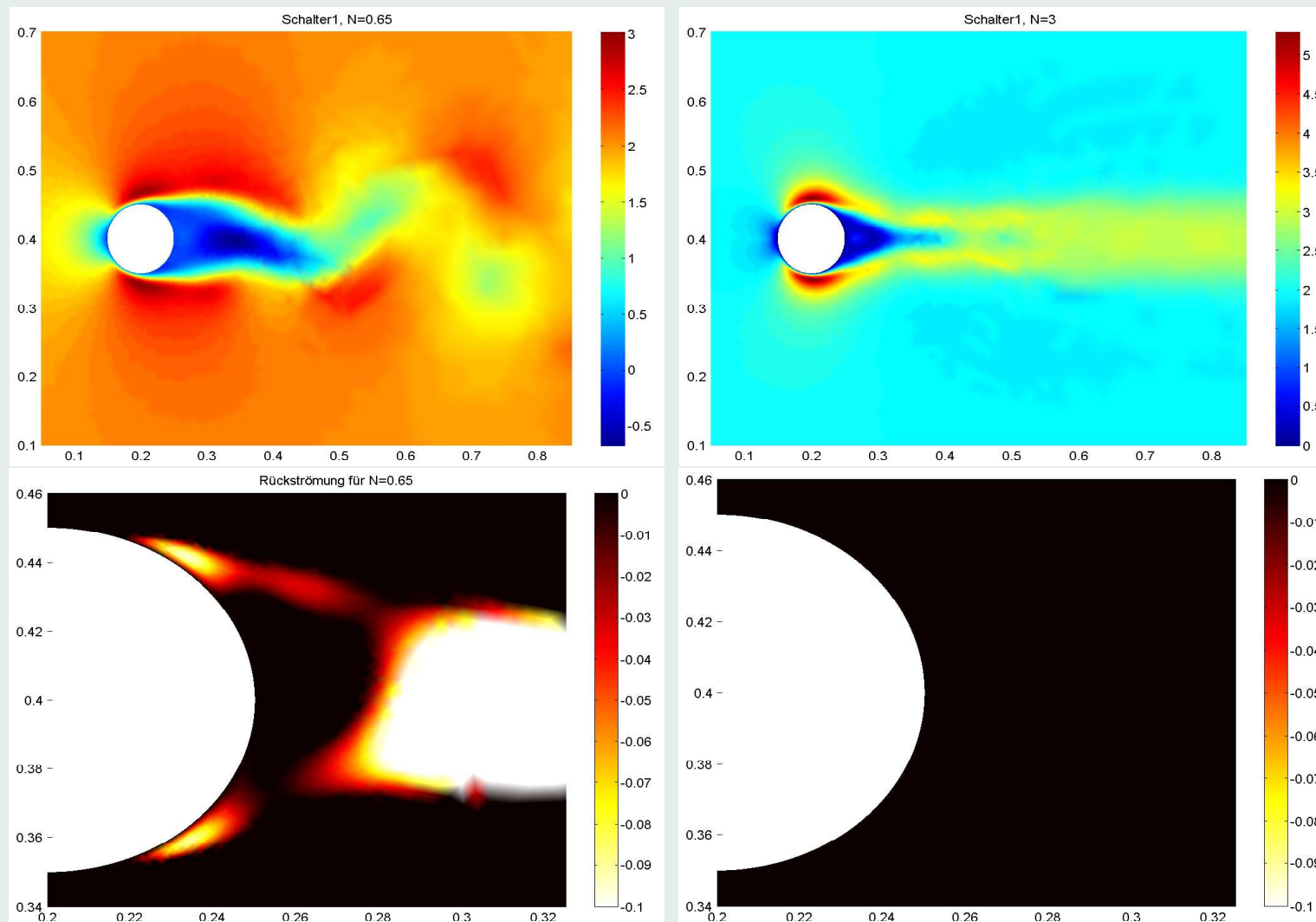


Drag Coefficients:

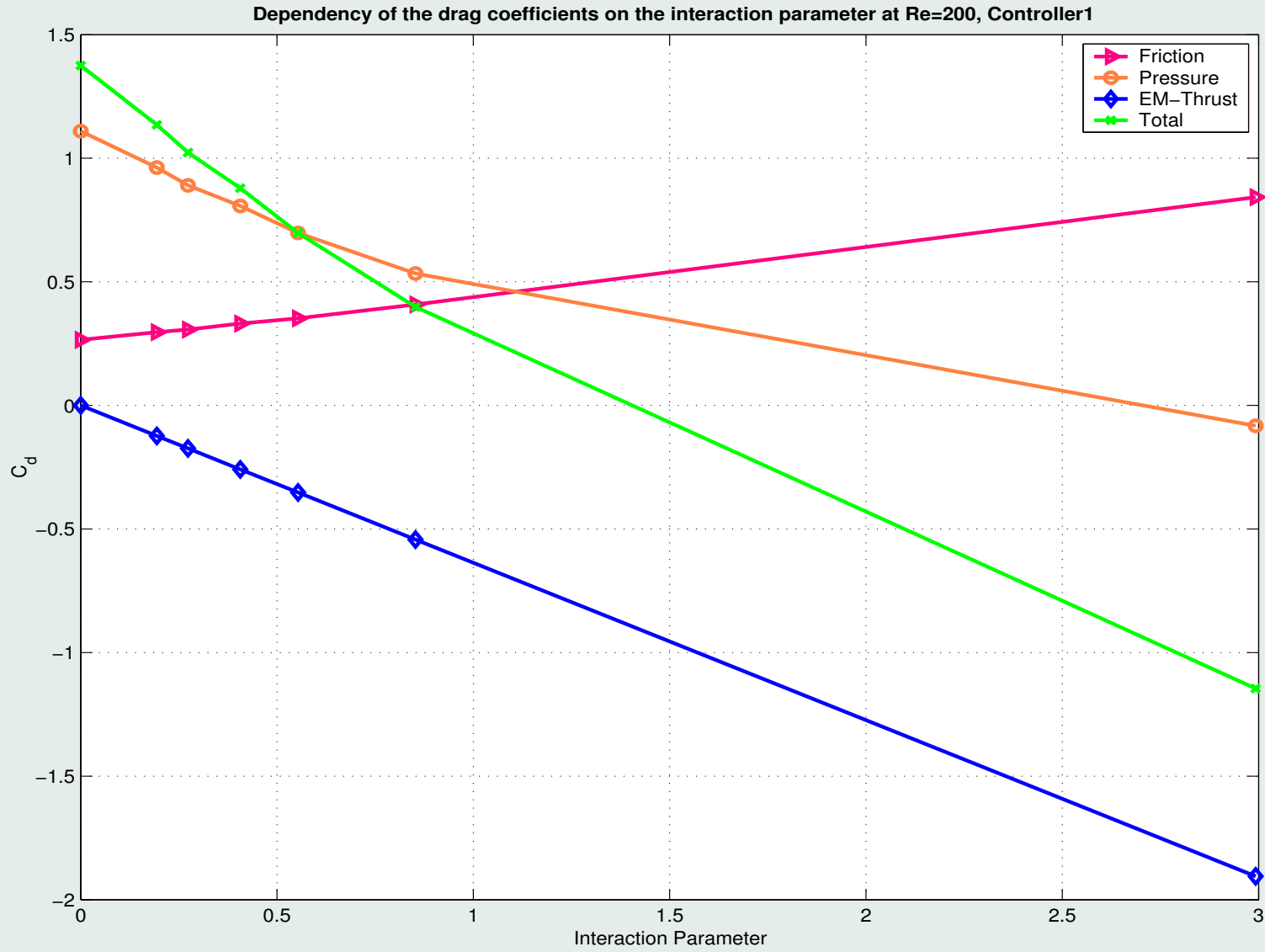
Figure 8: Dependency of the drag coefficients on the interaction parameter at Re=200
Comparison with Grundmann & Posdziech's results



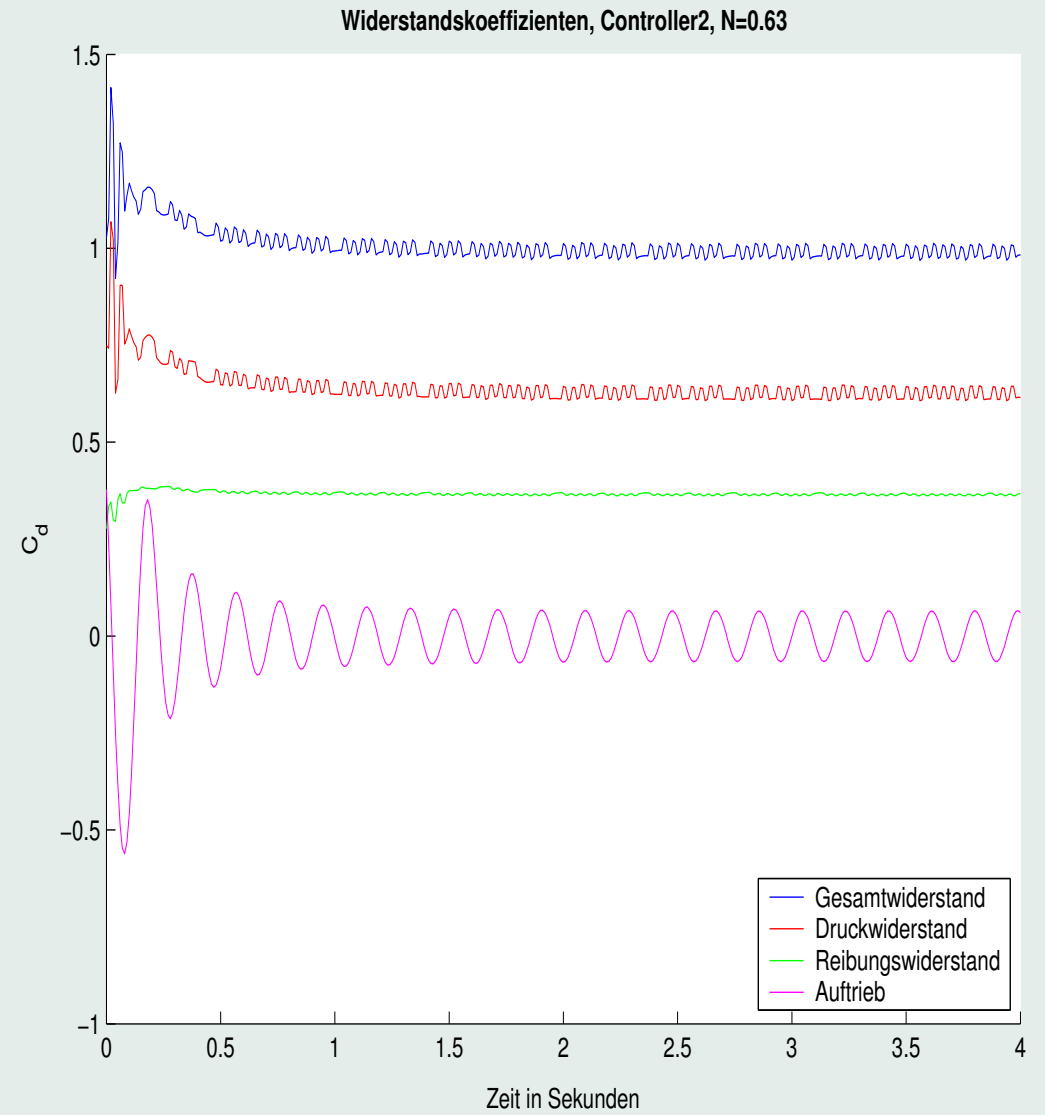
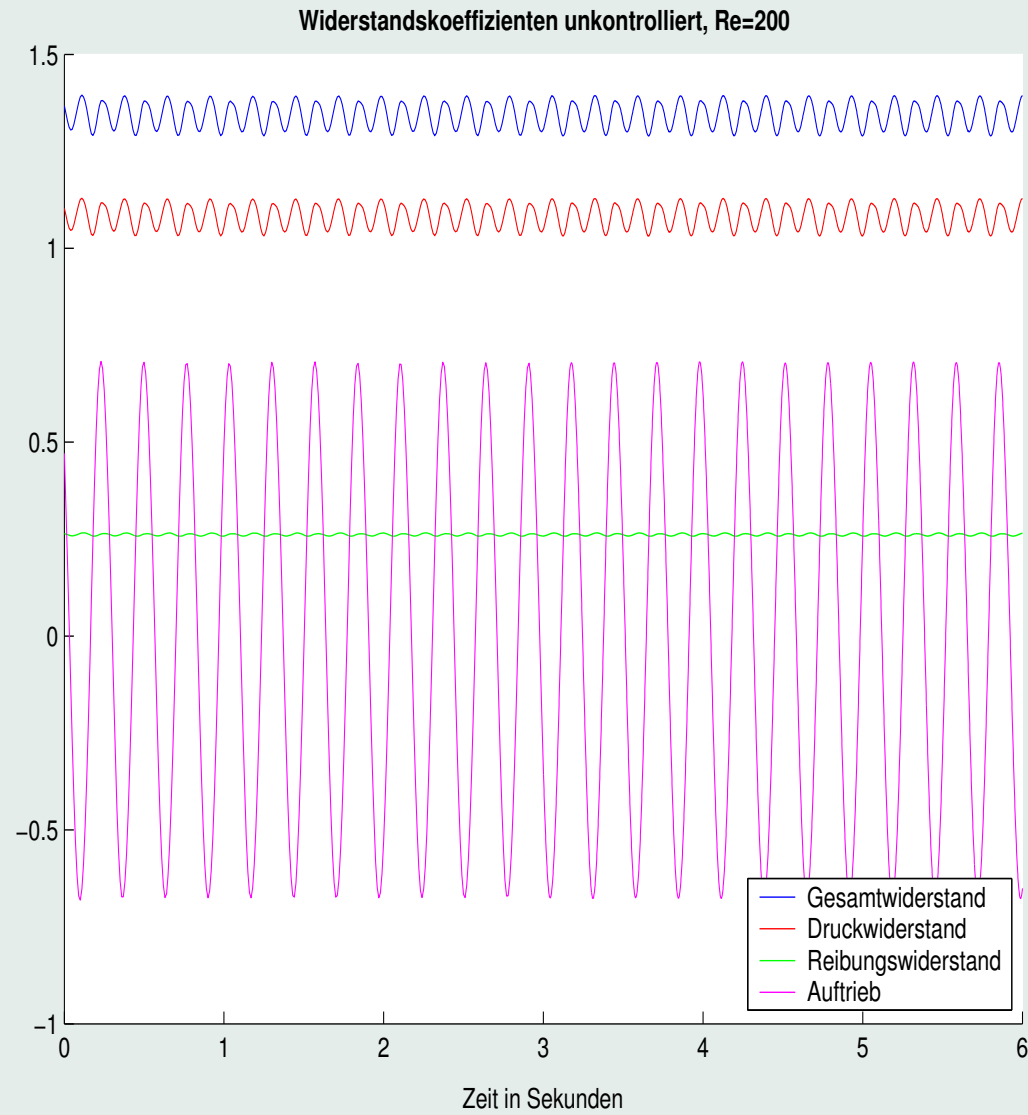
The Lorentz force is switched on, if the fluid flows back in a 10° sector around the rear stagnation point. Cf. Zhihua Chen's Dissertation *Electro-Magnetic Control of Cylinder Wake*



Drag Coefficients:



Time-dependent drag coefficients:



Part2: Nonlinear feedback control recipes

Model predictive control, instantaneous control

Receding horizon control and instantaneous control

Idea of receding horizon approaches

Receding horizon control and instantaneous control

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1 At time t_k compute an **optimal** time discrete control strategy u_{k+1}, \dots, u_{k+l} .

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Suboptimal variant: instantaneous control: Apply only one gradient step to the approximate computation of the optimal strategy for $l = 1$.

Receding horizon control and instantaneous control

Control target

Given some initial state x_0 , find a **control law** $\mathcal{B}u(t) = K(x(t))$ which steers the state $x(t)$ towards a given trajectory \bar{x} :

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Mathematical model:

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$$\begin{aligned}\dot{x}(t) + Ax(t) &= b(x, t) + \mathcal{B}u(t) \text{ state,} \\ y(t) &= \mathcal{C}x(t) \text{ observation,} \\ x(0) &= x_0\end{aligned}$$

Here

- \bar{x} desired stationary state, or
- \bar{x} a reference trajectory obtained from open loop optimal control.

From now onwards: $\mathcal{B} = ID$ und $\mathcal{C} = ID$.

Receding horizon control and instantaneous control

Discretization

Discretize the state equation w.r.t. time

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$$(I + hA)x^{k+1} = x^k + hb^k$$

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$$(P_k) \left\{ \begin{array}{l} \min J(u^{k+1}) = \frac{\gamma}{2}|u^{k+1}|^2 + \frac{1}{2}|x^{k+1} - \bar{x}^k|^2 \\ \text{s.t.} \\ (I + hA)x^{k+1} = x^k + hb^k \end{array} \right.$$

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Time discretization with implicit Euler.

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Feedback oracle

1. Set $x^0 = \phi$, $k = 0$ and $t_0 = 0$.

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$$u^{k+1} = \text{RECIPE}(u_0^k, x^k, \bar{x}^k, t_k)$$

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$$(I + hA)x^{k+1} = x^k + hb(x^k, t_k) + u^{k+1}.$$

4. Set $t_{k+1} = t_k + h$, $k = k + 1$. If $t_k < T$ goto 2.

Receding horizon control and instantaneous control

Instantaneous control

For instantaneous control the oracle RECIPE is given by

$$u = \text{RECIPE}(v, x^k, z, t_k)$$

Receding horizon control and instantaneous control

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- Solve $(I + hA)x = x^k + hb(x^k, t_k) + v$,

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This oracle realizes steepest descent for problem (P_k) .

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$$(I + hA)x = x^k + hb(x^k, t_k) + u$$

$$(I + hA)\lambda = z - x$$

$$\gamma u - \lambda = 0 \quad (\gamma u - \lambda \geq 0) \text{ in case of constraints.}$$

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This oracle realizes solution of problem (P_k) .

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$$u_0^{j+1} = 0.$$

- Instantaneous control rewritten

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This is the semi-discrete version of

$$\dot{x} + Ax = b$$

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- Instantaneous control rewritten

$$(I + hA)x^{k+1} = x^k + hb^k \underbrace{-\rho B^* B(x^k - \bar{x}^k) - h\rho B^* B(b(x^k) - A\bar{x}^k)}_{u^{k+1} =: K_I^d(x^k)}.$$

This is the semi-discrete version of

$$\dot{x} + Ax = b \underbrace{-\frac{\rho}{h} B^* B(x - \bar{x}) - \rho B^* B(b(x) - A\bar{x})}_{\frac{u}{h} =: K_I(x)}, \quad x(0) = x_0.$$

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Receding horizon control and instantaneous control

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Here

$$B := (I + hA)^{-1} \text{ and } S := \gamma(BB + \gamma I)^{-1}BB.$$

Receding horizon control and instantaneous control

Applications

1. $A = \text{Laplacian}$ Heat equation

Receding horizon control and instantaneous control

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Receding horizon control and instantaneous control

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Receding horizon control and instantaneous control

Stability of the closed-loop control operators, distributed control

Receding horizon control and instantaneous control

Stability of the closed-loop control operators, distributed control

Let h be small, \bar{x}, x_0 sufficiently smooth.

The discrete controllers K_I^d and K_O^d generate iterates $\{x^k\}$ which satisfy

$$\|x^k - \bar{x}(t_k)\|_{0,1} \leq c \kappa^k$$

with some $\kappa < 1$.

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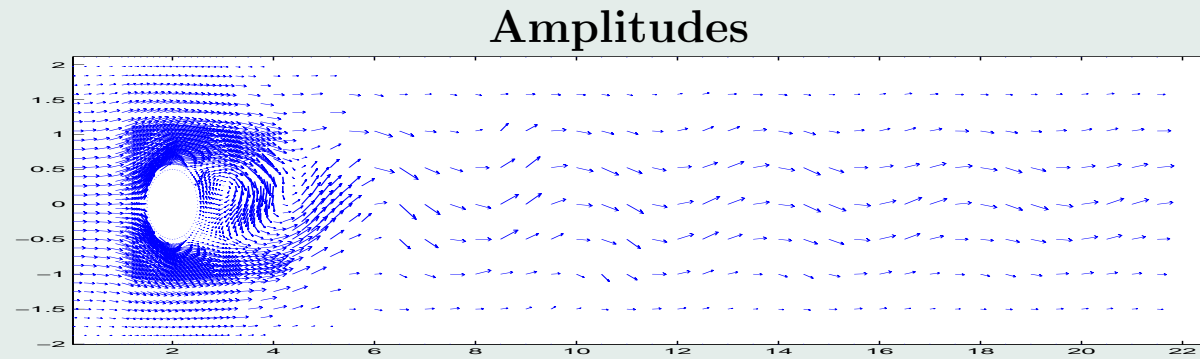
The continuous controllers K_I^d and K_O^d steer the system exponentially fast to \bar{x} :

$$|x(t) - \bar{x}(t)|_{0,1} \leq c \exp\left(-\frac{c}{h}t\right).$$

Control with near wall Lorentz force

Control: $u(t, x) = \sum_{i=1}^n u_i(t) f_i(x)$, $f_i(x) = e^{-|x-x_i|^2}$, x_i point on cylinder surface.

Desired flow \bar{x} : Stokes flow.

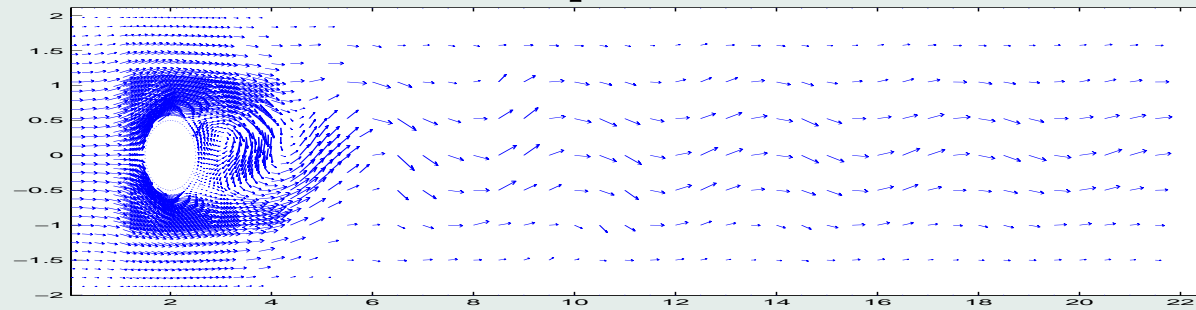


Control with near wall Lorentz force

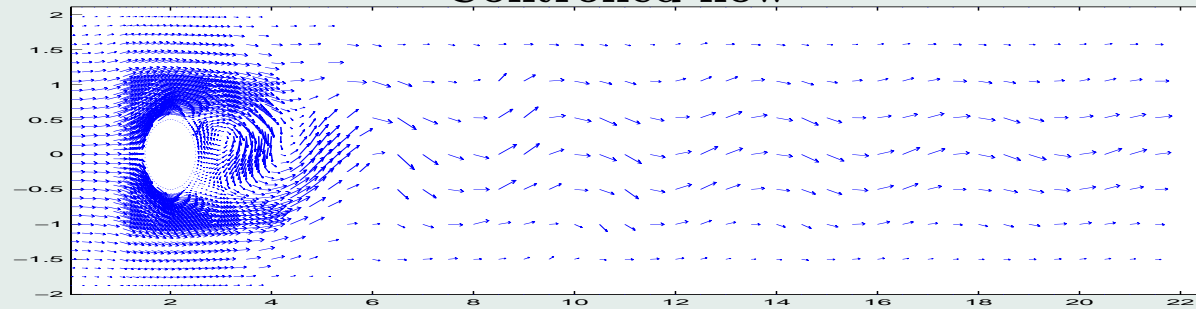
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Controlled flow

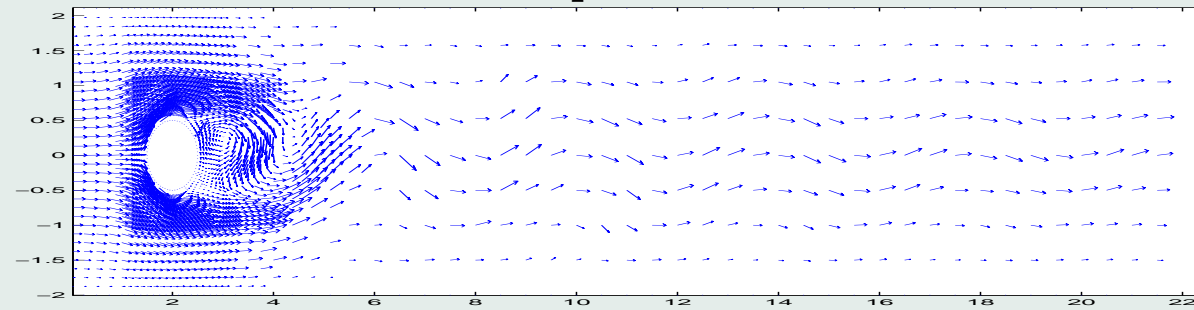


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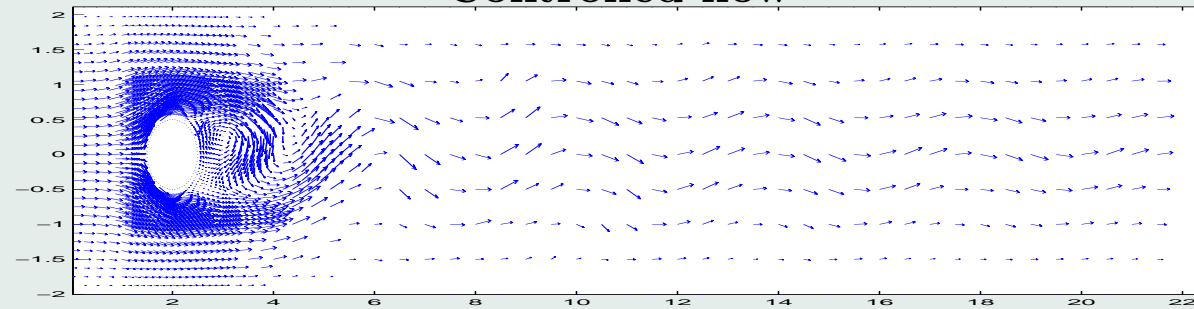
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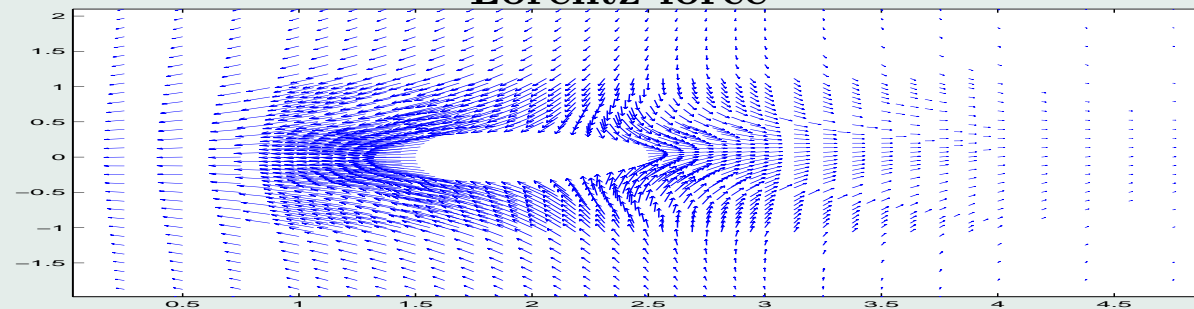
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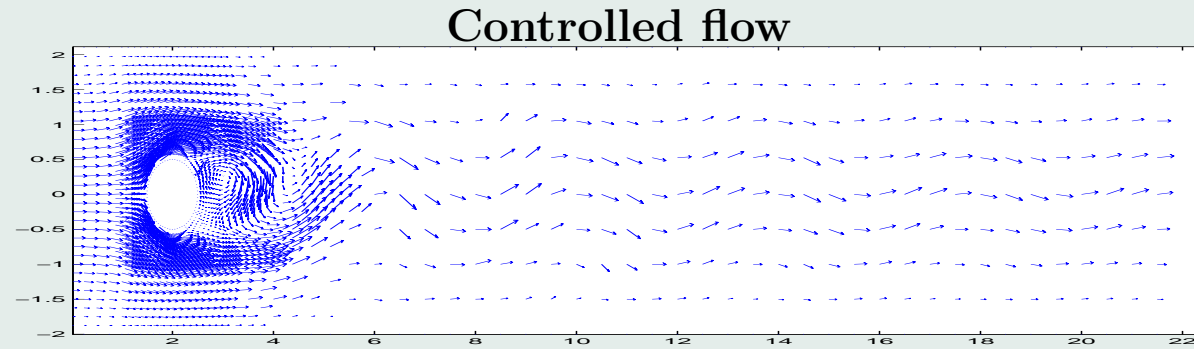
Lorentz force



Control with near wall Lorentz force

$$\text{Control: } u(t, x) = \sum_{i=1}^2 u_i(t) e_{\phi} f_i(x), \quad f_i(x) = g(\phi) e^{-\frac{\pi}{a} \cdot \text{dist}[(x,y), \text{cylinder}]}$$

Desired flow \bar{x} with $x_1 = 1$.

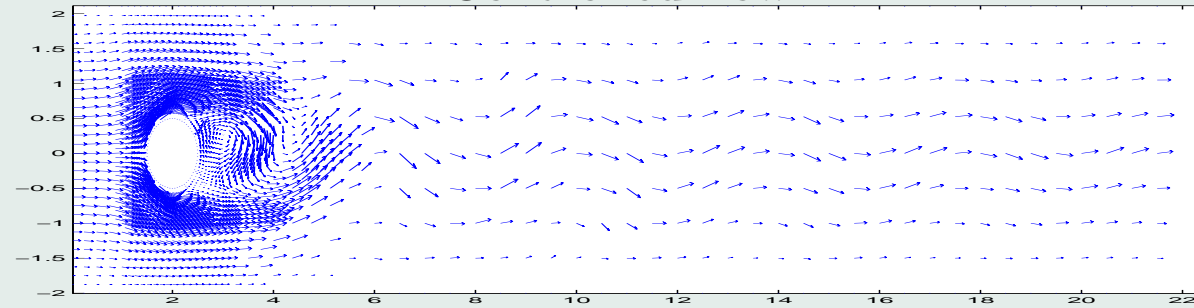


Control with near wall Lorentz force

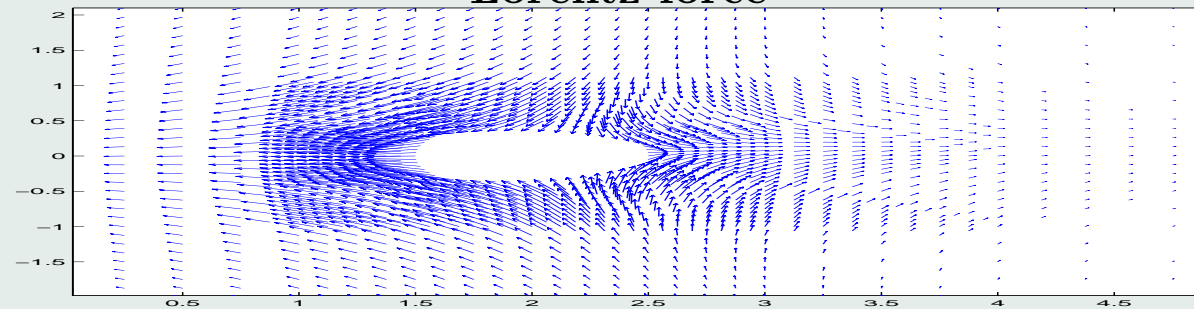
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CPU time needed to compute the instantaneous control strategy = **2,5 times** CPU time needed for one forward solve.

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This amounts to **1-2 %** of the CPU time needed to compute the optimal control trajectory.

Conclusions and future work

2position proportional controller, conducted by experiment

General construction recipe of feedback control operators for pde systems

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General construction recipe of feedback control operators for pde systems

- Instantaneous control
- Model predictive control
- Stability of corresponding controllers (distributed control)

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Apply methods to more complex configurations.

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Extend analysis to boundary control and/or observation

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Extend analysis to boundary control and/or observation

Provide numerical analysis.

Further informations on the topics

<http://www.math.tu-dresden.de/~hinze>